# Chemical Science

Accepted Manuscript



This is an *Accepted Manuscript*, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this Accepted Manuscript with the edited and formatted Advance Article as soon as it is available.

You can find more information about *Accepted Manuscripts* in the **Information for Authors**.

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard <u>Terms & Conditions</u> and the <u>Ethical guidelines</u> still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this *Accepted Manuscript* or any consequences arising from the use of any information it contains.



## Chemical Science

**RSCPublishing** 

### **EDGE ARTICLE**

Cite this: DOI: 10.1039/x0xx00000x

# A New pH Sensitive Fluorescent and White Light Emissive Material through Controlled Intermolecular Charge Transfer

Received 00th January 2012, Accepted 00th January 2012

DOI: 10.1039/x0xx00000x

www.rsc.org/chemicalscience

Y. I. Park,  $^{a, \pm}$  O. Postupna,  $^{b, \pm}$  A. Zhugayevych,  $^{b, \pm}$  H. Shin,  $^c$  Y. S. Park,  $^a$  B. Kim,  $^c$  H.-J. Yen,  $^a$  P. Cheruku,  $^a$  J. S. Martinez,  $^d$  J. W. Park,  $^c$  S. Tretiak  $^{b, *}$  and H.-L. Wang  $^{a, *}$ 

A new, pH dependent and water-soluble, conjugated oligomer (amino, trimethylammonium oligophenylene vinylene, ATAOPV) was synthesized with a quaternary ammonium salt and an aromatic amine at the two ends of a  $\pi$ -conjugated oligomer, thus creating a strong dipole across the molecule. A very unique white light LED is successfully fabricated from a stimuli responsive organic molecule whose emission properties are dominated by the pH value of the solution through controlled intermolecular charge transfer.

### Introduction

Conjugated materials have been extensively studied in recent years due to their tunable physical, optical and electronic properties and their applications in optoelectronics<sup>1</sup> and bio-sensor technologies.<sup>2</sup> Of particular interest are the p-phenylene vinylene (PPV) derivatives, one of the most extensively studied conjugated materials due to their high quantum yield and tunable structures that lead to desired physical and optical properties.<sup>3</sup> There is a long history of developing pH sensitive chromophores for biomedical research. In recent years, the development of pHdependent conjugated chromophores have shown promises in bioimaging of living cells<sup>4</sup> and sensing of proteins<sup>5</sup> and DNA.<sup>6</sup> Previously reported pH-dependent organic materials are xanthene, <sup>7</sup> cyanine, <sup>8</sup> fluorene, <sup>4</sup> acridine <sup>9</sup> and BODIPY <sup>4</sup> derivatives that show strong pH dependent fluorescence properties on the pH. Among all these chromophores, xanthene derivatives are of special interest as they exhibit white light emission under specific pH values; <sup>7</sup> Specifically, *Huvnh et al.* recently synthesized a dithienophosphole derivative that has white light emission from a single molecule species. 10 Despite this recent example, there has been limited success in developing a white emissive material resulting from a single molecule, 11 because the majority of molecular materials obey Kasha's rule exhibiting single-band fluorescence from the lowest singlet excited state. <sup>12</sup> In addition, even when light emission is generated due to a mix of multiple fluorescence peaks is generated, there is a tendency for energy cascading from high to low energy, which leads to a dominating emission color corresponding to the lowest energy. To overcome this low energy emission, several approaches have been attempted including those employing excited-state intramolecular proton transfer (ESIPT), 13 excimer, 14 and isomer in different pHs. 10

In this manuscript we report synthesis and characterization of a conjugated oligomer (ATAOPV) with pH-dependent fluorescence properties. We show that at pH 5.0, this oligomer exhibits dual fluorescence of blue and red light

leading to white light emission in aqueous solution, solid films and OLED devices. Our conjugated oligomer, ATAOPV, with a strong permanent molecular dipole moment, is unique as we can control (turn on/off) the excited state charge transfer via protonation/deprotonation. In addition to the white light emission, this ATAOPV can produce blue and orange emissions, tuned by pH. Our electronic structure modeling results suggest that the lower energy peak, formed at high pH values, is primarily attributed to intermolecular charge transfer. We have successfully fabricated LED devices that emit blue, orange or white light. Although the yield of the LED is relatively low, our study reveals a very unique white light LED fabricated from a stimuli responsive organic molecule whose emission properties are dominated by the pH value of the solution through controlled intermolecular charge transfer.

### Methods

Synthesis. The reaction scheme of ATAOPV is shown in Scheme 1. Compound A is synthesized by mono-reduction of the dinitro-OPV compound, 15 which is achieved using a procedure that involved chloromethylation of 1,4-dimethoxybenzene, converted into the phosphonate and then coupled with a vinyl group via Horner-Wadsworth-Emmons reactions.<sup>3</sup> The reduction of the nitro group into the amine was achieved by reacting dinitro-PPV with pyridine and sodium sulphate, under relatively mild reaction conditions. Further treatment of the product with excessive amount of methyl iodide leads to formation of the quaternary ammonium group. The remaining nitro group is then converted to the amine using a strong reducing agent, stannous chloride in acidic condition. The end product has a protonated amine salt which is then neutralized to the amine using triethylamine in methanol solution. The as-synthesized PPV oligomer is then purified via recrystallization resulting in a pure yellow solid. The <sup>1</sup>H and <sup>13</sup>C NMR spectra of ATAOPV are illustrated in Figure S1 and agree well with the proposed

**Journal Name** 

molecular structures. Detailed information on synthesis, characterization, and OLED fabrication is provided in Supporting

Scheme 1. Synthetic scheme for ATAOPV.

