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## COMMUNICATION

# Full Solution Processed Mesostructured Optical Resonators Integrating Colloidal Semiconductor Quantum Dots

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Herein we show a solution based synthetic pathway to obtain a resonant optical cavity with embedded colloidal semiconductor quantum dots (CSQDs). The optical cavity pore network, surrounded by two dense Bragg mirrors, was designed *ad hoc* to selectively host the quantum dots, while uncontrolled infiltration of those in the rest of the layered structure was prevented. Coupling between the optical resonant modes of the host and the natural emission of the embedded nanoparticles gives rise to the fine tuning of the luminescence spectrum extracted from the ensemble. Our approach overcomes, without the need of an encapsulating agent and exclusively by solution processing, the difficulties that arise from the low thermal and chemical stability of the CSQDs. It opens the route to achieving precise control over their location and hence over the spectral properties of light emitted by these widely employed nanomaterials. Furthermore, as the porosity of the cavity is preserved after infiltration, the system remains responsive to environmental changes, which constitutes an added value of the proposed structure.

## COMMUNICATION

**Introduction**

Colloidal semiconductor quantum dots (CSQDs) present intense photoemission that can be precisely tuned by means of particle size and shape, as well as by modification of their capping layer.<sup>1,2,3</sup> Another means of tailoring (spectrally or directionally) the photoluminescence of these nanomaterials is by engineering their photonic environment. In this regard, several works report on the modification of the luminescence of CSQDs infiltrated in porous 2D or 3D photonic crystals (PCs). In the case of 2D PCs, the effect of cavity modes on both emission intensity and lifetime have been investigated by integrating CSQDs either in the bulk of the structure<sup>4,5,6</sup> or in built-in point-like and line defects.<sup>7,8</sup> A variant of this approach makes use of the integration of CSQDs in 1D PC waveguides.<sup>9</sup> In all these cases, sophisticated microfabrication techniques are employed to prepare the luminescent resonators (radio frequency and magnetron sputtering, electron beam lithography, molecular beam epitaxy), and combined with solution processing of CSQDs (in solvents, liquid crystals or polymers). These approaches significantly limit the integration of luminescent optical cavities based on quantum dots over large areas and onto arbitrary composition substrates, hence restricting the impact of the micrometre size structures developed to date to mere proofs of concept. In the case of 3D self-assembled PCs, CSQDs are in general homogeneously distributed throughout the structure and luminescent properties are the results of the coupling to band edges, low dispersion modes or band gaps at high energies.<sup>10-12</sup> Structures present a high density of intrinsic defects that smears the quality and magnitude of the targeted effects.

Interestingly, and although it would be a simpler alternative way to attain large area coatings with tailored photoemission, as it has been demonstrated for other luminescent species,<sup>13-15</sup> very few works address the inclusion of CSQDs in planar 1D PC cavities, i.e., a slab surrounded by two multi-layered mirrors.<sup>16,17</sup> The reason for this is the low compatibility between the commonly employed preparation methods to build a Bragg stack of high dielectric contrast thin films and the usually low thermal and chemical stability of CSQDs. In general, a multilayer structure is built by sequential deposition of the constituent films from either gas, by physical or chemical deposition methods, or liquid phase, by approaches based on sol-gel techniques. Typically, CSQDs do not withstand the chemical reactions occurring during deposition or the later thermal or chemical stabilization of the slabs. Hence, apart from the abovementioned cases, optical cavity multilayers made of liquid crystal films have been used as hosts for CSQDs,<sup>18</sup> this being the only example of full liquid processing of an optical resonator containing CSQDs. Related to this, thin layers of CSQDs have been deposited onto multilayers.<sup>19</sup>

