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Layer-modulated synthesis of uniform tungsten disulfide nanosheet using gas-phase precursors

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The synthesis of layered transition-metal-disulfide (MS₂, M = Mo, W) nanosheets with layer controllability and large-area uniformity is an essential requirement for their application in electronic and optical devices. In this report, we describe a synthesis process of WS₂ nanosheets with layer controllability and high uniformity using chemical vapor deposition (CVD) and WCl₆ and H₂S as gas-phase precursors. Through this process, we can systematically modulate the thickness of WS₂ nanosheets by controlling the duration of the reaction between WCl₆ and H₂S. The CVD-grown WS₂ nanosheets exhibit good stoichiometry as well as dependencies of a clear Raman shift and bandgap on the number of layers. These properties are confirmed by X-ray photoemission spectroscopy, Raman spectroscopy, and photoluminescence measurements. The number of layers of WS₂ nanosheets is confirmed by atomic force microscopy. Finally, we demonstrate the fabrication and performance of a photodetector based on a hybrid structure consisting of graphene and a WS₂ nanosheet.

Introduction

Recently, two-dimensional (2D) transition-metal disulfides (TMDs) have received much attention as new materials in logic circuits, thin-film transistors, sensors, catalysts, and optoelectronics. Furthermore, the combination of 2D materials (heterostructures) has attracted scientific as well as technical interest in emerging applications. Nano flakes of transition-metal disulfide (MS₂, M = Mo, W) can be easily exfoliated from bulk MS₂ and transferred onto an arbitrary substrate using the scotch-tape method; this is feasible because MS₂ crystals are composed of sulfur–transition-metal–sulfur (S–M–S) layers that are stacked on top of each other by van der Waals interaction. Using this method, the physical, electrical, and optical properties of MS₂ have been intensively investigated and found to depend significantly on thickness: while monolayered MS₂ has a direct bandgap, multilayered or bulk MS₂ has an indirect bandgap because of the broken symmetry in the atomic basis. Moreover, the bandgaps of MS₂ nanosheets vary with thickness, resulting in thickness-dependent photodetection in a wide range of spectral responses. Therefore, there is a critical need to synthesize thickness-modulated and uniform MS₂ nanosheets. Because mechanical exfoliation only produces small-sized flakes whose thickness cannot be controlled, considerable efforts have been devoted to synthesize MS₂-based materials, including sulfurization of metal or metal-oxide films and annealing with Mo- and S-containing solutions. However, these methods produce films with low carrier mobility (0.004–0.04 cm²V⁻¹s⁻¹) and they involve large complexity in the synthesis processes. More recently, chemical vapor deposition (CVD) methods using WO₃ and MoO₃ with sulfur powder have been demonstrated. Nanosheets of MoS₂ and WS₂ produced via this method show highly crystallinity and reasonable electrical properties; however, this synthesis process still suffers from a lack of reliability since it is dominated by many process parameters that cannot be easily controlled, such as sulfur concentration, distance between the substrate and the crucible containing the S powder, and outgoing flow of gases from the chamber resulting in discontinuous and non-uniform nanosheets. Consequently, an alternative approach is required to synthesize thickness-modulated and uniform nanosheets for future applications. We recently reported the synthesis of thickness-modulated WS₂ nanosheets using sulfurization methods with plasma-enhanced atomic layer deposition (PE-ALD) of WO₃ films, but successful use of direct deposition methods is not reported.

In this study, we developed a synthesis process with good thickness controllability and wafer-scale uniformity, through which single-, bi-, and multilayered WS₂ were systematically synthesized using the gas-phase CVD method. In particular, to more easily control the amount of precursors utilized, we used gas-phase precursors of WCl₆ and H₂S rather than WO₃ and S powders. The gas phase precursors only reported about the nanotube and nanoparticle not nanosheet. With this strategy, we systematically modulated the thickness of WS₂ nanosheets...
by controlling the reaction time between WCl₆ and H₂S. We analyzed WS₂ nanosheets of different thicknesses using optical microscopy (OM), Raman spectroscopy, high-resolution transmission electron microscopy (HRTEM), atomic force microscopy (AFM), and photoluminescence (PL) measurement. Furthermore, we demonstrated the fabrication and performance of a graphene–WS₂ hybrid photodetector.

