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Improving the magnetic moment of Ca₂Ge and promoting the conversion of semiconductors to diluted magnetic semiconductors using Mn-doping

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The dilute magnetic properties of materials have important potential applications in the field of electronic science and technology. Intrinsic Ca₂Ge is a new environmentally friendly semiconductor material, and exhibits cubic and orthorhombic phases. The crystal structure characteristics of Ca₂Ge indicate that the modulation of its dilute magnetic properties can theoretically be achieved by doping with magnetic elements. The study of band structures shows that Ca₂Ge is a semiconductor, while Mn doped Ca₂Ge is a semi-metal. The results of density of states and atomic population analysis show that Mn doped Ca₂Ge exhibits ferrimagnetism with a magnetic moment of 5 μ_B, and the orbital splitting energy of the Mn atom is 1.0 eV. Mn-doping changes the cubic crystal field of Ca₂Ge, and the charge transfer and electron polarization of Ca d and Ge p orbitals are affected by Mn atoms. The Ca d orbital is split into d_{zz}, d_{yz}, d_{zx}, d_{xx-yy} and d_{xy} orbitals, and the contribution of spin of each d split orbital to the magnetic moment of the Ca d orbital is in the order d_{xy} > d_{yz} > d_{xz} > d_{xx-yy} > d_{zz}. The Ge p orbital is split into p_x, p_y and p_z orbitals, and the spin contribution of each p orbital to the magnetic moment of the Ge p orbital is in the order p_y > p_z > p_x. The analysis of atom populations shows that the charge transfer and spin of Ca and Ge change with Mn doping, and the difference between spin up and spin down increases, improving the magnetism of Ca₂Ge and forming a dilute magnetic semiconductor.

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1. Introduction

Diluted magnetic semiconductor materials (DMS) have semiconductor and electromagnetic properties, as a result of regulation of the charge freedom and spin freedom. Electronic devices using these properties have a wide range of potential applications in information processing, storage, transmission and other fields, and are of great significance to the development of electronic information technology.^{1–3} Semi-metallic materials are a particular type of DMS with a particular band structure, high electron polarization at the Fermi level, a total magnetic moment that is a multiple of the Bohr magneton, and a high Curie temperature T_C . Therefore, semi-metallic materials are preferred materials for manufacturing magnetic ion-implanted DMS electronic devices.^{4–8} Recently, investigations of DMS have mainly focused on III–V, IV, II–VI, oxide and sulfide semiconductors, such as GaAs, GaN, Si, Ge, ZnO, MoS₂, etc.^{9–13} The preparation of DMS mainly uses transition metal atoms or rare earth ions, which have spin magnetic moment,

doped to substitute some atoms in the semiconductor material.^{14–16} The magnetic moment of DMS is produced by the exchange of p state electrons of the anions and d state electrons of the magnetic ions, which are changed by the external temperature and doping concentration, so they may be ferromagnetic, sub ferromagnetic and antiferromagnetic.^{17,18}

As a new type of II–VI narrow band gap semiconductor, Ca₂Ge has high T_C , high carrier mobility, low dielectric constant, and excellent optical and other properties.¹⁹ Furthermore, it is compatible with conventional silicon processes and overcomes the physical quantum limitations of Si (conventional Si integrated circuit IC in the range of 2–20 nm), and has higher intrinsic hole mobility than GaAs and Si, so has potential application value in sensors and optoelectronic devices.^{20,21} The band structure and photoelectric characteristics of Ca₂Ge have been studied using the first principles method.²² It is found that cubic phase Ca₂Ge is a direct band-gap semiconductor with a band gap of 0.56 eV, and the lattice constant is 0.719 nm.²³ The stability of Ca–Ge alloys is affected by pressure and temperature, while the vibrational free energy depends only on temperature.^{24,25} Study of the structure, phonon spectrum, and lattice dynamics of Ca₂Ge found that the phonon correlation of Ca₂Ge is higher than that of Ca₂Si, which provides a theoretical basis for the use of Ca₂Ge in low power devices.²⁶

Recently, theoretical research has shown that the valence electron of Ca 3p⁶4s² is unstable in the Ca₂Ge electron system.