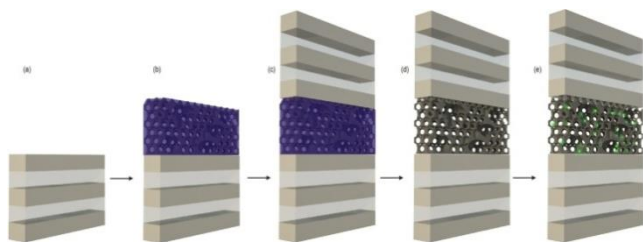
Herein we demonstrate a solution based synthetic route to prepare optical resonators with CdSe/ZnS core-shell colloidal QDs embedded at the desired depth. We make use of recent progress

in the field of supramolecularly templated mesoporous layers<sup>20</sup> to create a highly accessible optical cavity, sandwiched between two dense optical multilayers, with pores large enough as to selectively integrate organically capped CdSe/ZnS nanocrystals *a posteriori*, hence preventing their chemical or thermal degradation. Although inclusion of CSQDs within mesoporous layers can be achieved by *in situ* preparation through adsorption-reaction protocols,<sup>21</sup> this route usually leads to long wavelength emission, irrespective of nanoparticle size, due to surface defects.<sup>22</sup> Our approach overcomes this obstacle by directly incorporating pre-formed nanoparticles with well-defined luminescence properties (i.e., band position and intensity). Infiltration of these CSQDs is achieved by soaking a highly accessible resonator in the corresponding nanocrystal suspension. In this resonator, only the middle defect layer, which acts as optical cavity, presents an open interstitial structure and its pores are designed to host relatively large particles (maximum diameter  $\approx 10$  nm). This improves the selectivity and simplifies the analysis of the optical response caused by variations in the environment with respect to structures in which the CSQDs are homogeneously infiltrated in the whole photonic structure.<sup>11,12,15</sup> The photoemission of CdSe/ZnS nanocrystals integrated in the resonator is strongly modified by the interplay with the cavity modes. A direct correlation is observed between the intensity and shape of the luminescence spectra measured at different collection directions and the angular dependence of the cavity resonant mode. Furthermore, as the middle layer remains open after CSQD infiltration, the resonator responds to changes in the environment with a different luminescent emission, as we herein demonstrate. Our approach opens the route to tailoring light emission of CSQDs by using several square centimetre large nanostructured coatings compatible with current sensing and optoelectronics technology.<sup>23</sup>

**Results and discussion**

The scheme illustrating both the sequential procedure and intermediate layered materials realized to prepare the resonator and its nanostructure is depicted in Figure 1. Periodic layered structures were attained by dip-coating different substrates with silicon dioxide (SiO<sub>2</sub>) and titanium dioxide (TiO<sub>2</sub>) precursors. By this means, we obtain a dense one dimensional photonic crystal,<sup>24-26</sup> which has been shown to be compatible with the integration of mesostructured films to yield more complex structures with devised optical properties.<sup>27,28</sup> In this case, the deposition of an intercalated mesostructured film was realized also by dip-coating a hydrolyzed TiO<sub>2</sub> precursor mixed with a supramolecular template.<sup>29,30</sup> A combination of surfactants, Pluronic F127 and poly(propylene glycol) (PPG), in a THF-butanol solution was chosen. It has been previously proven that this method permits to attain a bimodal, low order, interconnected mesopore network once the organic mold is eliminated.<sup>31</sup> A detailed description of the preparation methods are provided in

the experimental section. We chose this material to form the optical cavity because its reported pore network allowed us foreseeing a suitable diffusion of small particles through it. From the perspective of the optical design, other middle layer compositions would also give rise to effects similar to those herein reported. Quantum dots were arbitrarily selected to emit in the middle of the visible range.