Computational methods. We performed quantum-chemical modeling of protonated and non-protonated molecules as well as their dimer structures. CAM-B3LYP density functional 16 combined with the 6-31G\* basis set was used as a default method for all molecular modeling, carried out using the Gaussian 09 program.<sup>17</sup> Density Functional Theory (DFT) and time-dependent DFT (TDDFT) frameworks<sup>18</sup> have been used for calculations of ground and excited state electronic properties, respectively. CAM-B3LYP is a long-range corrected hybrid functional providing a correct description of neutral electronic excitations and charged states of extended  $\pi$ -conjugated systems.<sup>19</sup> Because excitation energies are highly sensitive to the choice of density functional, performance of CAM-B3LYP was compared to that of B3LYP and ωB97X functionals, see Supporting Information. The 6-31G\* basis set is known to provide adequate results for geometry and energy values, for relatively large molecules consisting of the first and second row elements, 20 in a noticeably short CPU time compared to that of larger basis sets. Notably, inclusion of polarization orbitals is important for excited states and charged molecules. 18 The introduction of diffuse orbitals is known to enhance the accuracy of the acidity prediction, <sup>18</sup> however, for the given system the improvement is found to be insignificant.

Frontier orbitals (including HOMO and LUMO), natural orbitals (NO), and natural transition orbitals (NTO)<sup>21</sup> are used for visual analysis of the electronic wave function of the ground and excited states. Natural (transition) orbitals give the best representation of a one-electron (transition) density matrix in terms of molecular orbitals, irrespective of the quantum chemistry method by which that density is obtained. Consequently, a comparison of NO/NTO with the frontier orbitals allow for a robust identification of the nature of the excited/anionic/cationic states, provided that the occupation numbers of natural orbitals are close to integers (which is true for all the considered molecules/states). In particular, for an anion/cation we look at the half-occupied NO; the occupation numbers of the two adjacent NOs should be close to 2 and 0 in order for such NO analysis be meaningful. In the NTO analysis of an excitation transition we look at the highest occupied (hole NTO) and lowest unoccupied (electron NTO) orbitals, whose NTO weights should be close to 1 to identify such transition with a single excitation. In the NO analysis of an excited state corresponding to a single excitation we look for a pair of orbitals at the boundary between "doubly occupied" and "vacant" NOs. The difference in occupations of these two NOs characterizes the amount of charge transfer in the excited state, whereas the spatial separation of the corresponding hole and electron NTOs characterizes the charge separation itself.

To implicitly account for an influence of the media on the investigated properties the polarizable continuum model (PCM)

with the appropriate static and optical dielectric constants for water has been used for geometry optimization and calculation of optical spectra. Solvation of excited states is simulated using the state-specific (SS) approach.<sup>22</sup> An alternative, linear response (LR) method<sup>23</sup> leads to nearly the same result for the first absorption peak. Calculation of acid dissociation constant (see details in Supporting Information) is carried out using the SMD solvation model,<sup>24</sup> which also includes non-electrostatic solvation effects. This model is the recommended choice for computing Gibbs free energy of solvation, due to its fine-tuned parameters. Explicit solvation was used in classical molecular dynamics (MD) simulations with MM3 force field.<sup>25</sup> These calculations are carried out using the TINKER program.<sup>26</sup> The MM3 water box has been obtained by annealing the TIP3P water box. The simulations are conducted at 400K (at 300K the thermodynamic properties of MM3-water do not match experiment, as it is a high-density liquid).

In molecular calculations, counter-ions that could be possibly present in a solvent environment have been neglected, and the groups NMe<sub>3</sub> and NH<sub>3</sub> are assumed to be positively charged in a solution. This assumption is tested as follows. First, the only stable position of the iodide bound to the non-protonated molecule is in the NMe<sub>3</sub><sup>+</sup> "pocket". Then the kinetic stability of this configuration has been tested by MD in water: at 400 K the iodide dissociates from the molecule in tens of picoseconds. Finally, even if iodide is placed in the "pocket", its direct influence on absorption/emission spectra as well as on direct electron transfer is negligible, see Table S1.

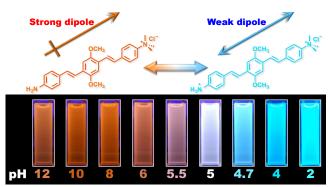
Ground state molecular geometries have been optimized in vacuum and water. In both cases the geometries are topologically equivalent to each other. The lowest excited state structures have been subsequently optimized starting from the ground state geometry in corresponding media. The  $\pi$ -conjugated system is planar in the excited state and non-planar in the ground state, with the terminal benzene rings slightly rotated around the long axis of a molecule, in agreement with previous studies.<sup>27</sup> Only the most energetically favorable molecular conformation is considered here, since it was shown previously<sup>3</sup> that the optical properties of various conformations differ insignificantly. Selected dimer configurations are obtained by  $\pi$ -stacking two molecules at 3.5 Å with zero slip-stack displacement in the orientation minimizing electrostatic and steric interactions. Dimer geometry in the ground state has been optimized by using  $\omega B97X$  functional, which accounts for dispersive interactions. Finally, statistical sampling of dimer geometries is performed by 10 ns MD in a water box using MM3 force field.

### Results and discussion

Observed pH-dependent optical properties. Figure 1 shows the chemical structures of non-protonated and protonated ATAOPV molecules (top panels) and their emission properties in aqueous solution at various pHs (bottom panel). The fluorescence color of ATAOPV spanning from red to blue colors is coupled to the decrease of aqueous solution pH from 12 to 2. Judging from the titration curve (Figure S2), we can reasonably assume that the fluorescence at pH 12 and pH 2 originate from the non-protonated and pronated molecular species, respectively. It is interesting to note that white light emission is attained when the pH value reaches 5.0.

**ARTICLE** 

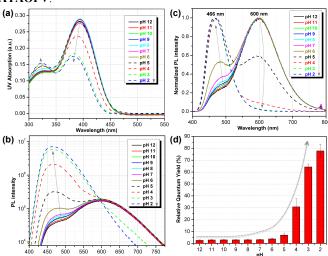
Page 3 of 7 Chemical Science



**Journal Name** 

**Figure 1.** (Top) Diagram of the protonation of **ATAOPV**: on the left is the non-protonated molecule whose chemical formula is abbreviated as  $NH_2$ -OPV- $NMe_3^+$ , on the right is the protonated molecule  $NH_3^+$ -OPV- $NMe_3^+$ . (Bottom) Photographs of **ATAOPV** multicolor emission as a function of pH 12-2 in solution ( $10^{-5}$  M in  $H_2O$ ).

