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We computationally investigate the hydrogen storage properties of calcium-decorated C₄₈B₁₂ boron-carbon heterofullerene molecules, and compare them to C_{60} (allcarbon) fullerene decorated with calcium. We employ density functional theory (DFT) on the lowest energy configurations of the $C_{48}B_{12}$ molecules and find that these molecules have the following properties. (1) The most stable $C_{48}B_{12}$ isomers have an electron affinity that is 0.93-1.04 eV higher than their carbon only counterpart. (2) The binding of a Ca atom to $C_{48}B_{12}$ is ~ 2.2 eV stronger than its binding to C₆₀. (3) Unlike C₆₀Ca_x, x = 1-6, C₄₈B₁₂Ca_x is stable with respect to decomposition into the fullerene molecules and Ca bulk metal. (4) $C_{48}B_{12}Ca_x$ binds up to six hydrogen molecules per metal center, leading to a gravimetric density of up to 7.1 weight percent (wt %). The hydrogen binding energies of up to \sim 0.24 eV open a prospect of hydrogen storage at ambient temperature.

Hydrogen is a clean energy carrier that shows promise to decrease our dependence on fossil fuels. The US Department of Energy (DOE) suggests that by the year 2020 a successful onboard hydrogen storage system has a hydrogen storage capacity of at least 5.5 wt % at moderate pressures, and should operate at delivery temperatures between -40 and 85° C.¹ The (un)loading of hydrogen (from)to the storage material should be at least as fast such that 5 kg of hydrogen should be loaded within 200 seconds.¹ Physisorption is preferred to chemisorption as no bonds need to be made or broken during the process.

Physisorption materials with open or porous structures allow for fast hydrogen kinetics, which facilitates loading hydrogen into the storage system. In physisorption materials, H₂ molecules are adsorbed onto substrates, typically via weak, van der Waals type, interactions. Thus, binding hydrogen molecules to the host materials requires very low temperatures and high hydrogen gas pressures. Such conditions are generally unsuitable for mobile applications, which require safe and practical means of storing hydrogen.

To improve the binding of hydrogen to physisorption materials, recent research has focused on dispersing metal atoms onto these materials. A stronger interaction with the hydrogen molecules has important consequences for the thermodynamics, and increases the operating temperatures of the storage materials. To restrict the increase of the total system weight, resulting from the addition of metals, only a limited number of metal atoms can be used to decorate lightweight hosts. Atoms of the simple, electropositive alkali or alkaline earth metals, form very polar bonds with carbon-based (or boron-based) materials. The electrostatic interactions between the (partially charged) metal atoms and the induced dipoles or multipoles on the hydrogen molecules result in a binding energy that is 2-3 times larger than a typical van der Waals bond between hydrogen and a carbon-based material.^{2,3} Alkaline earth metals, such as Ca, have a larger effect on the binding energy than alkali metals, such as Li or Na. Additionally, recent studies show that functionalizing lightweight hosts with Ca atoms has important benefits. 4-19

Dispersing transition metal atoms, such as Sc or Ti, on carbon-based substrates makes the adsorption energy of hydrogen molecules even larger.²⁰⁻²⁴ However, not all hydrogens are then bonded equally. Upon loading, the first hydrogen molecules are strongly chemisorbed in atomic form, and the subsequent ones are physisorbed in molecular form. This means that the full storage capacity can not be addressed at a single operating temperature. Moreover, clustering of metal atoms will almost certainly occur, because of the relatively large cohesive energy of these transition metals (≥ 4 eV/atom). This, in turn, degrades the storage capacity even further.²⁵ The Ca metal has a relatively small cohesive energy $(\leq 2 \text{ eV/atom})$, so Ca atoms have smaller tendency to cluster on the host material once they are deposited. Ca atoms are expected to bind all hydrogen molecules reversibly via phvsisorption. Moreover, Ca is an abundant element, which is

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relatively low-cost when compared with transition metals such as Sc or Ti.

The experimental and theoretical studies done so far substantially explored carbon-based systems (fullerenes, nanotubes, graphene, graphyne and carbyne).^{7–19} In this work, we study the Ca decoration of carbon-boron heterofullerenes. The first heterofullerenes realized in experiments were borondoped carbon fullerenes.^{26,27} Since then, a variety of heterofullerenes with different boron concentrations have been produced on a macroscopic scale.^{28–37} We show that incorporating boron stabilizes the Ca-decorated fullerene, and also increases the interaction of hydrogen molecules with this system. The optimal interaction strength can be obtained from the rationalization that follows.