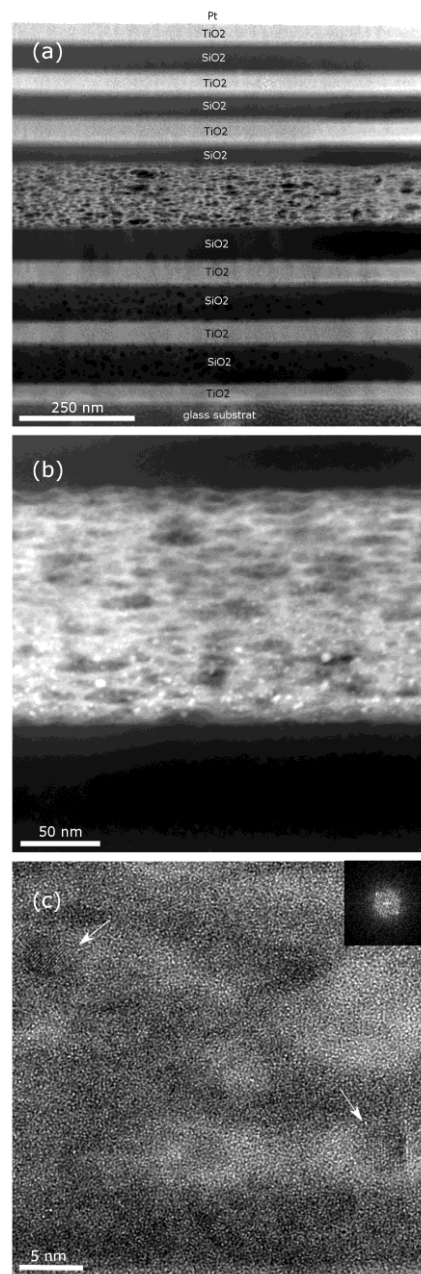
**Fig. 1** Scheme of the process to integrate CSQD in a mesoporous optical resonator. (a) initial multilayer structure, (b) deposition of a bimodal mesostructured films, (c) deposition of the upper multilayer, (d) template removal, and (e) colloidal quantum dot infiltration

Chemical element sensitive annular dark field (ADF) scanning transmission electron microscopy (STEM) and high resolution TEM (HRTEM) images of the cross section of the multilayer made by sandwiching a mesostructured  $\text{TiO}_2$  layer in between two 3-unit cell periodic dense  $\text{SiO}_2$ - $\text{TiO}_2$  layers and a detail of the porous optical cavity are shown in Figures 2.a and 2.b, respectively. The contrast between the electronic density of both types of layers allows to clearly identify them as bright ( $\text{TiO}_2$ ) and dark ( $\text{SiO}_2$ ) fringes in Figure 2.a. Regarding their crystalline phase, X-ray diffraction (XRD) showed that  $\text{SiO}_2$  layers were vitreous while  $\text{TiO}_2$  ones presented anatase structure, with an average crystallite size of a few tens of nanometers. Both layers present a very low porosity (near 3%), as environmental ellipsometric porosimetry (EEP) measurements demonstrate (Figure S1). Crystallization of the originally amorphous  $\text{TiO}_2$  layer gives rise to a small amount of homogeneously distributed micro-cracks. The bimodal pore structure of the middle mesostructured layer can be readily appreciated in Figure 2b. The effect of incorporating PEG and THF into the precursor solution is to produce large, interconnected mesopores (75 nm average diameter) with local ordering, while the smaller pores (15 nm average diameter) present longer range ordering. Most importantly for the goal herein pursued, the interconnecting windows between pores present an average size of  $\approx 10$  nm, significantly larger than the size of the CSQDs employed, whose hydrodynamic diameter nominal value is 8.1 nm including the organic capping. (See supporting information, Figure S2) Details of the EEP analysis and the actual sorption-desorption curves are provided in the experimental and the supporting information sections (Figure S3). Consistently, multilayers in which disordered mesostructured mid layers were built up showed no sharp peaks in the low-angle diffractograms, in agreement with HRTEM images (Figure 2b). The thickness of such mid-layer could be precisely tuned in the range comprised between 50 nm and 150 nm varying either the concentration of the precursor or the substrate withdrawal speed.

So-built mesostructured optical resonators were used as matrices to incorporate luminescent particles from a liquid dispersion.

Please notice that the integration of quantum dots *selectively*

within an optical cavity of the sort herein described cannot be performed during the deposition of the constituent films, since Q-dots are not stable versus the thermal treatments usually needed to endow mechanical robustness to the multilayer structure. An approach to incorporate quantum dots in photonic structures built by sol-gel was previously developed by Jasieniak et al.<sup>32</sup> In that case, an asymmetric optical resonator was completed using a silver mirror to prevent the thermal treatment of the sample once the CSQDs were incorporated into an upper layer deposited onto a dense  $\text{TiO}_2$ - $\text{SiO}_2$  Bragg mirror.