The optical properties of **ATAOPV** in water (10<sup>-5</sup> M), as a function of pH, are summarized in Figure 2. The UV/Vis absorption of the aqueous solution of ATAOPV at pH 12 shows two distinct peaks at 324 nm and 394 nm (Figure 2a). Decreasing the pH leads to a blue shift of the absorption peak (from 394 nm to 384 nm) and a concomitant decrease of the peak intensity. The photoluminescence (PL) spectra of ATAOPV (Figure 2b and 2c) generally has two broad peaks at 466 nm and 600 nm, particularly pronounced at pH 5 resulting in white light emission. The absolute intensity of the 466 nm peak increases dramatically with decreasing pH (Figure 2c). The measured quantum yield (QY) decreases from 77.8% to 2.6% as the pH increases from 2 to 12 (Figure 2d). Such dramatic decrease of the quantum yield accompanied by the huge PL-redshift suggests a charge transfer (CT) character of the lowest excited state of the non-protonated ATAOPV. 13a, 28

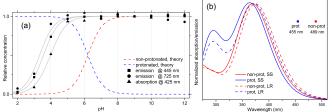


**Figure 2.** pH-dependent optical properties of **ATAOPV** ( $10^{-5}$  M in  $H_2O$ ): (a) UV-Vis absorption spectra; (b) Normalized PL spectra; (c) Absolute PL spectra in logarithmic scale; (d) Relative quantum yield using diphenylanthracene (DPA) in ethanol, as standard.

Statistical analysis of the pH-dependence of absorption and emission spectra (see SI page S9) confirms that these spectra can be represented as a linear superposition of the spectra of two species whose relative concentration is described by the formula

$$\frac{[\text{nonprotonated}]}{[\text{protonated}]} = 10^{pH-pKa}.$$

From this analysis the estimated  $pK_a$  is 3.3-3.6 from the PL spectra and 3.6-4.0 from the absorption spectra (**Figure S3**). These results are consistent with the titration curve (**Figure S2**), which gives the upper limit for the  $pK_a$  to be roughly 4.0. *Ab initio* estimate for the  $pK_a$  is 6.2, so that the calculated protonation energy of 0.36 eV is overestimated by 0.15 eV, probably due to the use of an implicit solvation model in our simulations. In any case, both experiment and theory suggest that protonated molecules are the dominant species at pH 2, whereas at pH 12 their concentration is negligible (**Figure 3a**).



**Figure 3.** (a) Calculated dependence of the concentration of protonated and non-protonated species on pH of the solution, compared to data extracted from experimental absorption spectrum at 425 nm, and emission spectrum at 445 and 725 nm. Gray curves are obtained by fitting Eq. 1 to these data. (b) Calculated absorption spectra and PL peak position for protonated and non-protonated molecules in water. In SS calculations only the first state was optimized. Spectral line shapes are obtained by Gaussian broadening of vertical electronic transitions.

Deeper insight can be obtained by analyzing the shape of the absorption and emission peaks. The results of such analysis (**Figure S4**) demonstrate that absorption spectra at pH 2 and pH 12 can be fitted by a Gaussian broadened vibronic progression, with the parameters typical for functionalized OPV molecules. <sup>19a</sup> The emission spectra are structureless, with less broadening than that for absorption. At pH 2, the blue PL band has a single component, whereas at pH 12 there are two gaussian bands: red and blue. Thus obtained vertical excitation energies and Stokes shifts are summarized in **Table 1**.

	Absorption, eV		Stokes shift, eV	
	exp.	theory	exp.	theory
reference	3.21	3.30	0.68	0.52
pH 2	3.26	3.32	0.68	0.54
pH 12	3.16	3.28	0.63	0.66
pH 12 red			1.17	1.16

**Table 1.** Experimental and calculated absorption energy and Stokes shift corresponding to vertical electronic transitions. The "reference" molecule is  $NMe_3^+$ -OPV- $NMe_3^+$  in water. At pH 12 two emission components are given: blue and red. For the first three rows the theoretical calculations are performed for a single molecule in water (SS solvation). In the last row emission energy is calculated as E(+1,+1)-E(0,+1)+E(-1,-1)-E(0,-1), where  $E(Q_1,Q_2)$  is the energy of the molecule with extra charge  $Q_1$  in the geometry optimized with charge  $Q_2$  (dimer approximated as infinitely separated cation and anion). An extended version of this table is given is Supporting Information (Table S1).

The time-resolved fluorescence decay lifetime measurements (Figure 4) made separately in blue (450±20 nm) and red (580±10 nm) spectral regions confirm that there are two emission components with distinct lifetimes (1.3 and 2.5 ns) and, consequently, without an efficient energy transfer between them. Assuming the Forster type energy transfer to be the only energy exchange channel between molecular species, experimental observations across the entire pH range are then consistent with the presence of two molecular species with negligible energy transfer. The latter is rationalized by (an) shiring spectral overlap between the absorption at pH 12 (394 nm, non-protonated species)

**Journal Name** 

and emission at pH 2 (466 nm, protonated species, see Figure **S5**). <sup>13a</sup>

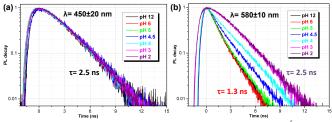


Figure 4. Fluorescence lifetime decay curves of ATAOPV (10<sup>-5</sup> M in  $H_2O$ ) as a function of pH in two spectral windows: (a)  $450 \pm 20$  nm and (b)  $580 \pm 10$  nm. The excitation wavelength is 405 nm.