The adsorption of molecules, such as H_2 , under common conditions does not usually involve a first-order phase transition. This means that when operating at a fixed temperature T, the storing and unloading of hydogen have to be done at two different gas pressures P_1 and P_2 , respectively. Assuming that physisorption of hydrogen can be described by the Langmuir adsorption model,³⁸ the fraction f of the total adsorption capacity that is reversibly loaded and unloaded at these pressures is maximally

$$f = \frac{\sqrt{P_1/P_2 - 1}}{\sqrt{P_1/P_2 + 1}},\tag{1}$$

if the change in free energy ΔG_0 , which determines the adsorption reaction constant, is equal to

$$\Delta G_0 = \Delta H_0 - T \Delta S_0 = \frac{RT}{2} \ln \left[\frac{P_1 P_2}{P_0^2} \right], \qquad (2)$$

where $P_0 = 1$ bar is the standard pressure. Using Monte Carlo simulations, Bhatia and Myers³⁸ have argued that for hydrogen adsorption on a range of carbon substrates the entropy change at standard pressure $\Delta S_0 \approx -8R$. It means that, e.g., at room temperature for $P_1 = 30$ bar and $P_2 = 1.5$ bar, the optimum adsorption enthalpy is $\Delta H_0 \approx -0.15$ eV/ H₂ molecule. As intrinsic properties, such as vibrations, of hydrogen molecules hardly change upon physisorption, one may approximate ΔH_0 by the T = 0 hydrogen binding energy, and neglect the contributions of zero-point motions. On substrates other than the carbon based materials, a variance in the entropy change at standard pressure (ΔS_0) shifts the enthalpy change at standard pressure (ΔH_0). Recent experiments on porous solids with cationic metal centers, such as cation-exchanged zeolites and metal-organic frameworks with coordinatively unsaturated cationic metal atoms, suggest a positive and nonlinear correlation between ΔS_0 and ΔH_0 . These new studies estimate the ideal value of hydrogen binding energies $(-\Delta H_0)$ to be between 0.23 and 0.26 eV/H₂.^{39,40} Accordingly, for the materials considered in the present work we aim for molecular a)

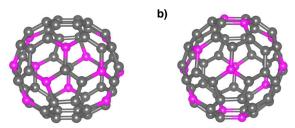


Fig. 1 Perspective views of DFT optimized $C_{48}B_{12}$ heterofullerenes: (a) the S_6 isomer, and (b) the C_i isomer. Carbon and boron atoms are shown as grey and magenta spheres, respectively.

hydrogen binding energies that are between 0.15 and 0.26 eV/ H_2 .

We calculate such adsorption energies within density functional theory (DFT) calculations with the PW91 functional of the generalized gradient approximation (GGA),⁴¹ using the Vienna Ab Initio Simulation Package (VASP).⁴²⁻⁴⁴ Projector augmented waves (PAW)^{45,46} are employed, treating for B and C the electrons in the 1s shell as frozen core, and for Ca all electrons up to and including the 3s shell. The PAW technique and the GGA-PW91 functional are sufficiently accurate for studying the interaction between molecular hydrogen and metals.^{2,3,6,47} To account for the dispersion (van der Waals) interactions we also carry out calculations using the DFT-D2 method of Grimme with its default forcefield parameters as implemented in VASP.⁴⁸ Grimme's empirical correction scheme to common density functionals has been applied successfully to many chemical problems including molecular hydrogen.^{4,49-54} We find however that Ca looses part of its valence electrons. The resulting Ca ion is more difficult to polarize, hence its van der Waals interaction with molecular hydrogen is probably weakened. Thus, PW91 results can be taken as lower bounds for hydrogen binding energies whereas DFT-D2 results can be taken as upper bounds. The kinetic energy cutoff for the plane waves is set to 400 eV and the conjugate gradient (CG) algorithm is used for the geometry optimizations. All molecules are optimized using a periodic cubic box and Γ -point sampling. A sufficiently large cell parameter of 30 Å is used to avoid significant interactions between the periodic images. The optimizations are assumed to be complete when the total remaining forces on the atoms are lower than 0.01 eV/Å. In the following we report the calculated numbers using the DFT-D2 method of Grimme. When discussing small binding energies, where van der Waals interactions could be significant, we report both DFT-D2 and PW91 results .