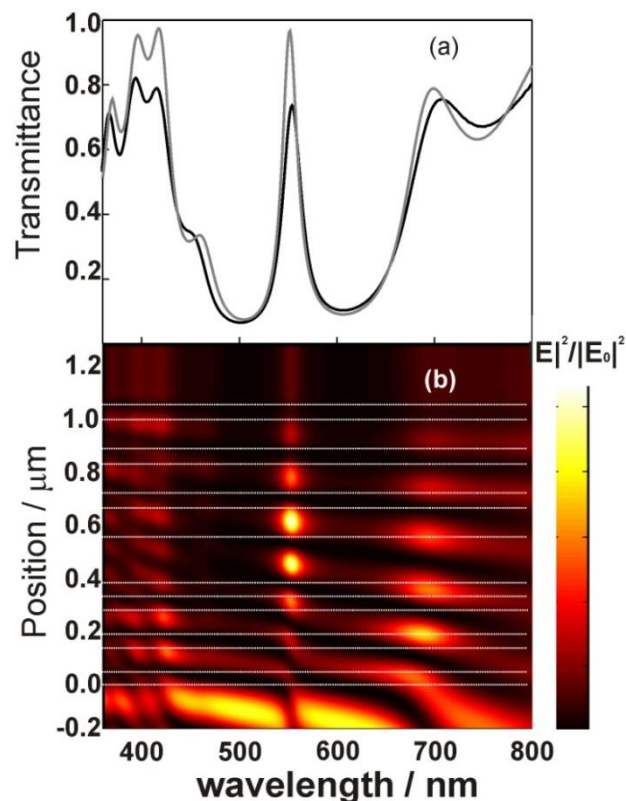
**Fig. 2** Cross sectional ADF STEM images of the (a) complete 13-layers optical resonator (b) ADF image of the porous layer sandwiched in the middle of the optical structure with CSQD appearing as bright spots due to their high atomic number. (c) HRTEM image of one region of the porous layer. The arrows indicate regions, whose image Fourier transform is depicted in the inset, in which crystalline planes of a nanoparticle can be seen.

In our case, we take advantage of the porosity of the resonator to  
 5 infill the CSQDs after the multilayer is finalized. The thicknesses  
 of the constituent films were chosen so that the spectral position  
 of the optical resonance coincides with the emission band of the  
 CdSe/ZnS quantum dots used for this experiment. The luminescent  
 10 nanocrystals were introduced within the mesopores of the middle  
 layer by immersion. After soaking the bimodal pore size distribution  
 resonators overnight in a suspension of said nanoparticles, we  
 found consistent evidence by electron microscopy and photoluminescence  
 15 that diffusion and hence integration of CSQDs in the porous optical  
 cavity has occurred. In Figure 2.b, the location of clusters of CSQDs  
 within the mesostructured middle layer can be seen from the chemical  
 contrast in the ADF STEM image. The HRTEM image in Figure  
 20 2c reveals the crystalline structure of the CSQDs present in the  
 pores. An EDX linescan performed over such particles within the  
 mesostructured layer shows a clear cadmium signal at the particles  
 position (Figure S4).

The loss of Se signal may be attributed to the TEM sample  
 25 preparation, which might cause oxidation or degradation of the  
 delicate CSQDs. Nevertheless, TEM images prove that CSQD are  
 inside the porous optical cavity layer. For the sake of comparison,  
 we prepared mesostructured resonators using only F127 as  
 macromolecular template and subjected them to the same  
 30 infiltration process. In that case, no trace of luminescence was  
 detected from the resonators, which indicated that no luminescent  
 particles were incorporated, in good agreement with the  
 interconnecting window or neck size estimated from the  
 corresponding EEP analysis, which reveals sizes of 8.6nm and  
 35 4.6nm for pores and necks respectively. (Figure S5). So, even  
 when pore sizes in the resonator prepared using only F127 in the  
 precursor solution are large enough to host the CSQDs, pore  
 necks are smaller than the average emitter size and thus  
 infiltration is prevented. This dramatic effect has also been  
 40 observed in the adsorption of bulky enzymes in the whole pore  
 volume of this type of films,<sup>33,34</sup> demonstrating one of the  
 advantages of the use of co-templates towards larger mesopore  
 systems with improved interconnectivity.