Measured PL lifetimes  $(\tau)$  and QY allow for estimation of radiative and non-radiative times using the formulas

 $\tau_{\rm rad} = \tau/{\rm IQY}$ ,  $\tau_{\rm nrad} = \tau/(1-{\rm IQY})$ ,

where IQY is the internal QY. If IQY is close to QY then the calculated times are  $\tau_{rad}$ =2.5-3 ns,  $\tau_{nrad}$ ≥10 ns at 466 nm, and  $\tau_{rad}$ ≥ 50 ns,  $\tau_{nrad}$ =1.3 ns at 600 nm. The radiative lifetime at 466 nm corresponds to the oscillator strength ~1, which is typical for a single OPV molecule emission.<sup>3</sup> The radiative lifetime at 600 nm corresponds to the oscillator strength ~0.1, which may be a signature of weakly emitting CT state. It should be noted that the kinetics of PL decay at 600 nm for pH 12 is not single exponential: at large times the slope of PL decay approaches that for pH 2 (Figure 4b), which, by magnitude, is consistent with the tails of the blue PL band in Figure 2c.

To summarize, the experimental observations are fully consistent with the assumption that at pH 2 the optical properties are dominated by the protonated highly emissive ATAOPV molecules in a dilute aqueous solution, whereas at pH 12 two distinct non-protonated species are present: one with optical properties similar to the protonated molecule (very minor concentration) and another one (dominating concentration) having the CT lowest excited state. To elaborate this hypothesis and determine the nature of the CT state we have performed first principle modeling of molecules and their aggregates.<sup>29</sup>

Ab initio calculations of single-molecule absorption and emission spectra. Compared to substitutions with neutral donor/acceptor groups, the influence of NMe<sub>3</sub><sup>+</sup> and NH<sub>3</sub><sup>+</sup> charged groups on electronic properties of the functionalized OPV molecule has important differences. This can be rationalized by examining various orbital plots. In particular, strong acceptors such as NO<sub>2</sub> (see Figure S6) influences the host molecule mainly by hybridizing its low-energy  $\pi$ -system with that of the OPV molecule, whereas the effect of creating an extra electric field for the  $\pi$ -system is less important. <sup>19a</sup> In contrast, the charged groups create much stronger electric field, whereas the  $sp^3$  hybridization of the nitrogen atom pushes its frontier  $\pi$ -orbitals out of resonance with the  $\pi$ -system of the host molecule. That is why wavefunctions of the molecules functionalized by NMe<sub>3</sub><sup>+</sup> and NH<sub>3</sub><sup>+</sup> groups are essentially the same as that for the molecules with hydrogen in place of these groups, i.e.  $\pi$ -isoelectronic molecules OPV-NH<sub>2</sub> and OPV (see Figures S7-S10). At the same time, the energetic characteristics, which are sensitive to charge and electric field distribution, such as solvatochromic shift and CT state transition energy, are different (Figure S7). These considerations help to interpret first-principle calculations results presented below

Calculated single-molecule absorption spectra are in a good agreement with experiment (Figure 3b and Table 1). The main peak is well resolved and is composed of a single electronic transition. NTO analysis (Figure 5) clearly identifies it as HOMO to LUMO  $\pi$ - $\pi$ \* transition of the conjugated system with a large transition dipole directed along the molecule. The pH dependence of both the peak position and its intensity is qualitatively reproduced by calculations (quantitatively theory underestimates the changes) as shown in Figure 3a. The second absorption band is comprised of several competing electronic transitions. The more pronounced character of this band for the protonated molecules compared to the non-protonated species is well reproduced by calculations (compare Figure 2a and Figure 3b). All above features exactly replicate the corresponding trends for the  $\pi$ isoelectronic molecules OPV and OPV-NH2, respectively.

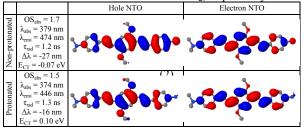


Figure 5. Calculated properties of the first excited state in water of protonated and non-protonated molecules. Natural transition orbitals are shown for the lowest excitation in the relaxed geometry. OS is oscillator strength,  $\tau_{\text{rad}}$  is radiative lifetime,  $\Delta\lambda = \lambda_{\text{emi}}(\text{water}) - \lambda_{\text{emi}}(\text{vacuum})$  is solvatochromic shift,  $E_{CT} = E_{cation} + E_{anion} - E_{exciton} - E_{ground}$  is the intermolecular CT energy relative to intramolecular exciton (relaxed geometries, equilibrium solvation).

Calculated single-molecule emission peak positions and fluorescence lifetimes (Table 1 and Figure 5) match the observed characteristics of the blue PL band (for both pH 2 and pH 12). Therefore, we can associate the blue PL band at pH 2 and pH 12 with the emission from protonated and non-protonated single molecules, respectively. Strictly speaking, the true origin of the blue PL band at pH 12 is unclear, because such weak signal can be also produced by contaminants such as residue protonated molecules (which should be absent at pH 12).

The red PL band is presumably attributed to a CT state because of a large fluorescence lifetime, low QY, and strong solvatochromism. Restricted and unrestricted TDDFT calculations show no intramolecular CT state for any of the tested geometries: ground, excited, or hypothetical forced CT state (in the latter, opposite sides of a molecule are optimized with opposite extra charges), as well as for the lowest triplet state. The measured PL spectrum in tetrahydrofuran (Table S1) also rules out the existence of an intramolecular CT state: the experiment shows positive solvatochromism, i.e. stabilization in a solvent, whereas the putative CT state should be destabilized in a more polar solvent because of the screening of the NMe<sub>3</sub><sup>+</sup> charge (for the same reason non-CT state also shows negative solvatochromism). An intramolecular non-CT state on a charged molecule (electron transfer from the counter-ion) is also excluded because such a molecule would emit at 800 nm. Therefore, experimentally observed red PL band has been attributed to an intermolecular CT

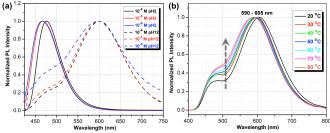
**Origin of the red PL band.** There are at least three possible donor-acceptor complexes able to host an intermolecular CT state: dimer of two molecules, molecule + water, molecule + counterion. A rigorous theoretical study of each of these possibilities requires an explicit solvation of CT states as well as proper consideration of electronic processes in water, which themselves are challenging theoretical problems.<sup>30</sup> For this reason we limit our investigation to analysis of only structural and energetic

Page 5 of 7 Chemical Science

Journal Name ARTICLE

feasibility of a CT state. First, the prerequisite for an efficient CT formation is an availability of the ground state configuration of the donor-acceptor complex. Second, since the Coulomb interaction stabilizes the CT state, the energy of two infinitely separated charges gives a robust upper estimate of the lowest CT state energy.