We calculate the binding energy $E_{\rm b}^{\rm Ca}$ of a Ca atom to a (hetero)fullerene molecule from

			$E_{\rm b}^{\rm H_2}({\rm eV})$					
isomer	$E_{\rm b}^{\rm Ca}({\rm eV})$	Q(e)	1 H ₂	$2 H_2$	3 H ₂	$4 H_2$	5 H ₂	6 H ₂
S ₆	1.616	+1.47	0.175	0.163	0.164	0.156	0.151	0.152
$S_6(DFT-D2)$	1.564	+1.52	0.265	0.255	0.255	0.249	0.246	0.245
C_i	1.626	+1.46	0.161	0.161	0.154	0.156	0.149	0.149
$C_i(DFT-D2)$	1.575	+1.51	0.243	0.243	0.239	0.246	0.243	0.242

Table 1 Calculated Binding Energies of Ca atoms to $C_{48}B_{12}$ isomers with Respect to Bulk Ca Metal (E_b^{Ca}) , Bader Charges on the Ca Atoms (Q), Average Binding Energies $(E_b^{H_2})$ of H₂ molecules to $C_{48}B_{12}$ Ca (see Fig. 2).

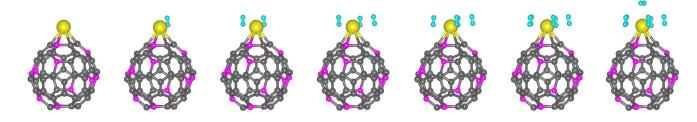


Fig. 2 Optimized structures of $C_{48}B_{12}Ca(H_2)_y$, y = 1-6 of the S₆ isomer. Carbon, boron, calcium and hydrogen atoms are shown as grey, purple, yellow and turquoise spheres, respectively.

$$E_{\rm b}^{\rm Ca} = \frac{1}{x} \left[E_{\rm C_{48}B_{12}} + x E_{\rm Ca_{(bulk)}} - E_{\rm C_{48}B_{12}Ca_x} \right], \tag{3}$$

where $E_{C_{48}B_{12}}$ is the total energy of a (hetero)fullerene molecule, $E_{Ca_{(bulk)}}$ is the total energy per atom of calcium bulk metal, and $E_{C_{48}B_{12}Ca_x}$ is the total energy of a (hetero)fullerene molecule decorated with x Ca atoms. Note that a positive value of E_b^{Ca} means that the Ca-decorated (hetero)fullerene molecule is stable.

We first consider the interaction of a single Ca atom with a (hetero)fullerene molecule. We find that the Ca atom does not bind strongly to carbon-only fullerene. Using the total energy $E_{Ca_{(atom)}}$ of an isolated Ca atom instead of the bulk energy $E_{Ca_{(bulk)}}$ in Eq. 3, we find a binding energy of 1.437 eV, in accordance with a previous study.¹⁹ However, if we use $E_{Ca_{(bulk)}}$ as a reference energy, as is more appropriate, then $E_b^{Ca} = -0.666$ eV. Such a negative value indicates that Ca-decorated C₆₀ molecules are thermodynamically not stable against decomposition into C₆₀ and Ca bulk. Therefore, there is an intrinsic driving force towards the clustering of Ca atoms, which severely impairs the potential use of C₆₀Ca as a hydrogen storage material.

The binding between Ca and C_{60} is polar. If we perform a Bader charge analysis ^{55,56} to estimate the effective charges on the atoms, we find that the Ca atom attains a charge of +1.38e. To increase the binding between Ca and fullerene, we consider substitution of fullerene carbons with boron atoms in the latter. This creates an electron-deficient molecule and hence a bet-

ter acceptor than C₆₀. Previous studies have shown such substitutions are possible up to boron concentrations of 20 atom %, while retaining the fullerene structure.^{26,28–30,57–60} C₄₈B₁₂ has two low energy isomers of S₆ and C_i symmetry, respectively.^{57,58,60} We find that the S₆ isomer is 0.836 eV/molecule lower in energy that the C_i isomer, making the former thermodynamically much more stable. The increased acceptor character of these boron-substituted fullerenes is illustrated by their electron affinity. The calculated electron affinity of C₆₀ is 2.77 eV, in good agreement with the earlier theoretical⁵⁷ (2.40 eV) and the experimental⁶¹ (2.69 eV) results. The calculated electron affinities for the S₆ and C_i isomers of C₄₈B₁₂ are 3.70 eV and 3.81 eV, respectively.

