


Cite this: *RSC Adv.*, 2021, **11**, 4035

Manipulating the Raman scattering rotation via magnetic field in an MoS₂ monolayer†

Yi Wan,^{†*ab} Xing Cheng,^{‡bc} Yanfang Li,^d Yaqian Wang,^d Yongping Du,^a Yibin Zhao,^a Bo Peng,^{id d} Lun Dai^{*bc} and Erjun Kan^{id *a}

Magneto-optical effects, which originate from the interactions between light and magnetism, have provided an important way to characterize magnetic materials and hosted abundant applications, such as light modulators, magnetic field sensors, and high-density data storage. However, such effects are too weak to be detected in non-magnetic materials due to the absence of spin degree of freedom. Here, we demonstrated that applying a perpendicular magnetic field can produce a colossal Raman scattering rotation in non-magnetic MoS₂, for A-mode representing the out-of-plane breathing vibration. Our experimental results show that linearly polarized scattering light is rotated by $\mp 125^\circ$, more apparent than the valley Zeeman splitting effect (∓ 1.2 meV) under the same experimental conditions (± 5 T), near room temperature. A detailed and systematic analysis on the polarization-resolved magnetic field-dependent micro-zone Raman intensity offers a feasible way to manipulate the inelastically scattered light via a magnetic technique. This explored phenomenology and physical mechanism arouse a new ramification of probing burgeoning magneto-optical effects in the field of two-dimensional laminar materials.

Received 3rd November 2020
Accepted 30th December 2020

DOI: 10.1039/d0ra09350e

rsc.li/rsc-advances

Introduction

Two-dimensional (2D) materials have received considerable attention and interest since the successful isolation of graphene for the first time.¹ In recent years, transition metal dichalcogenides (TMDCs) have become the flagship materials after graphene. The considerable extraordinary properties of TMDCs include the direct bandgap in visible and near-infrared wavelengths, suitable for exploring their potential in optoelectronics,² and the fascinating mechanism concerning the spin and the valley pseudospin, enabling unprecedented electronic devices or techniques.³ However, due to the limited thickness of such van der Waals (vdWs) materials, particularly for the monolayer, it is difficult to characterize the intrinsic properties

experimentally, which is inappropriate to the rapid development in this field.

Raman spectroscopy primarily focuses on the measurement of the inelastic scattering radiation generated from collective quasi-particle excitations.⁴ The crystal structure, atomic vibrating mode, and polarization angles between the incident light and detected light all have significant impacts on the intensity of Raman modes, making it an exceedingly effective technique for studying the crystal property and the property variation attributed from external perturbations. As reported in literature, Raman spectroscopy has been performed to examine the changes in the material property influenced by temperature,^{5,6} humidity,⁷ strain,⁸ externally applied electric field,⁹ or magnetic field.^{10,11} Along with the research upsurge in 2D magnets, Raman spectroscopy^{12–17} has been broadly adopted to characterize the magnetic property based on the surface sensitivity and remarkable magneto-optical effects. The successful applications of inelastic light scattering in such magnets has provided important information in studying magnetic impurity light scattering,^{18,19} defect magnetism,²⁰ spin-flip inelastically scattered light in diluted magnetic semiconductors,^{21,22} and lattice vibrations in colossal magneto-resistive materials.^{23–25} However, in most of the fabricated 2D vdWs materials, the absence of spin moments can only provide limited electronic or structural information relying solely on Raman techniques.

Molybdenum disulfide (MoS₂), a representative 2D semiconductor in the TMDC family, has recently drawn extensive attention.²⁶ In laminar MoS₂, the bulk crystals are formed by

^aDepartment of Applied Physics and Institution of Energy and Microstructure, Nanjing University of Science and Technology, Nanjing 210094, China. E-mail: wany@njust.edu.cn; ekan@njust.edu.cn

^bState Key Laboratory for Artificial Microstructure & Mesoscopic Physics, School of Physics, Peking University, Beijing 100871, China. E-mail: lundai@pku.edu.cn

^cCollaborative Innovation Center of Quantum Matter, Beijing 100871, China

^dNational Engineering Research Center of Electromagnetic Radiation Control Materials and State Key Laboratory of Electronic Thin Films and Integrated Devices, School of Microelectronics and Solid State Electronics, University of Electronic Science and Technology of China, Chengdu 610054, China

† Electronic supplementary information (ESI) available. See DOI: 10.1039/d0ra09350e

‡ These authors contributed equally to this work, and they should be viewed as first authors.