Regarding the optical characterization of the resonator, both the  
 45 dense multilayers and the mesostructured middle layer were  
 designed to yield an optical cavity mode at wavelengths matching  
 the photoemission band of the CdSe/ZnS nanocrystals infiltrated  
 in the structure, so the effect of the interplay between both  
 phenomena could be studied. Reflectance maxima result from the  
 50 constructive interference of beams of light in a certain  
 wavelength range reflected at the different interfaces existing  
 between each pair of layers in a periodic structure. The position  
 of this Bragg peak depends on the refractive indices and  
 thicknesses of the layers. An intermediate layer of different  
 55 optical thickness breaks the symmetry of the system and is  
 capable of hosting resonant modes, which can be recognized by  
 the opening of a transmission window at photonic band gap  
 frequencies. In this way, the whole structure behaves as an optical  
 cavity, whose transmittance at normal incidence is plotted in  
 60 Figure 3. From the analysis of the transmittance peak at around  
 $\lambda=540$  nm, we estimate a cavity quality factor of  $\lambda/\Delta\lambda \approx 30$ .  
 Fitting of these spectra using a MATLAB code,<sup>35</sup> allowed us  
 estimating the thickness of each constituent layer, the unit cell of

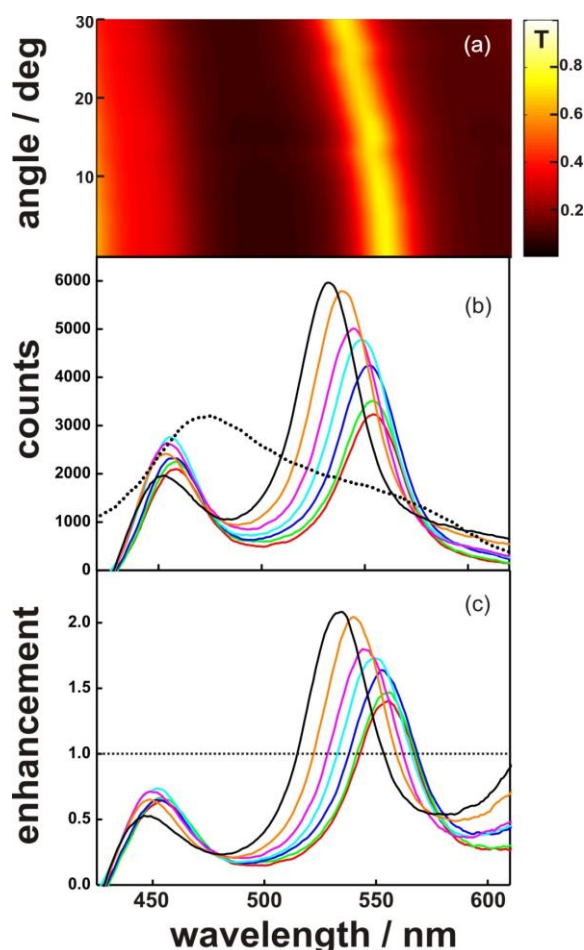
the dense Bragg mirror being composed of  $110 \pm 8$  nm thick SiO<sub>2</sub>  
 60 and  $58 \pm 6$  nm thick TiO<sub>2</sub> layers. The mesostructured middle layer  
 has a thickness of  $180 \pm 15$  nm and an average refractive index of  
 $n=1.70$  in the visible range. The wavelength dependent refractive  
 index and extinction coefficient curves employed for this fitting  
 are provided as supporting information (Figure S7). We use this  
 65 data to calculate the spatial distribution of the squared magnitude  
 of the electric field within the stack when a plane wave impinges  
 in an arbitrary direction onto the multilayer. In Figure 3b we plot  
 the particular  $|E|^2$  attained for normal incidence. Different optical  
 resonances can be identified as bright spots within the multilayer.  
 70 Specifically, field intensity reinforcement is observed at spatial  
 and spectral positions coincident with the photonic band gap  
 edges (approximately at  $\lambda=435$  nm and  $\lambda=665$  nm) and the  
 optical cavity mode (ca.  $\lambda=540$  nm).