Experimental

WS₂ Nanosheet Growth

For the synthesis of WS₂ nanosheets, SiO₂-coated Si wafers (referred to as SiO₂/Si substrates hereafter) were placed in the center of a tube furnace (3.05 cm/1.2 inch diameter). The temperature of the bubbler containing WCl₆ was kept at 100 °C to maintain an adequate vapor pressure, and gas-phase WCl₆ was transferred into the chamber by pure Ar (99.999 %) as the carrier gas.

Initially, the samples were heated to 700 °C for 30 min under flowing Ar gas (20 sccm) to remove organic contaminants on the surface. WCl₆ with a bubbling gas (Ar at 20 sccm) and H₂S were injected into the furnace at the same time, and the pressure of the chamber was maintained at 0.27 Torr during the process. At the end of the synthesis, the furnace was quickly cooled to room temperature under flowing Ar gas (20 sccm).

Transfer of WS₂ Nanosheets

The WS₂ nanosheets were covered with poly methyl methacrylate (PMMA) by spin coating at 2000 rpm for 30 s. After baking at 100 °C for 10 min, the samples were immersed in hydrofluoric acid for 1 s and then immersed in deionized water. Finally, PMMA-coated WS₂ nanosheets on SiO₂/Si substrates were removed and allowed to dry under ambient air. After that, PMMA was washed away from the samples with acetone. To determine the crystallinity of the nanosheets, TEM measurements were obtained with a Tecnai G2 F20 S-TWIN at an acceleration voltage of 200 kV.

Graphene–WS₂ Hybrid Photodetector

Graphene was synthesized by CVD on a copper foil in a quartz tube furnace using a H₂–CH₄ gas mixture. PMMA was then spin-coated onto the graphene/Cu substrate. The PMMA/graphene/Cu sample was placed in a container so that it floated on ammonium persulfate for 12 h in order to etch away the Cu foil. The resulting PMMA/graphene sample was then removed from the etch solution and was placed onto a CVD-grown monolayered WS₂ on a SiO₂/p++ Si substrate. The PMMA was removed with acetone and sample was then annealed at 300 °C for 3 h. We then fabricated source (S) and drain (D) electrodes using photolithography and lift-off processes, described as follows. First, photoresist (PR) layers were spin-coated on the sample. The substrate was then exposed to ultraviolet (UV) light with a photomask to form the S/D electrode pattern. After the pattern was developed, we first deposited Au/Cr for the S/D electrodes with a thermal evaporation system. The lift-off process was then completed with an acetone rinse. In addition, we patterned graphene using photolithography and oxygen-plasma etching.

Results and Discussion

Fig. 1 (a–c) SEM images showing the growth process of a WS₂ nanosheet from small triangles to a continuous film. (d) AFM image showing the second monolayer growing on the first monolayer.

The initial growth of WS₂ nanosheets was observed as a function of time by a scanning electron microscope (SEM). In Fig. 1 a–c, the changes in the growth process are demonstrated by showing the most commonly observed features. Initially, small triangular domains about 100 nm in size nucleated at random locations on the substrate surface. The nucleation sites then continued to grow and formed boundaries when two or more domains met, resulting in a partially continuous film. When the growth time was sufficiently long, this process eventually produced a WS₂ nanosheet with a large area, as shown in Fig. 1c. Fig. 1d. shows the nucleation of the second layer on the first monolayer, as observed by an AFM, which suggests a layer-by-layer growth mode of the WS₂ nanosheet.
OM images of WS$_2$ nanosheets transferred onto SiO$_2$ substrates for measurements of the (a) monolayer, (b) bilayer, and (c) tetralayer thickness. (d–f) Corresponding AFM images and inset height profiles of the transferred WS$_2$ nanosheets shown in (a–c).