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In order to form a stable state, the electron is transferred from the 4s orbital to the 3d orbital, forming a stable structure, $3d^1 4s^1$.²⁷ Due to the particularity of Ca 3d state electrons, it is speculated that there is a correlation between Ca 3d state electrons and Ge 3p state electrons in Ca_2Ge compounds. The outer electron configuration of the central Ca ion is $3d^1$, and its unfilled d state electrons provide the valence bonding for the doped magnetic ions, forming a new electronic state, resulting in changes in the original electronic orbitals, electronic state, local structure, *etc.* Theoretically, this material has the possibility to become a DMS with excellent performance. Therefore, it is meaningful to develop Ca_2Ge as a new DMS material.

In order to clarify the magnetization mechanism and the influence of the physical field on the magnetism, the electronic structure and magnetic mechanism of intrinsic Ca_2Ge and Mn-doped Ca_2Ge are studied using a first principles method. The influence of the mechanism of Mn doping on the electronic structure and magnetic properties of Ca_2Ge is investigated extensively, and the regulatory mechanism for the electronic structure and magnetic properties is analyzed. It is found that the electronic orbitals of Ca and Ge form spin magnetic moments upon Mn-doping.

2. Modelling and crystalline field theory

2.1 Modelling

Cubic phase Ca_2Ge belongs to the $Fm\bar{3}m$ space group with lattice constants $a = b = c = 0.719$ nm, and $\alpha = \beta = \gamma = 90^\circ$.^{22,23,26–28} Each cell contains 8 Ca atoms and 4 Ge atoms. The doping model of $\text{Ca}_8\text{Mn}_1\text{Ge}_3$ is established by replacing a Ge with a Mn atom. The structural model is shown in Fig. 1. The same method is used to construct the $\text{Ca}_7\text{Mn}_1\text{Ge}_4$ doping model by replacing a Ca with a Mn atom. Since Ca is in two unequal positions in the Ca_2Ge system, there are two structures of $\text{Ca}_7\text{Mn}_1\text{Ge}_4$. The band structure of $\text{Ca}_8\text{Mn}_1\text{Ge}_3$ shows that it is a semi-metal, while $\text{Ca}_7\text{Mn}_1\text{Ge}_4$ is a metal, so this work just discusses the $\text{Ca}_8\text{Mn}_1\text{Ge}_3$ DMS.

The pseudopotential plane wave method based on first principles is used for calculation, and the calculation is completed using the Cambridge Serial Total Energy Package (CASTEP).^{29,30} The Broyden, Fletcher, Goldfarb and Shanno (BFGS) algorithm is used to optimize the geometrical structure of Ca_2Ge , and then six different approximation methods for generalized gradient approximations in the scheme of Perdew, Burke and Ernzerhof (GGA-PBE), revised Perdew–Burke–

Ernzerhof (GGA-RPBE), Perdew–Wang 91 (GGA-PW91), Wu and Cohen (GGA-WC), Perdew *et al.* (GGA-PBESol), and local density approximation (LDA) are used for the exchange correlation energy for finding the most stable structure and the most suitable approximate calculation method. All the calculations below are performed on the basis of this structure terminally. On this basis, the electronic structure and electromagnetic properties of intrinsic Ca_2Ge and Mn-doped Ca_2Ge are investigated. The ultra soft pseudo-potential (USPP) method is used to deal with the interaction between ions and electrons. The convergence threshold of interatomic force is 0.001 eV nm^{-1} , the internal stress is not more than 0.02 GPa, the convergence threshold of atomic displacement is 5×10^{-5} nm, the convergence standard of energy between two iterations is 5.0×10^{-7} eV per atom, and the convergence standard of self-consistent field is 2×10^{-6} eV per atom. The truncation energy is selected as 380 eV, and the K point in the Brillouin region using the Monkhorst–Pack method is selected as $4 \times 4 \times 4$. The valence electrons involved in the calculation are Ca $3s^2 3p^6 4s^2$, Ge $4s^2 4p^2$ and Mn $3d^5 4s^2$. Due to Mn being a transition metal, the Mn 3d orbit is not filled, and the enhanced electron delocalization leads to inaccurate calculation results. In order to improve the accuracy of the result, the GGA + U_{eff} (Hubbard U) method is used to correct the Mn 3d orbital, where the effective Hubbard parameter, U_{eff} , of the Mn atom is set to 3.0 eV. The electronic structure and magnetic properties of Mn doped Ca_2Ge are studied with U set to 3.0 eV.