75 **Fig. 3** (a) Experimental (black solid line) and theoretical (gray solid line)  
 transmittance of the porous optical resonator. (b) Calculated spatial  
 (vertical axis) and spectral (horizontal axis) distribution of the normalized  
 electric field intensity enhancement ( $|E|^2/|E_0|^2$ ) along a cross section of  
 the mesoporous optical resonator. Calculations were carried out  
 80 considering the multilayer parameters extracted from the fittings to the  
 transmittance measurement. Horizontal dashed lines are guides to the  
 eye to delimit the interfaces between layers in the optical resonator. The  
 bottom of the images is the incoming medium (air) and the top, is the  
 glass side. Incidence is 0 degrees.

85 Photoluminescence (PL) spectra were collected at different exit  
 angles from the CSQD loaded optical resonator using an  
 excitation wavelength of 370 nm, which falls outside the spectral  
 range for which resonant modes are expected, impinging along a  
 fixed incident direction. It should be noted that similar PL spectra  
 90 were taken from spots located at different positions on the  
 resonator surface, including the central part of the film, hence  
 further confirming that the density of CSQDs within the porous

optical cavity is homogeneous as a result of the good lateral diffusion of the particles throughout the middle layer. On the grounds of Lorenz reciprocity principle,<sup>36,37</sup> the PL radiated from an emitter within the multilayer is expected to be enhanced along those directions and wavelengths for which a plane wave hypothetically traveling from the detector to the source would show a high field intensity at the position where the emitter is located. Analogously, the emission will be depleted at those wavelength ranges and along those extraction directions for which the corresponding field distribution display a dark region where the emitter is located.

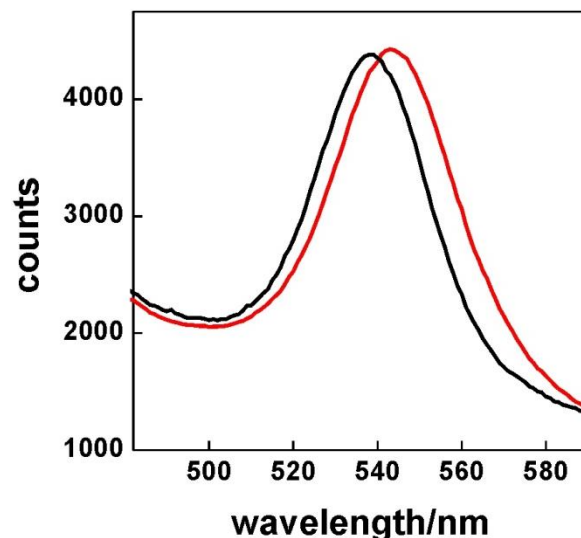


**Fig. 4** (a) Experimental ballistic transmittance spectra of the porous resonator collected at different angles. Measurements were obtained each 2 degrees. Color map picture is built interpolating two consecutive measurements. (b) Luminescence spectra of the porous optical structure collected at 0 (black line), 15 (red line), 20 (blue line), 25 (green line) and 30 (brown line) degrees. Luminescence at 20 degrees obtained for a reference sample is plotted in dashed line. (c) Spectral enhancement of the luminescence at different angles (color code is the same described in Figure 4b)

These effects are actually observed in our measurements and are displayed in Figure 4. In Figure 4a we plot the experimental transmittance of a CSQD loaded resonator *versus* angle of incidence with respect to the multilayer surface normal. As wavelengths at which hot spots appear in the field distribution pattern coincide with those of transmission maxima at each specific angle, PL enhancement is expected to blue-shift with the

angle of collection with respect to surface normal, just like the cavity mode does in Figure 4a. Figures 4b shows the actual PL spectra at different angles of collection and the ratio between the PL of the resonator and that of an unstructured reference sample, respectively. This consists of a CSQD embedded mesostructured film, similar to that sandwiched between the two multilayers in the resonator, deposited on a glass substrate.