Fig. 2a–c. show OM images of the synthesized WS$_2$ nanosheets after growth periods of 10, 15, and 25 min. The synthesized WS$_2$ nanosheets were uniform and continuous; the apparent colors of the transferred WS$_2$ nanosheets changed from pale violet to dark violet with increasing growth time. To confirm the thickness with AFM measurements, the WS$_2$ nanosheets were transferred to new SiO$_2$ substrates. Fig. 2d–f. show the AFM images of the transferred WS$_2$ nanosheets. The measured thicknesses of the synthesized WS$_2$ nanosheets were approximately 1 nm, 1.7 nm, and 3 nm for CVD durations of 10, 15, and 25 min, respectively. These thicknesses correspond to mono-, bi-, and tetralayers of WS$_2$ nanosheets, considering that the height of the WS$_2$ monolayer on SiO$_2$ was approximately 1 nm and the spacing between the first and the second WS$_2$ monolayers was approximately 0.65 nm. The larger AFM-derived spacing between the first monolayer and the substrate is not surprising, since it involved distinct tip–sample and sample–substrate interactions. Similar effects have been observed in MoS$_2$ on graphene. Thus, we conclude that the current WS$_2$ CVD process can systematically control the number of layers by controlling the growth time.

Fig. 3 (a) Raman spectra of the mono-, bi-, and tetralayered WS$_2$ nanosheets on SiO$_2$ substrates. (b) Relative Raman peak intensities and peak distances of the E$_{1g}^{1}$ and A$_{1g}^{1}$ modes of the mono-, bi-, and tetralayered WS$_2$ nanosheets.

The WS$_2$ nanosheets were further characterized using Raman spectroscopy, X-ray photoemission spectroscopy (XPS), and photoluminescence (PL) spectroscopy. The Raman spectra (excitation wavelength $\lambda_{exc} = 633$ nm) for the mono-, bi-, tetralayered and Bulk WS$_2$ nanosheets are shown in Fig. 3a. The first-order E$_{2g}^{1}$ and A$_{1g}^{1}$ modes were observed, which correspond to in-plane and out-of-plane vibrations at 355.4 cm$^{-1}$ and 418 cm$^{-1}$, respectively. In addition, the second-order mode was also observed, which corresponds to the second-order longitudinal acoustic mode, 2LA(M), at 350 cm$^{-1}$. It has been reported in detail that the changes in the ratio of the relative peak intensity, I(E$_{2g}^{1}$)/I(A$_{1g}^{1}$), and peak distances between the E$_{2g}^{1}$ and A$_{1g}^{1}$ modes depend on the number of layers (one, two, or four for mono-, bi-, or tetralayers, respectively, of WS$_2$). To clearly check the agreement of the results with those in the previous report, the Raman peaks were fitted using a Lorentzian function and the peak intensities and positions were extracted for the overlapping 2LA(M), and E$_{1g}^{2}$ and A$_{1g}^{1}$ modes for different number of layers. Fig. 3b shows the relative-peak-intensity ratio and peak distance as functions of the number of layers. The value of the ratio increased from 0.5 to 0.76 as the number of layers increased from one to four, which was mainly affected by the scattering volume change. The relative peak distance also increased from 62.6 cm$^{-1}$ to 64.4 cm$^{-1}$ with increasing number of layers owing to the increase in A$_{1g}^{1}$-mode frequency and the decrease in E$_{1g}^{2}$-mode frequency. These variations in Raman frequency in the A$_{1g}^{1}$ and E$_{1g}^{2}$ modes with the number of layers can be explained by the weak interlayer interaction and reduced long-range Coulomb interaction, respectively, between the effective charges caused by the increase in dielectric screening.