2.2 Crystalline field theory

The bonding of metallic complexes is analogous to the action of positive and negative ions in ionic crystals. The metal–ligand interaction is of electrostatic character, and the ligand is treated as a point charge. The role of ligands is to establish a negative charge potential field, and the field is referred to as a crystalline field. Under the perturbation of the crystalline field, the d orbital of the metal undergoes energy level splitting, and the electrons fill the split d orbit from low to high, resulting in a decrease of the total energy after splitting. The splitting energy is called the crystal field stabilization energy.³¹

When the central ion with d electrons is considered, the Hamiltonian of the d electrons is written as:

$$\hat{H}_d = \sum_n \hat{h}(n) \quad (1)$$

Under the single-electron approximation, the electron–electron interaction term is not considered, and the Hamiltonian operator is defined as:

$$\hat{h}(n) = -\frac{1}{2} \nabla_n^2 - \frac{Z_c}{r_n} + (-1) \hat{V}(\vec{r}_n) = \hat{h}_0 - \hat{V}(\vec{r}_n) \quad (2)$$

where \hat{h}_0 is a Hamiltonian operator of d electrons in free ions, and $\hat{V}(\vec{r}_n)$ is the electrostatic field of the ligand.

Under the point charge model, the charge of a ligand is defined as $-q$, and the metal–ligand distance is a . Using results of classical electrostatics, the following equation can be obtained:

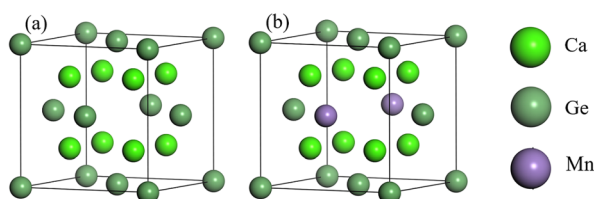


Fig. 1 Structure model of Ca_2Ge : (a) intrinsic Ca_2Ge , (b) Mn-doped Ca_2Ge .



$$\hat{V} = -\frac{6q}{a} - \frac{35q}{4a^5} \left(x^4 + y^4 + z^4 - \frac{3}{5}r^4 \right) \quad (3)$$

In a spherical coordinate system:

$$x^4 + y^4 + z^4 - \frac{3}{5}r^4 = r^4 \left[(\sin \theta \cos \varphi)^4 + (\sin \theta \sin \varphi)^4 + (\cos \varphi)^4 - \frac{3}{5} \right] \quad (4)$$

The single electron Schrodinger's equation is:

$$(\hat{h}_0 + \hat{h}') = \varepsilon_i \Psi_i \quad (5)$$

$$\hat{h}' = -\hat{V} = \frac{6q}{a} + \hat{V}_{O_h} \quad (6)$$

Using the perturbation method to approximately solve the single-electron Schrodinger's equation, the total electron wave function is obtained:

$$\Psi_d = N \Psi_1 \dots \Psi_n \quad (7)$$

The total electron energy is written as:

$$E_d = \sum_i \varepsilon_i \quad (8)$$

According to crystalline field theory, the symmetry of the crystalline field determines the mode of d orbital splitting, and the perturbation of the energy level splitting is defined as:

$$\hat{h}_0 \Psi_i^{(0)} = \varepsilon_i^{(0)} \Psi_i^{(0)} \quad (9)$$

$$\varepsilon_i = \varepsilon_i^{(0)} + \Delta \varepsilon_i \quad (10)$$

The first order perturbation correction is written as:

$$\Delta \varepsilon_i = \Delta \varepsilon_i^{(1)} = \langle \Psi_i^{(0)} | \hat{h}' | \Psi_i^{(0)} \rangle \quad (11)$$

The zero-order wave function is written as:

$$\Psi_i^{(0)} = \Psi_{nlm} = N_{nl} R_{nl}(r, \theta) Y_{lm}(\theta, \varphi) \quad (12)$$