We start with considering an intermolecular CT on the dimer, which is known to produce a large solvatochromic shift e.g. for merocyanine dyes.<sup>31</sup> Important parameters here are ionization potential (IP) and electron affinity (EA) as well as exciton energy for non-protonated molecule in water. These 3 quantities calculated for respective vibrationally relaxed species are 4.8 eV, 2.2 eV, and 2.7 eV, respectively. We immediately see that CT state on a dimer is lower in energy than the intramolecular exciton by 0.1 eV (see Tables 1 and S1). The binding energy of a dimer in its ground state is about 0.5 eV as obtained in both ab initio ωB97X/6-31G\* and MM3 force field calculations. Classical MD simulations in water at 400 K show that the dimer is stable for at least 10 ns, though its geometry fluctuates substantially. At large charge separation the estimated emission peak position exactly fits the observed red PL band (Table 1). Calculations for selected bound dimer configurations show essentially the same red-shifted emission (Table S1) as well as give very weak oscillator strengths consistent with the observed radiative lifetimes. At the same time, the optical absorption spectrum of a dimer should not deviate from the single-molecule absorption because of large geometry fluctuations suppressing delocalized intermolecular excitations. The dimer picture is also consistent with the observed decrease of PL intensity with dilution relatively to the blue emission band (see Figure 6a and S11). For protonated species the intermolecular CT state is energetically unfavorable, as well as there are no stable dimer configurations in the ground state.



**Figure 6.** (a) Experimental PL spectra normalized in such a way that the maximum PL intensity at pH 2 equals 1, whereas the spectra at pH 12 are additionally scaled up by a fixed constant. (b) Temperature-dependent PL spectra in a 10<sup>-5</sup> M ATAOPV aqueous solution reveal increase in PL intensity of blue emission with increasing temperature.

For charge transfer to solvent scenario there are two options. The variant with the non-protonated molecule being the acceptor is energetically unfavorable because the IP of liquid water is large, about 6.5 eV. $^{32}$  That IP is not exactly the energy of the solvated cation ( $\rm H_2O^+$  or  $\rm H_3O^+$ ) but corresponds to the light-driven electron detachment. The EA of water is 3-4 eV $^{33}$  making the electron transfer to water energetically feasible in the form of bulk or surface solvated electron. The solvated electron has short lifetime, hundreds of picoseconds,  $^{33b}$  and allows for plenty of nonradiative recombination channels,  $^{30a}$  that is consistent with low QY of red PL. Because the dispersion of energies of the solvated electron, 0.5 eV,  $^{33b}$ ,  $^{33c}$  is much larger than the observed emission bandwidth, 0.2 eV, this kind of charge transfer determines mainly nonradiative channel. In contrast to nonprotonated molecule, the protonated one has 0.5 eV higher IP and very different solvation shell that may suppress the charge escape to the solvent.

Finally, the charge transfer to iodide can be excluded: its IP/EA is too high/low. Also we did not find any electronically coupled configurations of molecule + iodide complex similar to that observed for other systems.<sup>34</sup> In summary, the intermolecular CT state is structurally and energetically feasible in two scenarios: dimer and charge transfer to solvent, with the former being the primary radiative channel and the latter being the mainly nonradiative channel.

To further test the dimer model, the temperature dependence of PL spectra was measured (Figure 6b). The observed increase of the relative intensity of the blue emission component, with temperature increases, is consistent with the increased thermal dissociation of dimers. At the same time, this trend is inconsistent with intramolecular CT hypothesis: the relaxation of initial high-energy excitation to a putative low-lying intramolecular CT state should be faster at higher temperatures, yielding the opposite temperature behavior of the blue-to-red PL intensity ratio. Also, the observed negative thermochromism of the red emission peak has no simple explanation in an intramolecular CT model, whereas in the dimer model it is a consequence of the increased intermolecular separation leading to an upshift of the lowest excitation energy.

White emission and electroluminescence. One very interesting property of ATAOPV in pH 5 aqueous solution is the white light emission resulting from the combination of two emission colors at 466 and 600 nm. The pH-dependent optical properties in solution are reflected also in films and organic light emitting diodes (OLED) fabricated by spin-coating from ATAOPV solution. As shown in Figure 7a, the films from different pH solutions are shown as blue, white, and orange under UV excitation. In OLED devices, as-synthesized ATAOPV was used as an emitting material layer (EML). The EMLs were spincoated from 1 wt % of ATAOPV solutions with different ratios of formic acid (acid): ammonia (base) (4:0, 4:1, 4:2, 4:4) and the spin-coated thin films were dried overnight in a 110 °C vacuum oven to ensure complete evaporation of formic acid and ammonia. The other layers were vacuum deposited with device architectures that include a hole inject/transport layer, 4,4',4"-tris-(N-(naphthylen-2-yl)-*N*-phenylamine)triphenylamine N,N'-di-[(1-naphthyl)-N,N'-diphenyl]-1,1'-biphenyl)-4,4'-diamine (NPB), and a hole blocking/electron transporting layer, 1,3,5tris(N-phenylbenzimidizol-2-yl)benzene (TPBI). The complete multilayered device architecture is indium tin oxide (ITO)/2-TNATA/NPB/ATAOPV/TPBI/LiF/A1 (Figure Electroluminescence (EL) spectra (colors) of LED devices are determined by the formic acid and ammonia (base) ratios (Figure 7b and 7c). The shift in the EL spectra corresponds to the change in the CIE chromaticity diagram coordinates from (0.18, 0.27) to (0.46, 0.48), see Figure 7d. Overall, this trend in the EL spectra is in a good agreement with the shifts from blue to the red emission region observed in the PL spectra of the ATAOPV in solution, as a function of pH. Of particular interest is the white emission EL device as CIE (0.28, 0.34) with a color rendering index (CRI) of 81.9 and the colors of the EL emission match very nicely to that of the thin film under UV excitation (see Figure 7a and 7b). The turn-on voltage (V) and external quantum efficiency, EQE (%) for the EL devices are 3.0 V, 0.08%; 2.25 V, 0.001%; 4.5 V, 0.007% for blue, white, orange EL devices, respectively. The efficiency of these LED devices is low mainly because of the ionic nature of the emitter. The ionic species lead to shallow or deep traps inside HOMO and LUMO levels that quench excitons.<sup>35</sup> Although the EQE for the white light LED is relatively low, the research to optimize the device performances is