PL spectra of WS$_2$ with different number of layers are shown in Fig. 4a. The monolayered WS$_2$ exhibited a strong PL
signal at 2.01 eV, which is higher than any reported value for synthesized WS$_2$ but is consistent with the reported bandgap of monolayered WS$_2$. A weak and wide PL signal was observed at 1.97 eV for the bilayered WS$_2$ nanosheet, while there was a very weak signal related to the tetralayered WS$_2$ nanosheet. The red shift and weakness of PL signals as the number of layers increased from one to two were due to the transition from a direct bandgap to an indirect one, which agrees well with previous results on the PL signals’ dependence on the number of layers. These PL results again confirmed the layer controllability of the current process scheme.

Fig. 4b shows the XPS of monolayered WS$_2$ nanosheets. The W 4f core-level spectrum shows three peaks at 34.8, 32.7, and 38.0 eV, which correspond to the W 4f$_{5/2}$ and W 4f$_{7/2}$ levels for W$^{+4}$ and the W 5p$_{3/2}$ state, respectively. The S 2p core-level spectrum exhibits two peaks at 163.5 and 162.3 eV, which correspond to the S 2p$_{3/2}$ and S 2p$_{1/2}$ states. The samples showed no ClI-related peaks, indicating the high purity of the film. Based on the XPS, the W to S ratio was calculated to be 2, indicating good stoichiometry of the WS$_2$ nanosheet.

The crystallinity of the monolayered WS$_2$ nanosheet was analyzed by HRTEM. The HRTEM images in Fig. 5a–d show periodically arranged atoms in a selected region of the monolayered WS$_2$. The selected-area electron diffraction (SAED) pattern in the inset of Fig. 5a indicates that the monolayered WS$_2$ had a polycrystalline structure. Fast Fourier transforms (FFTs) were conducted to obtain the original image of the WS$_2$ nanosheet’s crystal structure (Fig. 5b), and six-fold coordination symmetry was clearly observed. To check the lattice constant of the WS$_2$ nanosheet, we performed inverse FFT by applying a mask. The lattice spacing (Fig. 5c) was determined from the intensity profile to be 0.32 nm, which is consistent with the value shown in a previous report.

Fig. 6. (a) Large-area (1 × 7 cm$^2$) mono-, bi-, and tetralayered WS$_2$ nanosheets on SiO$_2$ substrates. (b) Raman spectra collected from six different areas of the mono-, bi-, and tetralayered WS$_2$ nanosheets on SiO$_2$ substrates.

Mono-, bi-, and tetralayered WS$_2$ were then synthesized on (1 × 7)-cm$^2$ SiO$_2$ substrates (Fig. 6a) to confirm the good and high uniformity of CVD-WS$_2$. The as-grown WS$_2$ nanosheets showed a distinct color depending on the number of layers. The Raman spectra were measured at six different positions along the length of the nanosheet to confirm its large-area uniformity. As shown in Fig. 6b, it can be seen that the peak positions of the 2LA(M) and the E$^{1g}$ and A$^{1g}$ modes remain similar in all the spectra of the mono-, bi-, and tetralayered WS$_2$. The profiles of the Raman peaks, such as the full width at half maximum (FWHM) and the intensity, show negligible differences, as can be seen in Fig. 6b. This indicates that the synthesized thin film indeed possessed large-area uniformity over the (1 × 7)-cm$^2$ SiO$_2$ substrate.