Both of the $C_{48}B_{12}$ isomers have pentagons containing a single B atom (Fig. 1). The S₆ isomer contains eight all-carbon and twelve C_4B_2 hexagons (Fig. 1a). Six of the latter contain B in para positions, whereas the others contain B in meta positions. The C_i isomer contains two all-carbon, twelve C_5B_1 , and six C_4B_2 hexagons (Fig. 1b). All six of the latter have B in para positions.

We find that the most favorable binding site of Ca is over the C₄B₂ hexagon centers, both in the S₆ and the C_i isomers, as shown in Fig. 2. The calculated binding energies of a single Ca atom, x = 1 in Eq. 3, are 1.564 eV and 1.575 eV to the S₆ and C_i isomers, respectively. These values clearly demonstrate that binding of Ca to C₄₈B₁₂ is much stronger than to C₆₀. In contrast to C₆₀Ca, C₄₈B₁₂Ca is stable with respect to clustering of Ca atoms, i.e. separation into Ca bulk and C₄₈B₁₂ molecules. Similar to C₆₀Ca, we perform a Bader charge anal-

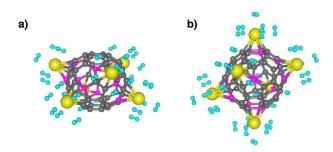


Fig. 3 Perspective views of the fully optimized Ca-decorated heterofullerenes in their fully hydrogenated states, $C_{48}B_{12}Ca_6(H_2)_{36}$; (a) and (b) correspond to the S₆ and C_i isomers, respectively.

ysis on $C_{48}B_{12}Ca$. The calculated Bader charges on Ca are +1.52e and to +1.51e for the S₆ and C_i isomers, respectively, indicating that the bonding has a substantial ionic character, and that the two isomers behave very similar regarding their bonding to Ca.

Next, we study the adsorption of H_2 molecules on $C_{48}B_{12}Ca$. We calculate the binding energy of hydrogen molecules to the $C_{48}B_{12}Ca_x$ using Eq. 4.

$$E_{\rm b}^{\rm H_2} = \frac{1}{y} \left[E_{\rm C_{48}B_{12}Ca_{\rm x}} + y E_{\rm H_2} - E_{\rm C_{48}B_{12}Ca_{\rm x}({\rm H_2})_{\rm y}} \right], \qquad (4)$$

where $E_{\rm H_2}$ and $E_{\rm C_{48}B_{12}Ca_x(\rm H_2)_y}$ are the total energies of the hydrogen molecule and the stepwise hydrogenated (hetero)fullerene-Ca system, respectively. A single Ca atom (x = 1) can bind up to six H₂ molecules, as illustrated in Fig. 2, with five H_2 in a ring around the protruding Ca atom, and the sixth H₂ molecule on top of the Ca atom. The first H_2 molecule is bonded with an energy of 0.265 eV/ H_2 and 0.243 eV/H₂, to the Ca-decorated S_6 and C_i isomers of C₄₈B₁₂, respectively. Additional H₂ molecules bind somewhat less strongly, see Table 1. At their fully hydrogenated states, the Ca-decorated S₆ and C_i isomers bind the hydrogen molecules with average binding energies of 0.245 eV/H_2 and 0.242 eV/H₂, respectively. The calculations using the PW91 functional give similar results to the DFT-D2 method, but the average hydrogen binding energies are approximately 0.1 eV lower for both of the isomers (Table 1).

As $C_{48}B_{12}$ is isoelectronic with C_{60}^{12+} , and the bonding with Ca atoms has a substantial ionic character, one expects that $C_{48}B_{12}$ accommodate six Ca atoms with a similar binding. This is achieved by putting a Ca atom over the centers of all $C_{4}B_{2}$ hexagons, both in the S₆ and the C_i isomers. Each of these Ca atoms is bonded strongly to $C_{48}B_{12}$ with average binding energies of 1.046 eV/atom and 1.082 eV/atom for the S₆ and C_i isomers, respectively (Eq. 3). Note that, since these binding energies are calculated with respect to Ca bulk, we ex-

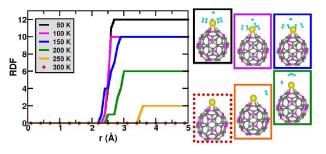


Fig. 4 The radial distribution functions (RDFs) of Ca–H at the end of 15 ps *ab initio* MD equilibration runs using the DFT-D2 method at 50, 100, 150, 200, 250 and 300 K for the S₆ isomer based $C_{48}B_{12}Ca(H_2)_6$ molecule. The boxes shown at the right show the structures at the end of each run.

pect $C_{48}B_{12}Ca_6$ to be stable against the clustering of Ca atoms, which is again in contrast to Ca-decorated C_{60} .