Emission reinforcement and depletion spectral regions match bright and dark regions in the transmittance plot of Figure 4a, as expected. For the sake of comparison, the emission of the CSQDs from the isolated mesostructured thin film (dashed line in Figure 4.b) remains approximately constant with the collection angle (Figure S6 in the supporting information). At the same time, PL quantum yield (ratio between the total number of emitted and absorbed photons) of the resonator embedded with CSQDs do not show any significant difference with that of the bare nanocrystals. This behaviour is the one expected for a planar structure for which no lateral confinement has been built up and hence for which no significant modifications of the Purcell factor are expected. In this case, channelling of emitted light along those directions for which there are more available modes does not imply that the intensity of emitted light integrated over all possible directions is larger. Please notice that the reciprocity principle can only provide a qualitative description of the PL enhancement and depletion phenomena observed, and that, in order to fully describe the effect of the inclusion of the CSQDs within the optical cavity on their emission properties, a detailed calculation of the local density of states at each spot from which emission takes place would be required.<sup>38</sup> In this regard, as the LDOS is expected to have a non-constant spectral and spatial profile within the resonator, the recorded PL spectra are inhomogeneous since they are the result of collecting light emitted from spots with different photonic environment.



**Fig. 5** Luminescence spectra of the porous optical structure collected at 0 degrees without (black line) and with (red line) ethanol infiltrated in the structure

Finally, we confirmed that the responsive character of the porous structure against changes in the environment was maintained after QD infiltration, revealing that the void network remained open and accessible. In Figure 5 we plot the photoluminescence spectra

of the CSQD embedded optical resonator when exposed to liquid ethanol (red line). The 6 nm red-shift observed in the emission peak position is the result of the interplay with the spectral shift in the resonant wavelength of the optical cavity, which is in turn due to the increase of the refractive index of the middle layer upon filling with the ethanol flowing from the lateral sides of the film.

## Experimental Section

### Preparation of the porous mesostructured optical resonators.

In order to prepare porous optical resonators, first a periodic structure made of alternate layers of  $\text{TiO}_2$  and  $\text{SiO}_2$  was deposited by dip-coating from sol-gel precursor suspensions. Then, this initial structure was coated with a mesostructured layer by dip-coating a supramolecularly templated precursor dispersion. Finally, a second periodic structure was grown onto the mesoporous layer.

### Synthesis and deposition of dense $\text{TiO}_2$ and $\text{SiO}_2$ layers

$\text{TiO}_2$  and  $\text{SiO}_2$  sols were prepared after procedures reported elsewhere.<sup>24,25</sup> Glass slides cleaned with a  $\text{HNO}_3$  solution (5%) and rinsed with water and ethanol were used. For the synthesis of the  $\text{SiO}_2$  sol, we used 4.5 ml of tetraethylorthosilicate (TEOS, Merck) and 34 ml of absolute ethanol (EtOH). After some minutes under vigorous stirring, 1.72 ml of  $\text{H}_2\text{O}$  and 0.08 ml of  $\text{HCl}$  0.05N were added. To prepare the  $\text{TiO}_2$  sol, 2.5 ml of titanium isopropoxide (TIPT, 97%, Aldrich), 0.156 ml of  $\text{HNO}_3$  (1M), 78  $\mu\text{l}$  of  $\text{H}_2\text{O}$  and 25 ml of isopropyl alcohol (iPrOH, Aldrich) were mixed. The dip-coater employed to deposit the multilayer from these precursors was a HWTL-01-A, from MTI Corporation. The withdrawal speed of the substrates, from the suspension it was immersed in, was varied between 0.5 and 2 mm/s. This permitted us to obtain thicknesses in the range between 50 nm and 200 nm. Stabilization temperature was fixed at 300°C for 1h, and each layer in the stack was subjected to this treatment before a new one was deposited on top of it.

### Synthesis and deposition of mesoporous layers

The porous layer precursor sol was synthesized using a novel recipe that allows to obtain bigger pore size using a low-cost polymer. A quantity of 12.6 grams of a  $\text{TiCl}_4$ :BuOH solution (molar ratio  $\text{Ti}/\text{BuOH}=1:40$ ) was used to dissolve 0.28 g of F127. After 10 minutes of vigorous mixing, 0.72g of  $\text{H}_2\text{O}$  were added and then 0.24 g of polypropylene glycol 4000 (Alfa Aesar GMBH&Co KG) was incorporated to the sol. Finally, 4 ml of tetrahydrofuran (THF) is mixed with the sol during 30 minutes. This precursor suspension was dip-coated onto the first dense multilayer, preserved at 50% relative humidity during 24h and then sequentially heated at 60°C (12h), 120°C (12h) and 200°C (2h). Once the mesostructure was stabilized, the second multilayer was deposited onto it following a similar method than the one employed to prepare the first one. The whole structure was thermally treated at 350°C for 1h (ramp 1°C/min) to remove the supramolecular compound used as template.