Shemical Science Accepted Manusc

ARTICLE Journal Name

ongoing. Altogether, these results from solution, film and device are experimental demonstration of multi-color emission though a pH-controlled interplay of intra- and inter-molecular fluorescent molecular species.

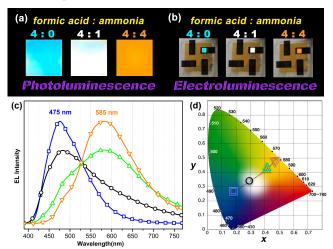


Figure 7. (a) Photographs of thin film emission depend on pH under UV lamp and (b) the image of LED devices under operation. (c) EL spectra and (d) CIE chromaticity diagram depend on formic acid to ammonia ratio ( $\square$ ) 4:0, ( $\bigcirc$ ) 4:1, ( $\triangle$ ) 4:2, and ( $\nabla$ ) 4:4.

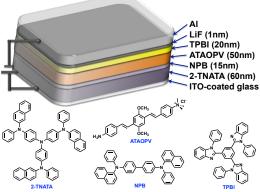


Figure 8. Schematic illustration of OLED structure.

### **Conclusions**

We have designed and synthesized a water-soluble oligomer ATAOPV with amine (electron donor) and ammonium salt (electron acceptor) on both ends of the molecule. This molecule shows dramatic change in fluorescence color from red to blue (bottom panel in Figure 1) and a 30 fold increase in quantum yield from 2.6 to 77.8 % as the pH decreases from 12 to 2 couple with a gradual change from the non-protonated to pronated molecular species (top panel in Figure 1). Both protonated and nonprotonated species have absorption spectra over the broad range from 300 to 450 nm. Our TDDFT quantum chemical simulations attribute these absorbance to  $\pi$ - $\pi$ \* intramolecular transitions of the conjugated system. Likewise, the blue band appearing in the PL spectra is assigned to intramolecular emissions of protonated and non-protonated molecules, as suggested by theoretical modeling and fits of time-resolved emission dynamics. In contrast, the red PL band appearing with increasing concentration of nonprotonated molecules is suggestive of appearance of weakly emissive intermolecular CT transition. In particular, this red PL dominates the fluorescence at pH 12. Our simulations ascribe this red PL to CT state in an antiparallel dimer of non-protonated molecules (determining mainly radiative decay channel) and

electron transfer to solvent (determining mainly non-radiative decay channel). Such emissive properties of ATAOPV are unique and different compared to the intramolecular CT transitions observed in many previously reported conjugated molecules with strong donor and acceptor substituents.<sup>36</sup> At pH 5.0, ATAOPV in aqueous solution exhibits dual color emission originating from both blue and red broad fluorescence bands. Unique emissive properties of ATAOPV have been further used to fabricate thinfilm LEDs displaying blue, white and orange colors as we vary the processing conditions (changing the ratio between formic acid and ammonium hydroxide). Specifically, the LED spectra variations correspond to the CIE coordinate change from (0.18, 0.27; blue) to (0.28, 0.34; white) to (0.46, 0.48; orange). Thus our joint experimental and theoretical study suggests pH controlled ratio of protonated and non-protonated species leading to tunable fluorescence across the entire visible region. Our experiments demonstrate multi-color emission resulting from a single molecular species in solution, thin film and LED device, which may have implication for light-emitting technologies based on organic semiconductor materials.

### Acknowledgements

The authors would like to acknowledge financial support by the Laboratory Directed Research and Development program, (YIP, HLW, ST, OP and AZ), for the synthesis and theoretical consideration of conjugated oligomers. Photophysical characterization of conjugated oligomers was supported by the Basic Energy Science, Biomolecular Materials Program, Division of Materials Science and Engineering (H.-L Wang and J.S. Martinez). A.Z. thanks Troy van Voorhis for discussion on charge transfer state and Jason Keith for helpful suggestions on computational methodology. This work was performed, in part, at the Center for Integrated Nanotechnologies, an Office of Science User Facility operated for the U.S. Department of Energy (DOE) Office of Science. Los Alamos National Laboratory, an affirmative action equal opportunity employer, is operated by Los Alamos National Security, LLC, for the National Nuclear Security Administration of the U.S. Department of Energy under contract DE-AC52-06NA25396.

### **Abbreviations**

CT, charge transfer; HOMO, highest occupied MO; LR, linear response; LUMO, lowest unoccupied MO; MD, molecular dynamics; MO, molecular orbital; NO, natural orbital; OLED, organic light-emitting diode; PL, photoluminescence; QY, quantum yield; SS, state specific; TDDFT, time-dependent density functional theory.