In particular:

$$nd_{xx-yy} = \sqrt{\frac{15}{16\pi}} R_{n2}(r) \sin^2 \theta \cos 2\varphi \quad (13)$$

$$nd_{xy} = \sqrt{\frac{15}{16\pi}} R_{n2}(r) \sin^2 \theta \sin 2\varphi \quad (14)$$

The first order perturbation correction is written as:

$$\Delta \varepsilon_{e_g} = \langle nd_{xx-yy} | \hat{h}' | nd_{xx-yy} \rangle \quad (15)$$

Similarly:

$$\Delta \varepsilon_{t_{2g}} = \langle nd_{xy} | \hat{h}' | nd_{xy} \rangle \quad (16)$$

3. Results and discussion

3.1 Structure optimization

The crystal structure of Ca₂Ge is optimized by GGA-PBE, GGA-RPBE, GGA-PW91, GGA-WC, GGA-PBESol and LDA approximation methods. The system energy optimized by different approximation methods is shown in Table 1. It can be seen from Table 1 that the lattice constant obtained by the GGA-RPBE method is in good agreement with the experimental value and the energy value calculated is -8454.32 eV, lower than the energy obtained under other approximation methods, indicating that the structure of Ca₂Ge obtained using the GGA-RPBE method is the most stable.¹⁶ Therefore, the band structure, electron density of states, atomic population and orbital population are calculated using the GGA-RPBE method.

3.2 Band structure

The integral path of cubic phase Ca₂Ge is W-L-G-X-W-K along the direction of highly symmetric points in the Brillouin region under the plane wave basis set group. Fig. 2(a)–(d) show the band structure of the intrinsic Ca₂Ge and Mn-doped Ca₂Ge with spin up and spin down. It can be seen from Fig. 2(a) and (b) that the top and bottom of the spin up and spin down valence bands are both located at point X, forming a direct band gap $E_g = 0.57$ eV. The band structures of the spin up and spin down states are exactly the same. That is, the intrinsic Ca₂Ge has no magnetism, which is consistent with the results in reference.² Fig. 2(c) and (d) show the band structure diagram of Mn doped Ca₂Ge. It can be seen from the figure that the band structure of Mn doped Ca₂Ge with spin up forms a direct band gap of 0.32 eV at point X, while the band structure with spin down presents metallic characteristics with a band gap of 0 eV, indicating that Ca₂Ge with Mn doping forms a semi-metal. Comparing Fig. 2(a–d), it can be observed that the impurity bands exist for spin up and spin down when Mn is doped. In the spin up band structure, the influence of impurity bands is relatively small, and the band distribution is similar to that of intrinsic Ca₂Ge, while in the spin down band structure, the distribution of impurity bands has a greater impact. As a result, the distribution in the spin-down band structure is significantly different from that of the spin-up band structure, indicating that magnetism appears in Ca₂Ge after Mn doping. In order to explore the mechanism of magnetism and its source, the density of states is analyzed.

3.3 Electron state density

Fig. 3(a–e) show the partial density states (PDOS) of intrinsic Ca₂Ge and Mn doped Ca₂Ge. Fig. 3(a and b) show the PDOS of Ca and Ge in the intrinsic Ca₂Ge system, respectively. It can be seen from Fig. 3(a) that in the energy range of $-2-0$ eV and $0.6-3.1$ eV, contribution is mainly of Ca 3d state electrons. The contribution of Ca 3s state electrons is smaller, and the contribution of Ca 3p state electrons is zero. It can be seen from Fig. 3(b) that the valence band of intrinsic Ca₂Ge is mainly contributed to by Ge 2p state electrons in the energy range of $-2-0$ eV and by Ge 2s state electrons in the energy range of $0.6-$



Table 1 The system energy optimized by different approximation methods

Method	PBE	RPBE	PW91	WC	PBEsol	LDA
Energy/eV	-8448.27	-8454.32	-8453.83	-8538.13	-8431.4	-8452.17
$a/\text{\AA}$	7.201	7.182	7.20	7.148	7.148	7.00

3.5 eV. Therefore, the valence band of intrinsic Ca_2Ge is mainly contributed to by electrons of Ca 3d and Ge 2p states, and the conduction band is mainly contributed to by Ca 3d state electrons. In the valence band, the density of Ca 3d and Ge 2p states with spin up and down form symmetrical distributions, and in the conduction band, the density of Ca 3d states with spin up and spin down also form symmetrical distributions. Therefore, Ca and Ge do not exhibit magnetism in the intrinsic Ca_2Ge system. These results are consistent with the analysis results of the band structure of the intrinsic Ca_2Ge .