stacking covalently-bonded S–Mo–S sandwiched trilayers *via* relatively weak vdWs interactions. As the thickness is thinned from bulk to monolayer, on account of the 2D quantum confinement effect, the bandgap of MoS₂ will transfer from indirect to direct, with remarkable increase in the luminescence quantum efficiency.^{26,27} Briefly, with the distinct crystal structure enabling the degenerate-but-inequivalent valleys,^{28–30} remarkable photoluminescence (PL) in the visible frequency range^{26,27} and appropriate Gibbs free energy for large-scale hydrogen production applications,^{31,32} monolayer MoS₂ has offered considerable potential for the exploration of excellent electronic^{33,34} and optoelectronic^{35–37} devices. Herein, by performing polarization-resolved magneto-Raman spectroscopy measurements, we discovered a gigantic scattering rotation of A-mode in the Raman spectra obtained from MoS₂ monolayers in the presence of a perpendicular magnetic field. Although the well-known valley Zeeman splitting under the same conditions is quite weak as expected, the maximum Raman scattering feature rotation is about $\pm 125^\circ$, which has never been observed. This obvious magneto-optical effect originates from the symmetry breaking when the magnetic field is perpendicularly applied, particularly for the electron motion during the inelastically scattering process, which implies that the non-magnetic laminar materials can still exhibit an apparent response to the magnetic field. Thus, our study validates the magneto-Raman spectroscopy as a sensitive technique to characterize properties in quantum material systems, refines the microscopic inelastic scattering theory and demonstrates new avenues for modulating the optical signals based on the phonon polarity-related symmetry, providing a platform for the innovative magneto-optical devices and techniques based on 2D laminar materials, such as phonon lasers and nanomechanical oscillators.

Results and discussion

MoS₂ samples were synthesized directly on the surface of Si/SiO₂ substrates *via* the commonly used chemical vapor deposition (CVD) method.^{38,39} From the optical microscopy image (Fig. 1a) and the atomic force microscope (AFM) topography image (Fig. 1b), we can see that most MoS₂ domains on Si/SiO₂ appear as triangles with the side length of $28.1 \pm 12.8 \mu\text{m}$ (ESI Fig. S1†), obeying the symmetry of the hexagonal crystal lattice. The height profile along the MoS₂ triangular domain shows a terrace of $\sim 0.66 \text{ nm}$, as an indication of the monolayer. When imaged at a higher resolution, as shown in the lateral deflection signal obtained by AFM (Fig. 1c), the MoS₂ lattice possessing a honeycomb-like hexagonal arrangement of spots can be clearly resolved. This high-resolution AFM image was obtained *via* a conventional tip (TR400PSA) operating at ambient conditions, readily shown in the AFM lateral deflection images but only occasionally in the AFM topography images. It should be mentioned that the lattice-resolution image does not exhibit the accurate atomic-level resolution; on the contrary, the actual contact area between the scanned-cantilever and the MoS₂ sample surface is up to the order of $10 \times 10 \text{ nm}^2$; however, the periodic stick-slip atomic-level frictional forces are adequate to

capture the lattice-level resolution, sufficiently enough to determine the crystallographic orientation.^{40,41}

Bulk MoS₂ belongs to the space group of D_{6h} , possessing a dozen of vibration modes at the high-symmetric Γ point,⁸ four of which are first-order active, named in terms of the vibration mode symmetry, as E_{2g}^2 , E_{1g} , E_{2g}^1 , and A_{1g} . The E_{1g} mode is forbidden, while E_{2g}^1 and A_{1g} modes are allowed when the incident excitation light propagates along the MoS₂ c -axis as the light path configuration used in our experiments. The low-wavenumber E_{2g}^2 mode is not in view in this study. For the MoS₂ monolayer, the A - and E -mode symmetries will become A'_1 and E' . Besides, the monolayer nature of MoS₂ samples can be easily identified by distinguishing the peak position interval between the A'_1 and E' modes, which is prevailingly adopted to determine the thickness of ultrathin MoS₂ nanoflakes⁴² and confirmed by PL spectrum as a supplement.^{26,27} As a direct-gap semiconductor, the PL spectrum of the MoS₂ monolayer on Si/SiO₂ exhibits a strong excitonic emission (ESI, Fig. S2†).