The optimized structures of the fully hydrogenated $C_{48}B_{12}Ca_6$ molecules are shown in Fig. 3. As for $C_{48}B_{12}Ca_6$, we find that a maximum of six H_2 molecules are adsorbed per Ca atom. Consequently, the fully hydrogenated molecules have the chemical formula $C_{48}B_{12}Ca_6(H_2)_{36}$. Each hetero-fullerene molecule then contains 7.1 wt % H_2 with average binding energies of 0.245 eV/ H_2 (0.152 eV/ H_2 with PW91) and 0.242 eV/ H_2 (0.149 eV/ H_2 with PW91), for the S₆ and C_i isomers, respectively.

For physisorbed molecules one may expect that these energies also determine the kinetics, and that no additional kinetic energy barriers are involved in the adsorption process. To illustrate this we carried out first-principles molecular dynamics (MD) simulations on the $C_{48}B_{12}Ca(H_2)_6$ molecule in the micro canonical ensemble.⁶² Due to the presence of lightweight elements, we used a short time step of 1.0 fs in the MD simulations. Using the average kinetic energy of the atoms as a measure for the temperature, the latter is stepwise increased from 50 K to 300 K in 50 K steps, where at each temperature the simulations are performed for 15 ps. Fig. 4 shows the radial distribution function for Ca-H pairs and the structures at the end of each equilibration step. In the simulation run, five hydrogen molecules stay adsorbed for temperatures \leq 150 K, and all are desorbed for temperatures \geq 300 K. This behavior is consistent with the absence of any kinetic barrier for desorption.

In summary, we studied the functionalization of boron heterofullerenes with Ca atoms for high density hydrogen storage. We find that six Ca atoms interact strongly with $C_{48}B_{12}$, and the dispersed Ca atoms are thermodynamically stable against metal clustering, unlike Ca atoms dispersed on C_{60} . The bonding between Ca and $C_{48}B_{12}$ has a substantial ionic character, and the well-exposed, (partially) charged Ca atoms bind up to six H_2 molecules per Ca with an average binding

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energy of $\sim 0.24 \text{ eV/H}_2$ (DFT-D2) and $\sim 0.15 \text{ eV/H}_2$ (PW91). These energies are in the range needed for onboard hydrogen storage, yielding 7.1 wt % H₂. Furthermore, the predicted hydrogen binding energies are only weakly dependent on H₂ loading.

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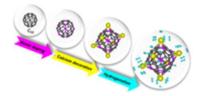
References