### **Notes and references**

- <sup>a</sup> Physical Chemistry and Applied Spectroscopy (C-PCS), Chemistry Division, Los Alamos National Laboratory, Los Alamos, New Mexico 87545. United States
- <sup>b</sup> Theoretical Division, Los Alamos National Laboratory, Los Alamos, New Mexico 87545, United States.
- <sup>c</sup> Department of Chemistry/Display Research Center, Catholic University of Korea, Bucheon 420-743, Republic of Korea.
- <sup>d</sup> Center for Integrated Nanotechnologies, Materials Physics and Applications Division, Los Alamos National Laboratory, Los Alamos, New Mexico 87545, United States.

**Chemical Science Journal Name ARTICLE** 

¥ Authors with equal contribution.

Electronic Supplementary Information (ESI) available: Text: experiments, characterization, devices fabrication. Figure: NMR, titration curve, estimation of pKa, single quantum mode fit, absorption/PL spectrum in pH 2/12, different orbitals for molecule A, half-occupied natural orbitals, PL spectra of ATAOPV at different concentrations under different pH, experimental and calculated absorption and emission. See DOI: 10.1039/b000000x/

- 1. (a) C. W. Tang and S. A. Vanslyke, Applied Physics Letters, 1987, 51, 913-915; (b) Y.-I. Park, J.-H. Son, J.-S. Kang, S.-K. Kim, J.-H. Lee and J.-W. Park, Chemical Communications, 2008, 2143-2145; (c) Y.-I. Park, J. S. Lee, B. J. Kim, B. Kim, J. Lee, D. H. Kim, S.-Y. Oh, J. H. Cho and J.-W. Park, Chemistry of Materials, 2011, 23, 4038-4044
- 2. L. H. Chen, D. W. McBranch, H. L. Wang, R. Helgeson, F. Wudl and D. G. Whitten, Proceedings of the National Academy of Sciences of the United States of America, 1999, 96, 12287-12292.
- 3. Y. I. Park, C.-Y. Kuo, J. S. Martinez, Y.-S. Park, O. Postupna, A. Zhugayevych, S. Kim, J. Park, S. Tretiak and H.-L. Wang, ACS Applied Materials & Interfaces, 2013, 5, 4685-4695.
- X.-X. Zhang, Z. Wang, X. Yue, Y. Ma, D. O. Kiesewetter and X. Chen, Molecular Pharmaceutics, 2013, 10, 1910-1917.
- G. Miesenbock, D. A. De Angelis and J. E. Rothman, Nature, 1998, **394**, 192-195.
- S. Ikeda, T. Kubota, M. Yuki, H. Yanagisawa, S. Tsuruma and A. Okamoto, Organic & Biomolecular Chemistry, 2010, 8, 546-551.
- Y. Yang, M. Lowry, C. M. Schowalter, S. O. Fakayode, J. O. Escobedo, X. Xu, H. Zhang, T. J. Jensen, F. R. Fronczek, I. M. Warner and R. M. Strongin, Journal of the American Chemical Society, 2006, 128, 14081-14092.
- M. S. Briggs, D. D. Burns, M. E. Cooper and S. J. Gregory, Chemical Communications, 2000, 2323-2324
- 9. M. G. Palmgren, Analytical Biochemistry, 1991, 192, 316-321.
- 10. H. V. Huynh, X. He and T. Baumgartner, Chemical Communications, 2013, 49, 4899-4901.
- 11. (a) J. Y. Li, D. Liu, C. W. Ma, O. Lengyel, C. S. Lee, C. H. Tung and S. T. Lee, Advanced materials, 2004, 16, 1538-1541; (b) M. C. Gather, A. Kohnen and K. Meerholz, Advanced materials, 2011, 23, 233-248; (c) G. Li, T. Fleetham and J. Li, Advanced materials, 2014, 26, 2931-2936; (d) X. Wu, L. Wang, Y. Hua, C. Wang, A. S. Batsanov and M. R. Bryce, *Tetrahedron*, 2014, 70, 2015-2019; (e) Q.-Y. Yang and J.-M. Lehn, Angewandte Chemie, 2014, 126, 4660-4665.
- 12. M. Kasha, Discussions of the Faraday Society, 1950, 9, 14-19.
- 13. (a) S. Kim, J. Seo, H. K. Jung, J. J. Kim and S. Y. Park, Advanced materials, 2005, 17, 2077-+; (b) S. Park, J. E. Kwon, S. H. Kim, J. Seo, K. Chung, S.-Y. Park, D.-J. Jang, B. M. Medina, J. Gierschner and S. Y. Park, Journal of the American Chemical Society, 2009, 131, 14043-14049; (c) J. E. Kwon and S. Y. Park, Advanced materials, 2011, 23, 3615-3642; (d) K. C. Tang, M. J. Chang, T. Y. Lin, H. A. Pan, T. C. Fang, K. Y. Chen, W. Y. Hung, Y. H. Hsu and P. T. Chou, Journal of the American Chemical Society, 2011, 133, 17738-17745.
- 14. (a) V. Adamovich, J. Brooks, A. Tamayo, A. M. Alexander, P. I. Djurovich, B. W. D'Andrade, C. Adachi, S. R. Forrest and M. E. Thompson, New Journal of Chemistry, 2002, 26, 1171-1178; (b) Y. Liu, M. Nishiura, Y. Wang and Z. M. Hou, Journal of the American Chemical Society, 2006, 128, 5592-5593.
- 15. U. Caruso, M. Casalboni, A. Fort, M. Fusco, B. Panunzi, A. Quatela, A. Roviello and F. Sarcinelli, Optical Materials, 2005, 27. 1800-1810.
- 16. T. Yanai, D. P. Tew and N. C. Handy, Chemical Physics Letters, 2004, 393, 51-57.
- 17. M. J. Frisch and etal, Gaussian 09, revision A.1, Gaussian, Inc., Wallingford CT, 2009.
- 18. C. Cramer, Essentials of computational chemistry, Wiley, 2004.
- 19. (a) A. Zhugayevych, O. Postupna, W. H.-L. and S. Tretiak, (submitted); (b) R. J. Magyar and S. Tretiak, Journal of Chemical