The schematic of the experimental set-up for micro-zone Raman signal collection under the magnetic fields is shown in Fig. 2a. From the top view, the linear polarizer in front of the notch filter enables the collection of the polarization-resolved optical signal. During the process of signal collection, we began along the parallel direction ($\varphi = 0^\circ$) and continuously rotated the polarizer with the step of 10° until a cycle of 360° was finished. For compactness and simplicity, only 19 measured Raman spectra, along with the polarization angles between the incident light and the scattered radiation with the step of 20° are plotted in Fig. 1d. Two dominant peaks at 384 and 404 cm^{-1} , corresponding to the E' and A'_1 modes, are clearly observed. As shown in this series of Raman spectra (Fig. 1d), the breathing mode along the c -axis of S atoms (A'_1) exhibits a noteworthy polar dependence and can be negligible when the polarization angle reaches 90° , where the polarized incident light and the scattered light are orthogonal. In comparison, the intensities of the in-plane shearing motion (E') vary very slightly. The polar behavior of these two modes, regardless of the MoS₂ crystalline orientation, only shows the relationship with the polarization angles between the incident and scattered lights when no magnetic field is applied (ESI, Fig. S3†), which can be easily understood by the optical selection rules based on the symmetry of MoS₂ Raman tensors (more discussion later).⁸

As depicted from the side view of the magneto-Raman measurement experimental set-up (Fig. 2a), a superconducting magnet (attoDRY2100, attocube) was used to maintain the perpendicular magnetic field of up to $\pm 5 \text{ T}$. In our experiments, the incident and scattered lights, and the direction of the magnetic field are all perpendicular to the surface of the MoS₂ sample, *viz* the c -plane. The positive magnetic field is defined as the upward direction. In the presence of magnetic field, the time-reversal symmetry breaking will induce the A'_1 mode in the MoS₂ monolayer to be presented in the form of A' , while the E' mode remains unchanged. The Raman characteristic peak located at 521 cm^{-1} corresponding to the Si substrate is recorded as reference throughout the whole measurement process (ESI, Fig. S4†). Polarization-resolved Raman spectra through a full 360° were collected from the MoS₂ monolayer at selected



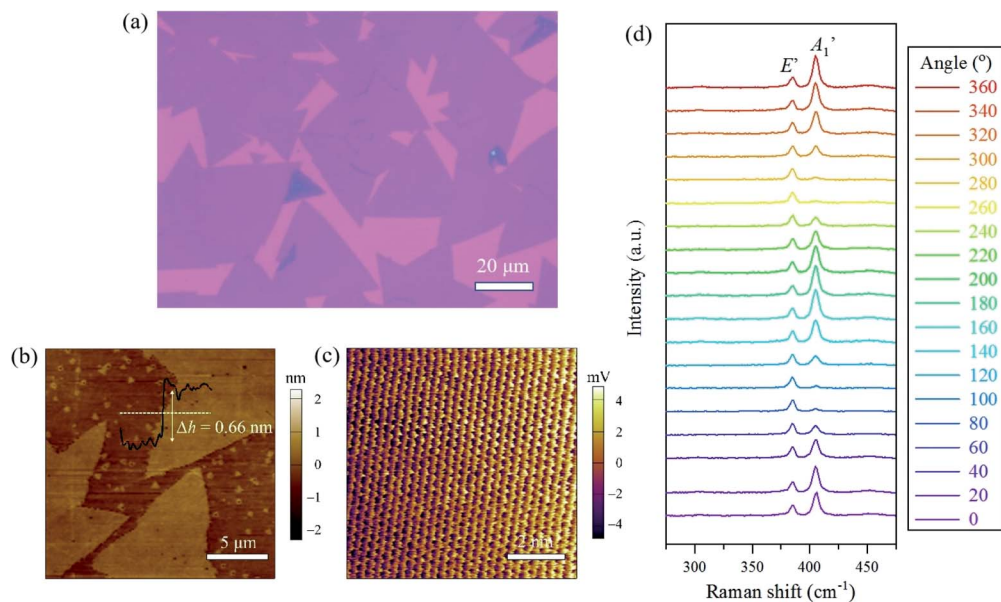


Fig. 1 (a) Optical microscopy image of the MoS₂ monolayers synthesized on the Si/SiO₂ substrate. (b) AFM topography image of MoS₂ monolayers synthesized on the Si/SiO₂ substrate, together with the corresponding height analysis. (c) Magnified AFM lateral deflection image of MoS₂, exhibiting the hexagonal lattice structure. (d) A series of Raman spectra of the MoS₂ monolayer, along with the polarization angles between the incident light and the scattered radiation.