- 1 US Department of Energy Fuel Cell Technologies Office Multi-Year Research, Development and Demonstration Plan, http: //energy.gov/sites/prod/files/2014/11/f19/fcto_ myrdd_storage.pdf, Accessed: 2014-12-11.
- 2 S. Er, G. A. de Wijs and G. Brocks, J. Phys. Chem. C, 2009, 113, 8997– 9002.
- 3 S. Er, G. A. de Wijs and G. Brocks, J. Phys. Chem. C, 2009, 113, 18962– 18967.
- 4 X. Chen, F. Yuan, Q. Gu and X. Yu, *Journal of Materials Chemistry A*, 2013, **1**, 11705.
- 5 F. Gao, Z. Ding and S. Meng, Sci. Rep., 2013, 3, 1882.
- 6 N. Park, K. Choi, J. Hwang, D. W. Kim, D. O. Kim, and J. Ihm, Proc. Natl. Acad. Sci. U.S.A., 2012, 109, 19893–19899.
- 7 T. Hussain, B. Pathak, M. Ramzan, T. A. Maark and R. Ahuja, *Appl. Phys. Lett.*, 2012, **100**, 183902.
- 8 P. B. Sorokin, H. Lee, L. Y. Antipina, A. K. Singh and B. I. Yakobson, *Nano Lett.*, 2011, **11**, 2660–2665.
- 9 C. Li, J. Li, F. Wu, S.-S. Li, J.-B. Xia and L.-W. Wang, J. Phys. Chem. C, 2011, **115**, 23221–23225.
- 10 C. Cazorla, S. A. Shevlin and Z. X. Guo, Phys. Rev. B, 2010, 82, 155454.
- 11 Q. Wang, Q. Sun, P. Jena and Y. Kawazoe, J. Chem. Theory Comput., 2009, 5, 374–379.
- 12 G. Wu, J. Wang, X. Zhang and L. Zhu, J. Phys. Chem. C, 2009, 113, 7052–7057.
- 13 G. Kim and S.-H. Jhi, J. Phys. Chem. C, 2009, 113, 20499-20503.
- 14 M. Li, Y. Li, Z. Zhou, P. Shen and Z. Chen, *Nano Lett.*, 2009, 9, 1944– 1948.
- 15 H. Lee, J. Ihm, M. L. Cohen and S. G. Louie, *Phys. Rev. B*, 2009, 80, 115412.
- 16 H. Lee, J. Ihm, M. L. Cohen and S. G. Louie, *Nano Lett.*, 2010, **10**, 793–798.
- 17 X. Yang, R. Q. Zhang and J. Ni, Phys. Rev. B, 2009, 79, 075431.
- 18 G. Kim, S.-H. Jhi, S. Lim and N. Park, Phys. Rev. B, 2009, 79, 155437.
- 19 M. Yoon, S. Yang, C. Hicke, E. Wang, D. Geohegan and Z. Zhang, *Phys. Rev. Lett.*, 2008, **100**, 206806.
- 20 T. Yildirim and S. Ciraci, Phys. Rev. Lett., 2005, 94, 175501.
- 21 Y. Zhao, Y.-H. Kim, A. C. Dillon, M. J. Heben and S. B. Zhang, *Phys. Rev. Lett.*, 2005, **94**, 155504.