Fig. 3(c–e) show the PDOS of Ca, Ge and Mn in the Mn-doped Ca_2Ge system. It can be seen from Fig. 3 that in the energy range of -2 – 0 eV, the spin up PDOS is mainly contributed to by Ca 3d, Ge 2p and Mn 3d state electrons, while the spin down PDOS is mainly contributed to by Ca 3d and Ge 2p state electrons. In the energy range of 0 – 4 eV, the spin up PDOS is mainly contributed to by Ca 3d, Ge 2p, and Mn 3d state electrons, while the spin down

PDOS is mainly contributed to by Ca 3d, Ge 2p, and Mn 3d electrons. It can be seen from Fig. 3(e) that the density of spin-up and spin-down states presents an asymmetric distribution. In particular, a large spin-up density peak appears near -0.4 eV, and a strong spin-down density peak appears near 0.6 eV. The energy interval between the peaks of the two states is about 1.0 eV. This shows that Mn 3d is magnetic and the orbital splitting energy of Mn is 1.0 eV. At the same time, it is worth noting that when Mn is doped in Ca_2Ge , the distribution of the Ca and Ge state densities is affected by Mn, and the densities of the Ca and Ge states are asymmetric. Therefore, it can be inferred that the magnetism of the Mn-doped Ca_2Ge system is not only derived from Mn atoms. In order to further clarify the influence of Mn atom doping on the density of Ca and Ge states and the influence on the magnetic properties of the Ca_2Ge system, it is necessary to analyze the electron orbitals of each atom in the system. The atomic population and orbital population are discussed below.