magnetic fields. There exists no obvious difference in the frequency values of E' and A' modes no matter how the magnetic field strength changes. Since the E' mode remains constant and shows no direct relationship to the magnetic field under all polarization angle configurations, we mainly focus on the wavenumber ranging from 400 cm⁻¹ to 410 cm⁻¹ where the

A' mode is located. As shown in Fig. 2b, the A' mode exhibits a conspicuous response to the magnetic field. As the perpendicular magnetic field ranges from negative (−5 T) to positive (+5 T), the polar behavior of the A' mode shows a monotonous evolution. Polar plots of the E' and A' integrated intensities as a function of the polarization angles between the incident light

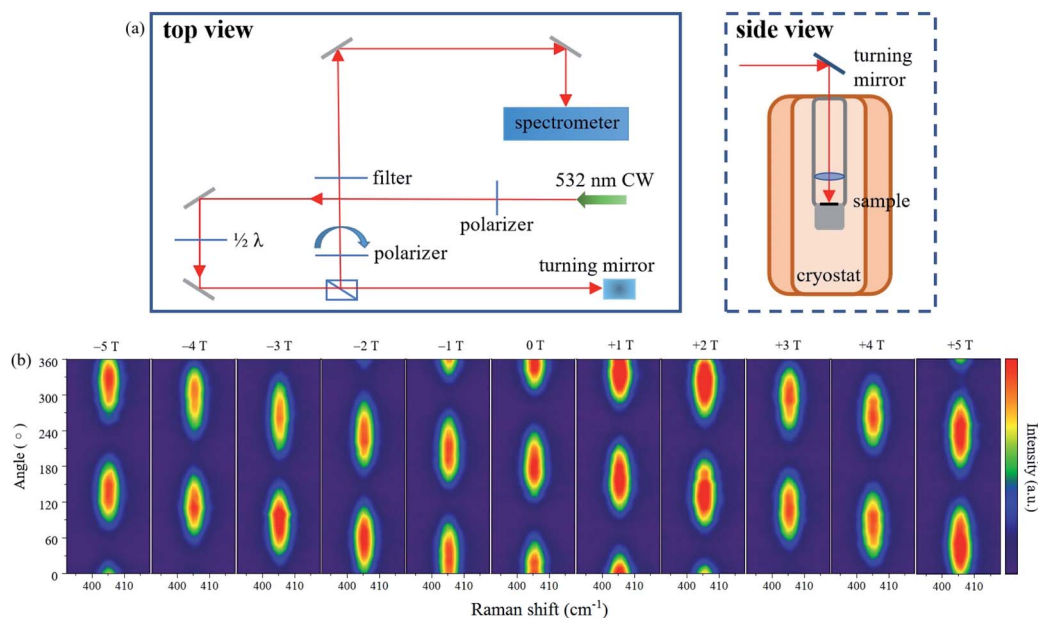


Fig. 2 (a) A schematic of the experimental set-up for micro-zone magneto-Raman measurements. Positive magnetic field is defined as the upward direction. The excitation laser is a 532 nm continuous wave (CW) laser. (b) Polarization-dependence of Raman spectra corresponding to the monolayer MoS₂ A' mode, measured at selected magnetic fields; the corresponding value of the applied perpendicular magnetic field is labelled above each contour image.

and the scattered radiation (Fig. 3a–k) can reveal the evolution trend of the polar behavior of these two modes more clearly. The Raman scattering pattern shows a gigantic rotation with the changes in the magnetic field strength. The linearly polarized inelastically scattered light rotates by $\sim \pm 125^\circ$ (± 2.2 rad) when the perpendicular magnetic field reaches ± 5 T (Fig. 3l).

MoS₂ has a honeycomb-like crystal lattice structure.¹⁰ The MoS₂ material possesses the point-group symmetry of D_{6h} in the bulk form and becomes D_{3d} in the bilayer form. The symmetry is further reduced to D_{3h} when the MoS₂ material decreases to monolayer, at which the center inversion symmetry is broken. The polarization dependence of the scattering intensity can be expressed as $I_s \propto |e_i \chi e_s|^2$, where e_i and e_s are the unit vectors describing the polarized states of the incident and scattered lights, respectively, and χ is the Raman tensor. In this study, the polarization of the incident light is fixed along the horizontal direction, while that of the scattered light is selected with an angle φ to the horizontal by a polarizer,

$$e_i = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, e_s = \begin{pmatrix} \cos\varphi \\ \sin\varphi \\ 0 \end{pmatrix}, \quad (1)$$