- 22 H. Lee, W. I. Choi and J. Ihm, Phys. Rev. Lett., 2006, 97, 056104.
- 23 G. Kim, S.-H. Jhi and N. Park, Appl. Phys. Lett., 2008, 92, 013106.
- 24 E. Durgun, S. Ciraci and T. Yildirim, Phys. Rev. B, 2008, 77, 085405.
- 25 Q. Sun, Q. Wang, P. Jena and Y. Kawazoe, J. Am. Chem. Soc., 2005, 127, 14582–14583.
- 26 T. Guo, C. Jin and R. E. Smalley, *J. Phys. Chem.*, 1991, **95**, 4948–4950.
 27 Y. Chai, T. Guo, C. Jin, R. E. Haufler, L. P. F. Chibante, J. Fure, L. Wang,
- J. M. Alford and R. E. Smalley, J. Phys. Chem., 1991, 95, 7564–7568.
 D. Golberg, Y. Bando, K. Kurashima and T. Sasaki, Appl. Phys. Lett.,
- 1998, 72, 2108.
 D. Golberg, Y. Bando, O. Stéphan, L. Bourgeois, K. Kurashima, T. Sasaki, T. Sato and C. Goringe, *J. Electron Microsc.*, 1999, 48, 701–709.
- 30 Y. J. Zou, X. W. Zhang, Y. L. Li, B. Wang, H. Yan, J. Z. Cui, L. M. Liu and D. A. Da, *Journal of Materials Science*, 2002, 37, 1043–1047.
- 31 H.-J. Muhr, R. Nesper, B. Schnyder and R. Kötz, *Chemical Physics Letters*, 1996, 249, 399–405.
- 32 B. Cao, X. Zhou, Z. Shi, Z. Gu, H. Xiao and J. Wang, Fullerene Science and Technology, 1998, 6, 639–648.
- 33 H. Lange, A. Huczko, P. Byszewski, E. Mizera and H. Shinohara, *Chemical Physics Letters*, 1998, 289, 174–180.
- 34 T. Nakamura, K. Ishikawa, K. Yamamoto, T. Ohana, S. Fujiwara and Y. Koga, *Phys. Chem. Chem. Phys.*, 1999, 1, 2631–2633.
- 35 J. C. Hummelen, C. Bellavia-Lund and F. Wudl, in *Fullerenes and Related Structures*, Springer Verlag Berlin Heidelberg, 1999, vol. 199, pp. 93–134.
- 36 A. A. Arie, B.-J. Jeon and J.-K. Lee, Carbon letters, 2010, 11, 127–130.
- 37 P. W. Dunk, A. Rodríguez-Fortea, N. K. Kaiser, H. Shinohara, J. M. Poblet and H. W. Kroto, *Angew. Chem. Int. Ed.*, 2012, **52**, 315–319.
- 38 S. K. Bhatia and A. L. Myers, *Langmuir*, 2006, **22**, 1688–1700.
- 39 E. Garrone, B. Bonelli and C. O. Areán, Chem. Phys. Lett, 2008, 456, 68-70.
- 40 C. O. Areán, S. Chavan, C. P. Cabello, E. Garrone and G. T. Palomino, *ChemPhysChem*, 2010, **11**, 3237–3242.
- 41 J. P. Perdew and Y. Wang, Phys. Rev. B, 1992, 45, 13244-13249.
- 42 G. Kresse and J. Hafner, *Phys. Rev. B*, 1993, **47**, 558–561.
- 43 G. Kresse and J. Furthmüller, Phys. Rev. B, 1996, 54, 11169–11186.
- 44 G. Kresse and J. Furthmüller, Comput. Mater. Sci., 1996, 6, 15-50.
- 45 P. E. Blöchl, Phys. Rev. B, 1994, 50, 17953-17979.
- 46 G. Kresse and D. Joubert, Phys. Rev. B, 1999, 59, 1758-1775.
- 47 Y. Y. Sun, K. Lee, L. Wang, Y.-H. Kim, W. Chen, Z. Chen and S. B. Zhang, *Phys. Rev. B*, 2010, **82**, 073401.
- 48 S. Grimme, J. Comput. Chem., 2006, 27, 1787–1799.
- 49 M. P. Waller, H. Kruse, C. Mück-Lichtenfeld and S. Grimme, *Chem. Soc. Rev.*, 2012, 41, 3119–3128.
 50 C. Li, L. Li, F. Wu, S. S. Li, L. R. Xia and L. W. Wang, *L. Phys. Chem. C.*
- 50 C. Li, J. Li, F. Wu, S.-S. Li, J.-B. Xia and L.-W. Wang, *J. Phys. Chem. C*, 2011, **115**, 23221–23225.
- 51 S. Lee, M. Lee, H. Choi, D. S. Yoo and Y.-C. Chung, Int. J. Hydrogen Energ., 2013, 38, 4611–4617.
- 52 Y. Wang, F. Wang, B. Xu, J. Zhang, Q. Sun and Y. Jia, J. Appl. Phys., 2013, 113, 064309.
- 53 Y. Wang, X. Li, F. Wang, B. Xu, J. Zhang, Q. Sun and Y. Jia, *Chem. Phys.*, 2013, **415**, 26–30.
- 54 S. Y. Willow and S. S. Xantheas, *Chem. Phys. Lett.*, 2012, **525–526**, 13– 18.
- 55 G. Henkelman, A. Arnaldsson and H. Jónsson, *Comput. Mater. Sci.*, 2006, 36, 354–360.
- 56 W. Tang, E. Sanville and E. Henkelman, *J. Phys.: Condens. Matter*, 2009, **21**, 084204.
- 57 R.-H. Xie, G. W. Bryant, J. Zhao, V. H. Smith Jr., A. D. Carlo and A. Pecchia, *Phys. Rev. Lett.*, 2003, **90**, 206602.
- 58 M. R. Manaa, Chem. Phys. Lett., 2003, 382, 194–197.

- 59 Y.-H. Kim, Y. Zhao, A. Williamson, M. J. Heben and S. B. Zhang, *Phys. Rev. Lett.*, 2006, **96**, 016102.
- 60 Q. Sun, Q. Wang and P. Jena, Appl. Phys. Lett., 2009, 94, 013111.
- 61 X.-B. Wang, H.-K. Woo and L.-S. Wang, J. Chem. Phys., 2005, 123, 051106.
- 62 D. Marx and J. Hutter, *Ab Initio Molecular Dynamics: Theory and Implementation*, John von Neumann Institute for Computing, Jülich, 2000, vol. 3, pp. 329–477.

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15x7mm (300 x 300 DPI)

Using first principles calculations we predicted new molecular based hydrogen storage systems, which compose of abundant elements, with interesting thermodynamics.