- Theory and Computation, 2007, 3, 976-987; (c) I. H. Nayyar, E. R. Batista, S. Tretiak, A. Saxena, D. L. Smith and R. L. Martin, The Journal of Physical Chemistry Letters, 2011, 2, 566-571; (d) A. Zhugayevych, O. Postupna, R. C. Bakus Ii, G. C. Welch, G. C. Bazan and S. Tretiak, The Journal of Physical Chemistry C, 2013, 117. 4920-4930.
- 20. (a) NIST Computational Chemistry Comparison and Benchmark Database, http://cccbdb.nist.gov; (b) E. Badaeva, M. R. Harpham, R. Guda, O. z. n. Süzer, C.-Q. Ma, P. Bäuerle, T. Goodson and S. Tretiak, The Journal of Physical Chemistry B, 2010, 114, 15808-15817; (c) A. Masunov and S. Tretiak, The Journal of Physical Chemistry B. 2003. 108, 899-907.
- 21. R. L. Martin, The Journal of Chemical Physics, 2003, 118, 4775-4777.
- 22. M. Cossi and V. Barone, The Journal of Chemical Physics, 2000, 112, 2427-2435.
- 23. M. Cossi and V. Barone, The Journal of Chemical Physics, 2001, 115, 4708-4717.
- 24. (a) C. J. Cramer and D. G. Truhlar, Chemical Reviews, 1999, 99, 2161-2200; (b) A. V. Marenich, C. J. Cramer and D. G. Truhlar, *The Journal of Physical Chemistry B*, 2009, **113**, 6378-6396.
- 25. (a) N. L. Allinger, F. Li, L. Yan and J. C. Tai, Journal of Computational Chemistry, 1990, 11, 868-895; (b) J.-H. Lii and N. L. Allinger, Journal of Computational Chemistry, 1998, 19, 1001-1016.
- 26. J. W. Ponder, TINKER 6.2, http://dasher.wustl.edu/tinker.
- 27. C. Katan, F. Terenziani, O. Mongin, M. H. V. Werts, L. Porrès, T. Pons, J. Mertz, S. Tretiak and M. Blanchard-Desce, The Journal of Physical Chemistry A, 2005, 109, 3024-3037.
- 28. (a) R. G. Bennett, Journal of Chemical Physics, 1964, 41, 3037-&; (b) A. P. Kulkarni, P. T. Wu, T. W. Kwon and S. A. Jenekhe, Journal of Physical Chemistry B, 2005, 109, 19584-19594.
- 29. (a) J. H. Hsu, W. S. Fann, P. H. Tsao, K. R. Chuang and S. A. Chen, Journal of Physical Chemistry A, 1999, **103**, 2375-2380; (b) S. R. Amrutha and M. Jayakannan, Journal of Physical Chemistry B, 2008, 112, 1119-1129.
- 30. (a) B. C. Garrett, D. A. Dixon, D. M. Camaioni, D. M. Chipman, M. A. Johnson, C. D. Jonah, G. A. Kimmel, J. H. Miller, T. N. Rescigno, P. J. Rossky, S. S. Xantheas, S. D. Colson, A. H. Laufer, D. Ray, P. F. Barbara, D. M. Bartels, K. H. Becker, K. H. Bowen, S. E. Bradforth, I. Carmichael, J. V. Coe, L. R. Corrales, J. P. Cowin, M. Dupuis, K. B. Eisenthal, J. A. Franz, M. S. Gutowski, K. D. Jordan, B. D. Kay, J. A. LaVerne, S. V. Lymar, T. E. Madey, C. W. McCurdy, D. Meisel, S. Mukamel, A. R. Nilsson, T. M. Orlando, N. G. Petrik, S. M. Pimblott, J. R. Rustad, G. K. Schenter, S. J. Singer, A. Tokmakoff, L.-S. Wang and T. S. Zwier, Chemical Reviews, 2004, 105, 355-390; (b) C. M. Isborn, B. D. Mar, B. F. E. Curchod, I. Tavernelli and T. J. Martínez, The Journal of Physical Chemistry B, 2013, 117, 12189-12201.
- 31. L. Lu, R. J. Lachicotte, T. L. Penner, J. Perlstein and D. G. Whitten, Journal of the American Chemical Society, 1999, 121, 8146-8156.
- 32. (a) A. Bernas and D. Grand, The Journal of Physical Chemistry, 1994, 98, 3440-3443; (b) D. M. Bartels and R. A. Crowell, The Journal of Physical Chemistry A, 2000, 104, 3349-3355.
- 33. (a) Á. Madarász, P. J. Rossky and L. Turi, The Journal of Chemical Physics, 2009, 130, -; (b) K. R. Siefermann, Y. Liu, E. Lugovoy, O. Link, M. Faubel, U. Buck, B. Winter and B. Abel, Nat Chem, 2010, 2, 274-279; (c) A. T. Shreve, T. A. Yen and D. M. Neumark, Chemical Physics Letters, 2010, 493, 216-219.
- 34. C. Reichardt and T. Welton, Solvents and solvent effects in organic chemistry, Wiley, 2011.
- 35. (a) V. I. Arkhipov, E. V. Emelianova and H. Bässler, Physical Review B, 2004, 70, 205205; (b) A. L. Burin and M. A. Ratner, The Journal of Physical Chemistry A, 2000, 104, 4704-4710; (c) C. Rothe, C.-J. Chiang, V. Jankus, K. Abdullah, X. Zeng, R. Jitchati, A. S. Batsanov, M. R. Bryce and A. P. Monkman, Advanced Functional Materials, 2009, 19, 2038-2044.
- 36. (a) S. A. Jenekhe, L. D. Lu and M. M. Alam, Macromolecules, 2001, 34, 7315-7324; (b) Z. R. Grabowski, K. Rotkiewicz and W. Rettig, Chemical Reviews, 2003, 103, 3899-4031; (c) S. J. Vella, J. Tiburcio, J. W. Gauld and S. J. Loeb, Organic Letters, 2006, 8, 3421-3424.