In consideration of

$$\chi_{E'}(D_{3h}) = \begin{pmatrix} d & 0 & 0 \\ 0 & -d & 0 \\ 0 & 0 & 0 \end{pmatrix}, \chi_{A'_1}(D_{3h}) = \begin{pmatrix} a & 0 & 0 \\ 0 & a & 0 \\ 0 & 0 & b \end{pmatrix}, \quad (2)$$

thus under zero field, E' and A'_1 modes follow d^2 and $a^2 \cos^2 \varphi$ dependence, respectively, and are clearly demonstrated by our fitting results, as shown in Fig. 3f. Here, d and a are the elements of the Raman tensor, and the absolute value of these two does not matter. For the MoS₂ monolayer, when $B \parallel c$ -axis, the point-group symmetry of D_{3h} will be further reduced compared to that of C_{3h} . In general, in the presence of magnetic field, the time-reversal symmetry breaking can add an anti-

symmetric component to the A-mode Raman tensor, shown as A' , which can partly explain the rotation in the scattered light polarization,

$$\chi_{A'}(C_{3h}) = \begin{pmatrix} a & c & 0 \\ -c & a & 0 \\ 0 & 0 & b \end{pmatrix}. \quad (3)$$

But it's unpredictable that this scattering rotation is remarkably large.⁴³

The preliminary analysis indicates that the giant rotation of the A' polar pattern closely related to the magnetic field requires a systematic investigation into the Raman scattering process. From a microcosmic viewpoint, the Raman scattering process includes the absorption of photons to activate the system to an excited state, the excitation of phonons *via* electron–phonon interactions in the intermediate step accompanied by decaying to a low-energy virtual state, and the emission of photons to return to the ground state. It is reasonable that if the intermediate states of electrons are perturbed by magnetic field, which is externally applied perpendicular to the MoS₂ c -plane, the property of radiated phonons during the Raman scattering process will make a difference. The intensity of the Raman scattering signal strongly relies on the second-order electronic susceptibility α , which can be expressed as a function of normal coordinates and external magnetic fields. The Lorentz force upon electrons induced from the magnetic field perpendicular to the sample c -plane will generate an additional transverse component to the in-plane electron motion. The behavior of the 2D electron motion can be described and quantified in the form of equations as

$$\begin{cases} \ddot{x} + \gamma\dot{x} + \omega_0^2 x + \frac{eB}{m}\dot{y} + \frac{e}{m}Ee^{i\omega t} = 0 \\ \dot{y} + \gamma\dot{y} + \omega_0^2 y - \frac{eB}{m}\dot{x} = 0 \end{cases}, \quad (4)$$