Porous mesostructured resonators were soaked in a toluene solution of CdSe/ZnS quantum dots (LumidotTM480, Aldrich) during 24 hours. After that, the samples were rinsed with fresh

toluene to eliminate excess of quantum-dots. Monolayers of same thickness as the middle layer used to build the optical cavity and possessing a similar porous mesostructure were treated in the same way in order to obtain a luminescent film whose emission can be compared to that of the resonator.

### Structural characterization

Pore and neck size of individual films were characterized through environmental ellipsometric porosimetry (EEP). A SOPRA GES5A apparatus was used to attain water adsorption-desorption isotherms at 298 K by EEP. From the fitting of the ellipsometric parameters in the visible range (400-800 nm) we estimated film thickness and refractive index (real part). The film refractive index was described according to a modified Cauchy equation. The WinElli 2 software (SopraInc) is employed to convert variations of  $n$  with  $P/P^\circ$  into filled pore volume. Kelvin equation is used to extract pore and neck size distributions.<sup>29</sup> For TEM analysis, a Lamella was cut with a FEI Helios Nanolab 400s with a subsequent milling in a Fischione Nanomill to remove destroyed surface layers. TEM and STEM experiments were conducted at FEI Tecnai F20 (ER-C) operated at 200 kV.

### Analysis of the optical properties and their environmental dependence.

Transmittance spectra were obtained using an UV-VIS scanning spectrophotometer (UV-2101PC, Shimadzu). Transmission angle measurements were acquired supporting the sample onto a rotating stage (part number M55-028, Edmund Optics). Excitation and emission spectra were collected using a spectrofluorometer (Fluorolog-3, Horiba JobinYvon) with a Xe450 W lamp source. PL experiments were made collecting light in front face configuration with a fixed angle of 22.5° between the entrance and the exit slits. Sample was placed in the same rotating stage described above to obtain the emission at different angles. Ethanol-soaking experiments were made by placing a drop of ethanol onto the sample and pressing it gently with a cover slide to force infiltration.

Samples were constructed on a 75 mm x 25 mm low fluorescence glass. Metal oxide precursors liquid dispersions used to fabricate dense and porous layers cover an 80% of that area. Then samples were cut in three equal parts (regions near to the borders of the glass were discarded) and immersed in q-dots toluene dispersions. Final size of the samples is 15mm x 15mm. In all the cases, different series of photoluminescence measurements were taken from the side to the center of the sample to probe the degree of penetration of the SCQDs throughout the middle porous layer. We find that luminescence is not dependent of the region of the sample measured, thus proving the homogeneous lateral infiltration of the resonator.

## Conclusions

We have demonstrated a synthetic route to selectively integrate colloidal semiconductor quantum dots in a porous optical cavity surrounded by dense photonics crystals. This is, to the best of our knowledge, the first time that precise spatial confinement of colloidal quantum dots in a resonator is achieved, as their incorporation in this configuration has been prevented before due to the low chemical and thermal stability of the organically

capped nanocrystals. In order to do so, the mesostructure of the resonator was carefully devised to permit the efficient lateral infiltration of the nanocrystals by soaking. We demonstrate that the photoemission spectrum of the confined quantum dots is determined by the cavity resonance, which opens the door to precise tailoring of the optical properties of this type of nanocrystals by confinement in porous matrices in which optical resonances are built up. We foresee our work may ease the integration of colloidal semiconductor quantum dots into new light emitting coatings with improved spectral and directional properties.

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### Notes and references

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<sup>†</sup> Footnotes should appear here. These might include comments relevant to but not central to the matter under discussion, limited experimental and spectral data, and crystallographic data.

Electronic Supplementary Information (ESI) available: [details of any supplementary information available should be included here]. See DOI: 10.1039/b000000x/

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