Fig. 7. (a) OM images of graphene on a WS$_2$ nanosheet and a graphene–WS$_2$ hybrid photodetector. (b) PL spectra of graphene on a WS$_2$ nanosheet and graphene only. (c) Source-drain current ($I_d$) versus backgate voltage ($V_g$) for graphene on WS$_2$ nanosheet in the absence and presence of light. Inset: $I_d$ versus $V_g$ from 0 to 20 V for graphene on WS$_2$ in the absence and presence of light. (d) $I_d$ versus $V_g$ for graphene on WS$_2$ in the absence and presence of light. Inset: $I_d$ versus $V_g$ from 0.09 to 0.11 V for graphene on WS$_2$ in the absence and presence of light.
In a previous study, an exceptional photoresponse from graphene–TMD heterostructures, with large quantum efficiency and photocurrent generation, was reported mainly because of their large optical absorption (over $1 \times 10^7$ m$^{-1}$) and a visible range in the bandgap of TMDs. Thus, it would be interesting to evaluate the performance of photodetectors constructed with such materials. In particular, there has been no report on the photoresponse of CVD-synthesized graphene–TMD heterostructures. Fig. 7a shows the top view of the graphene–WS$_2$ heterostructure produced simply by depositing graphene via CVD on a monolayered WS$_2$ nanosheet. Fig. 7a shows a clear difference in the optical contrast of graphene on the WS$_2$ nanosheet. The corresponding PL spectrum shows a PL signal at 2.02 eV, which is consistent with the monolayered WS$_2$ PL, as shown in Fig. 7b. Fig. 7c shows the source–drain current $I_d$ between contacts located on graphene as a function of the backgate voltage $V_g$ for a typical graphene–WS$_2$ hybrid device. The positive shift of the Dirac point was due to the luminescence of SiO$_2$ as a result of hole doping. The device was then illuminated with a green light (wavelength = 550 nm). The red trace in Fig. 7c represents the variation in $I_d$ with $V_g$ when the device was illuminated continuously. The photocurrent ($I_p$) can be calculated from the difference between $I_d$, dark (80.6 µA) and $I_d$, light (77.44 µA) at $V_g = 0$ V and $V_d = 0.1$ V, and $I_p = 3.16$ µA was obtained. The negative sign of the photoconductivity at $V_g << V_{\text{dirac}}$ can be easily understood from the transfer of photoelectrons generated inside WS$_2$ to the graphene layer. However, the bands in WS$_2$ remained flat, suppressing the unidirectional transfer of electrons and equilibrating the charge distribution between graphene and WS$_2$. As shown in the insets of Fig. 7c and d, the ratio of the photocurrent and dark current ($I_p/I_d$, dark $= 0.04$) was lower than the ratio reported for graphene–exfoliated-MoS$_2$ owing to the trapped water between graphene and WS$_2$ from the graphene transfer process. However, considering the large-area uniformity of the WS$_2$ nanosheet, we believe that the graphene–WS$_2$ hybrid photodetector is a promising candidate for a new class of scalable optoelectronic devices.

Conclusions

We have reported successful preparation of CVD-grown WS$_2$ using gas-phase precursors with WCl$_6$ and H$_2$S. The synthesis process of CVD-WS$_2$ nanosheets showed systematic layer controllability ranging from a monolayer to tetrалayers by adjusting the growth time, thus improving the accessibility for fundamental research studies and industrial applications. Using AFM measurements, we calculated a growth rate of 5 min per layer; the initial growth mode was confirmed by SEM and AFM data as functions of synthesis time, where small triangular domains nucleated at random locations on the substrate, after which the nucleation sites continued to grow and formed boundaries when two or more domains met, resulting in a partially continuous film. The CVD-grown WS$_2$ nanosheet exhibited Raman shift and bandgap shift that depended on the number of layers and their good qualities were demonstrated in the PL, XPS, and TEM results. It was confirmed that the WS$_2$ nanosheet was deposited in mono-, bi-, and tetralayers on large-area (1 × 7 cm$^2$) SiO$_2$ substrates. Furthermore, we fabricated a graphene–WS$_2$ hybrid photodetector and evaluated its properties, which show the possibility of its application in optoelectronics. Our work is proof that gas-phase CVD will provide a foundation to further develop processes for the synthesis of TMD nanosheets.

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