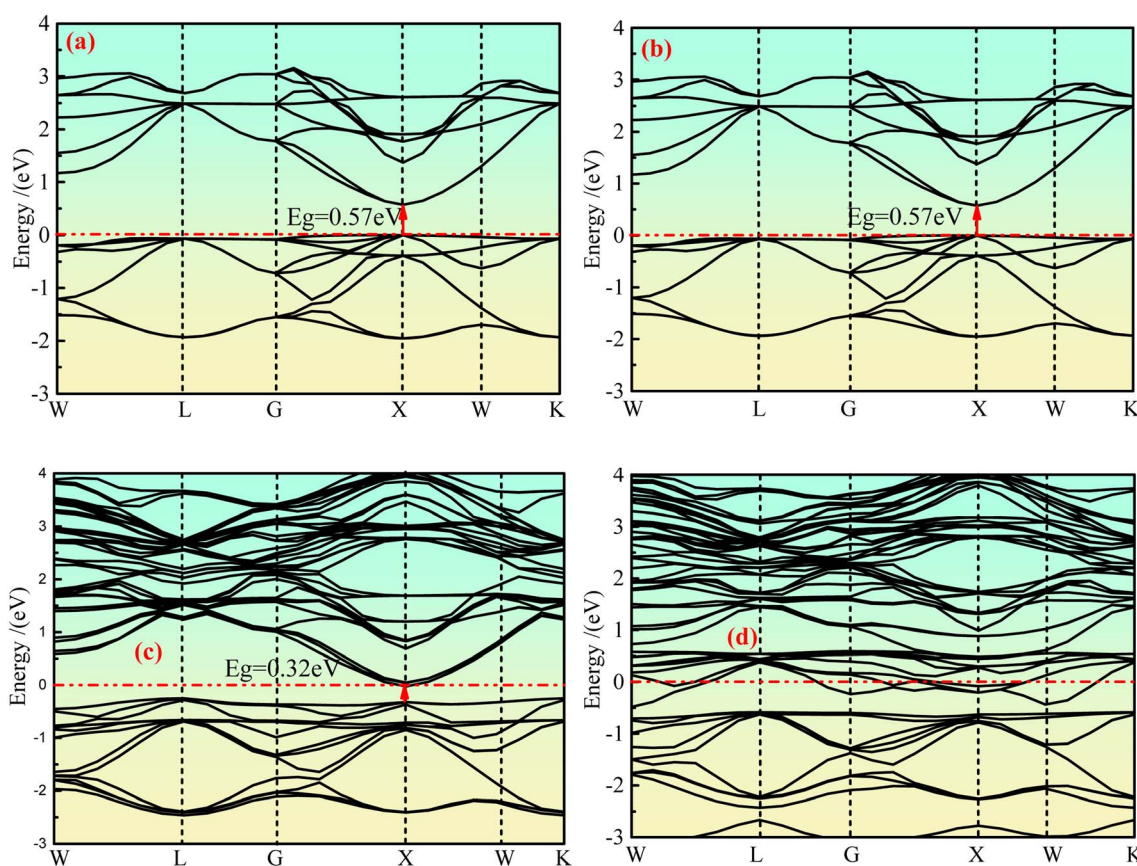


Fig. 2 The band structure of Ca_2Ge : (a) intrinsic with spin up, (b) intrinsic with spin down, (c) Mn-doping with spin up, (d) Mn-doping with spin down.



3.4 Atomic population

The atom population results show that the magnetic moment of the Mn atom is $5.09 \mu_B$, that of the Ge atom is $0.03 \mu_B$, and that of the Ca atom is $-0.12 \mu_B$. The results indicate that the main source

of magnetism is the spin polarization of Mn atoms. However, the polarization of Ge and Ca atoms is affected by the polarization of the Mn atom. The Ca atom exhibits antiferromagnetism, the Ge atom exhibits ferromagnetism, and Mn doped Ca_2Ge exhibits

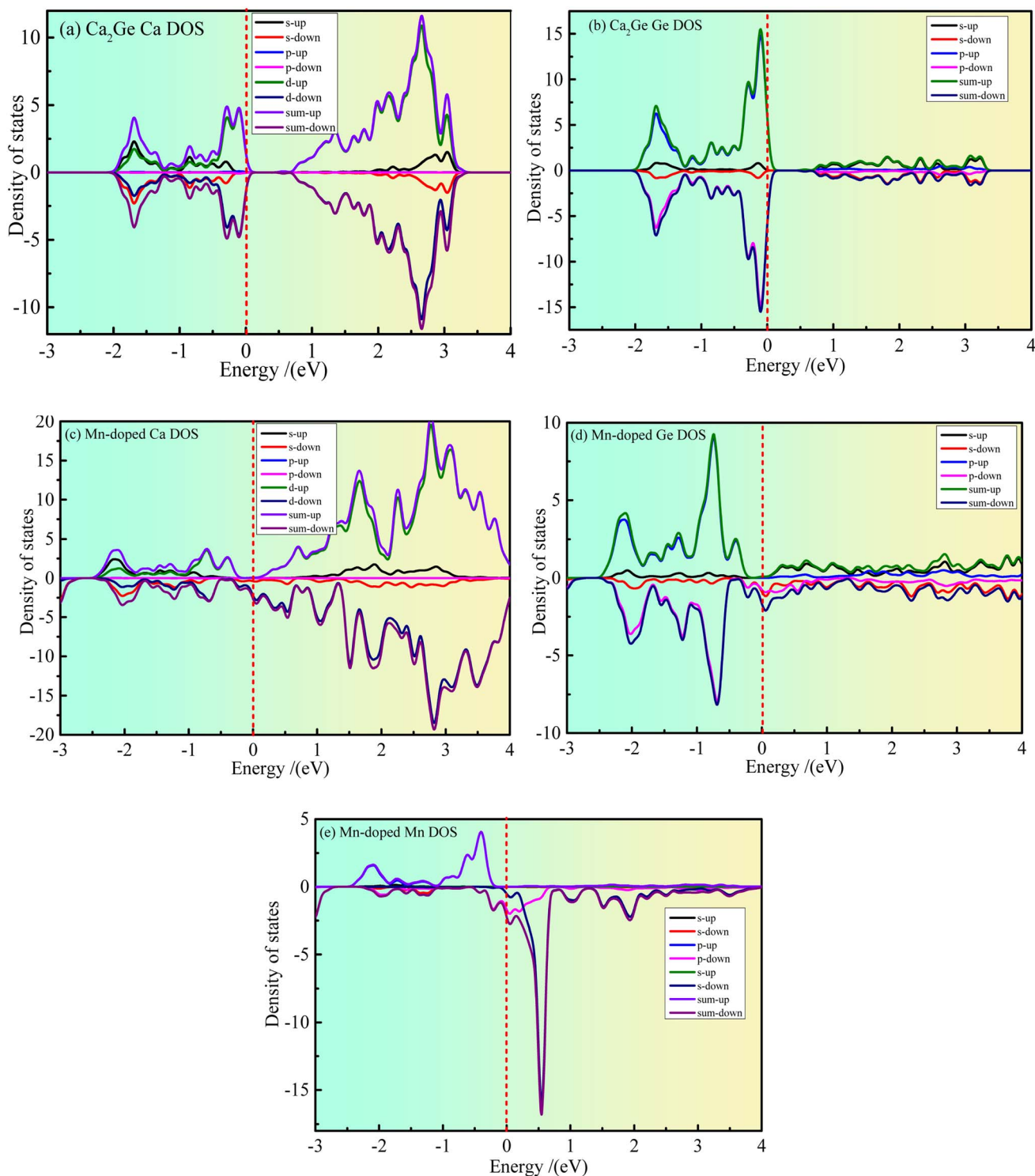


Fig. 3 The partial density of states of Ca_2Ge . (a) Ca partial density of states of intrinsic Ca_2Ge ; (b) Ge partial density of states of intrinsic Ca_2Ge ; (c) Ca partial density of states of Mn doped Ca_2Ge ; (d) Ge partial density of states of Mn doped Ca_2Ge ; (e) Mn partial density of states of Mn doped Ca_2Ge .



ferromagnetism. Fig. 4 shows the charge distribution of the electron orbital of Ca_2Ge . Fig. 4(a) and (c) show that the distribution of the electron orbital of Ca in the Ca_2Ge system changes from the original $3s^23p^64s^2$ to the more stable $3s^23p^63d^14s^1$. In the intrinsic Ca_2Ge , the spin up charges of the s, p and d orbitals are equal to those of spin down, for Ca and Ge. In the Mn doped Ca_2Ge system, the charges of the s, p and d electron orbitals with spin up and down Mn atoms are quite different. In particular, the Mn 3d orbital has the most significant difference of charge between spin up and spin down. In the doped Ca_2Ge system, the spin up and spin down charges of the s and p electron orbitals of the Ge atom are affected by the Mn atom, so that the charge of the s and p electron orbitals of each Ca atom is redistributed. For Ca 3d, the spin up charge increases from $0.25e$ to $0.34e$, and the spin down charge increases from $0.25e$ to $0.33e$, but the charges in the Ca 4s and Ca 3p orbitals are almost unaffected. However, Mn doping has different effects on each Ge atom, and the discussion of the change of the charge of the Ge atom is based on the average effect. For the Ge 3s orbital, the spin up charge decreases from $0.99e$ to $0.9e$, and the spin down charge decreases from $0.99e$ to $0.86e$. For the Ge 3p orbital, the spin up charge decreases from $2.04e$ to $1.91e$, and the spin down charge decreases from $2.04e$ to $1.84e$. The results show that the charge of Ca 3d is increased, and the charge of Ge 3s and Ge 3p decrease with Mn doping. Simultaneously, the charge revolutions of Ca and Ge are changed by Mn doping. In the intrinsic Ca_2Ge system, each Ca atom transfers and gains $1.02e$, and each Ge loses $2.05e$. However, each Ca atom gains

$0.85e$, the Mn atom loses $2.07e$, and the three Ge atoms lose $1.64e$, $1.57e$, and $1.58e$ with Mn doping.

3.5 Orbital populations

Fig. 5 shows the orbital energy level spin of Ca_2Ge . Fig. 5(a) shows that the Ca d orbital is split into d_{xy} , d_{yz} , d_{xz} , d_{xx-yy} and d_{zz} orbitals by the force of the crystal field in the intrinsic Ca_2Ge system. The maximum angle of the d_{xy} , d_{yz} and d_{xz} orbitals is between the axes, forming a plane-symmetric orbital and forming a t_{2g} orbital level with the same polarization. The maximum values of d_{zz} and d_{xx-yy} along the axial direction form an axially symmetric orbital and form an e_g orbital with the same polarization. t_{2g} and e_g have equal magnetic moments of spin up and spin down, so the d orbital does not exhibit magnetism. When Mn is doped, the spin up state of t_{2g} is split into d_{yz} , d_{xz} and d_{xy} orbitals, and d_{yz} is lower spin, while d_{xz} and d_{xy} are high spin. The spin up state of the e_g orbital is split into d_{zz} and d_{xx-yy} orbitals. The order of polarization of each orbital is as follows: d_{xy} , d_{xz} > d_{yz} > d_{zz} > d_{xx-yy} . The spin down states of the t_{2g} and e_g orbitals are split into d_{xy} , d_{yz} , d_{xz} , d_{xx-yy} and d_{zz} orbitals with different polarization. The order of polarization of each orbital is as follows: d_{yz} > d_{xz} > d_{xy} > d_{xx-yy} > d_{zz} . Thus, when Mn is doped, the contribution of spin splitting to the magnetic moment of the d orbital is in the following order: d_{xy} > d_{yz} > d_{xz} > d_{xx-yy} > d_{zz} . Fig. 5(b) shows the orbital level spin of Ge p. In the intrinsic Ca_2Ge system, the spin polarizations of the p_x , p_y and

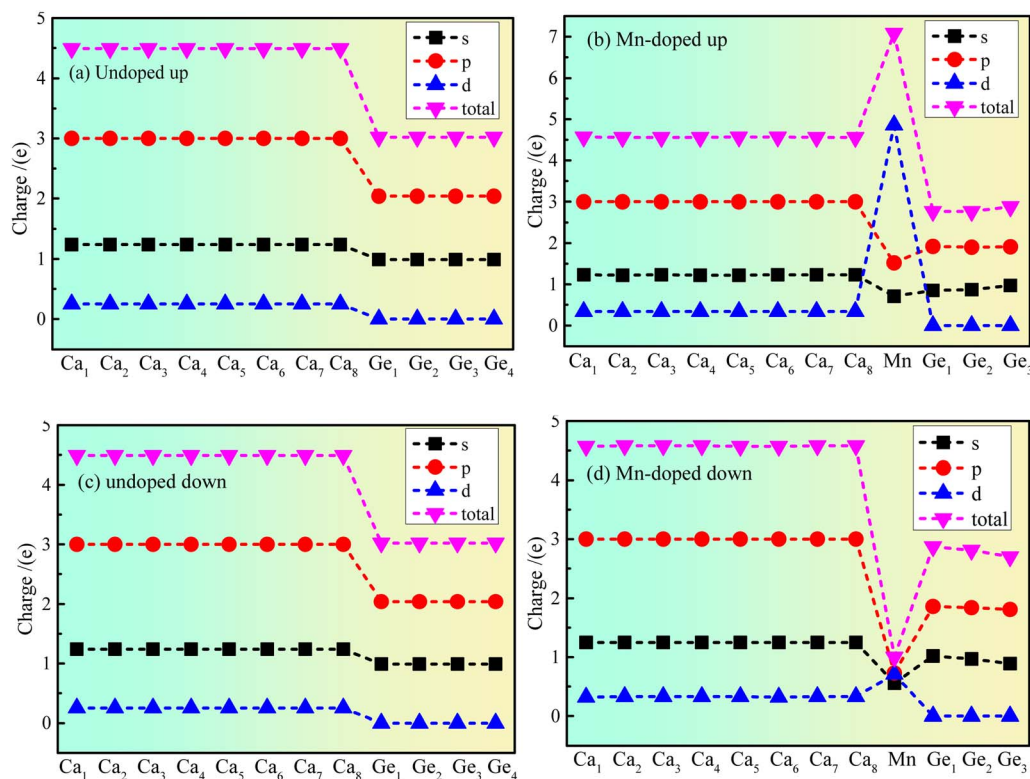


Fig. 4 The charge distribution of electron orbitals of Ca_2Ge : (a) spin up of intrinsic Ca_2Ge , (b) spin up of Mn doped Ca_2Ge , (c) spin down of intrinsic Ca_2Ge , (d) spin down of Mn doped Ca_2Ge .



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