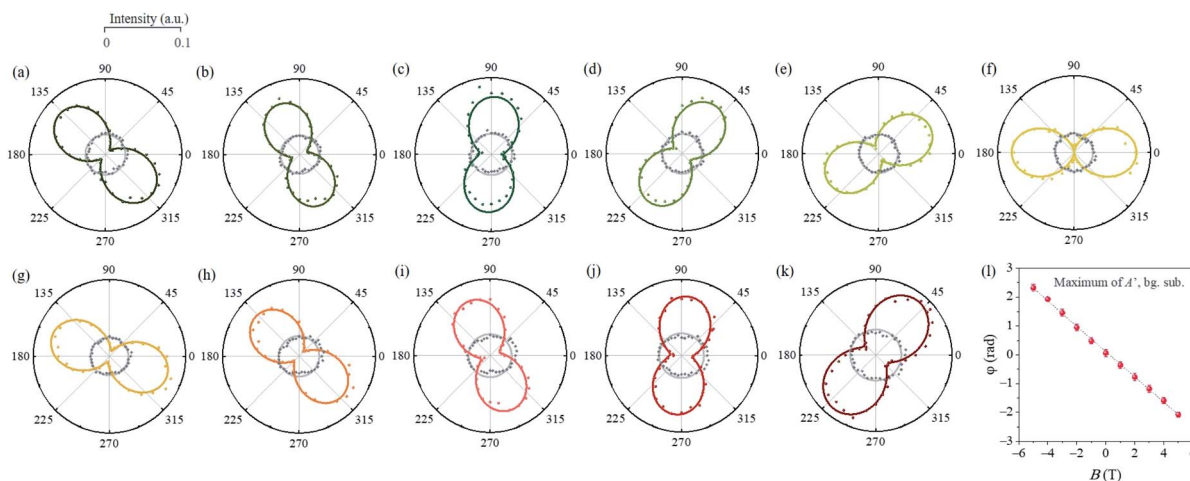


Fig. 3 Polar plots of the integrated intensities of E' and A' modes as a function of polarization angles between the incident light and the scattered radiation at $B =$ (a) -5 T, (b) -4 T, (c) -3 T, (d) -2 T, (e) -1 T, (f) 0 T, (g) $+1$ T, (h) $+2$ T, (i) $+3$ T, (j) $+4$ T, and (k) $+5$ T. (l) The polarization angle where the maximum of A' mode lies *versus* the perpendicular magnetic field strength. Systematic error induced by Faraday effects or Faraday rotation has been considered and excluded.



where E and B are the electric field vector of the incident light normal to the c -plane and the magnetic field, respectively. The rest of the parameters, e , m , γ , ω , and ω_0 are the charge, effective mass, scattering rate, incident light frequency, and excited electron restoring frequency, respectively.

Taking the Lorentz force into consideration, the dependence of the second-order electronic susceptibility on the magnetic field, $\alpha(B)$, can be derived based on the expressions provided in eqn (4) as

$$\begin{cases} \alpha_{xx}(B) = \alpha_{yy}(B) = -\frac{e\chi}{E} \frac{B_0}{B_0^2 + B^2}, B_0 = \frac{m\gamma}{e} \left(1 - i \frac{\omega_0^2 - \omega^2}{\omega\gamma}\right) \\ \alpha_{xy}(B) = \alpha_{yx}(B) = -\frac{e\gamma}{E} \frac{B}{B_0^2 + B^2} \end{cases} \quad (5)$$

where B_0 can be regarded as a resonant magnetic field where the influence on the polarization-resolved scattering light caused by the applied magnetic field reaches the peak. We can then obtain the complete forms of Raman tensors under a magnetic field, considering the crystal symmetry and the magnetic field influence together. The refined Raman tensor is represented in the following form as

$$\begin{aligned} \chi_{A'}(C_{3h}, B=0) &= \begin{pmatrix} a & c & 0 \\ -c & a & 0 \\ 0 & 0 & b \end{pmatrix} \rightarrow \chi_{A'}(C_{3h}, B \neq 0) \\ &= \begin{pmatrix} \alpha_{xx}(B) \cdot a & \alpha_{yx}(B) \cdot c & 0 \\ -\alpha_{xy}(B) \cdot c & \alpha_{yy}(B) \cdot a & 0 \\ 0 & 0 & b \end{pmatrix}. \end{aligned} \quad (6)$$

The polarization-resolved Raman intensity of the A' modes can finally be derived as

$$I_s \propto (a \cos \varphi)^2 \rightarrow I_s \propto \left| a \frac{B_0}{B_0^2 + B^2} \cos \varphi + c \frac{B}{B_0^2 + B^2} \sin \varphi \right|^2. \quad (7)$$

Since the integrated intensities of Raman modes rely on numerous factors, including excitation power density, integral time, and collection efficiency, the absolute values of the Raman tensor elements, a and c , make no great sense. By exploiting eqn (7) to fit the polarization-dependence in Fig. 3a–k, we obtained the value of $|B_0| = 3.25 \pm 0.23$ T, where the influence from the magnetic field on the polarization-resolved optical signal intensity approaches the maximum. It should be pointed out that if the magnetic field B reaches very high, the influence of Landau levels (LLs, $\hbar\omega_c = \hbar eB/m^*$) will become non-neglected (ESI, Fig. S5†). In short, the emergence of LLs with high degeneracy under magnetic fields will subtly modify the original band structures and perhaps remarkably affect the optical transitions, which may result in a Raman resonance if the allowed optical transition energetically matches LLs. In magneto-optical measurements, the magnetic field induces a beam deflection for light beams traversing along the magnetic field due to the Faraday effect, which can also be called Faraday rotation.^{43,44} We determined this Faraday rotation to be $-1.09^\circ \text{ T}^{-1}$ (ESI Fig. S6†) in our measurement system by analyzing the linearly polarized scattering light rotation on bare Si/SiO₂ substrates under various magnetic fields. The polarization

angle where the maximum of MoS₂ A' mode lies *versus* the perpendicular magnetic field strength is plotted in Fig. 3l, where the Faraday rotation has already been compensated.

We also performed the magneto-PL measurements on the MoS₂ monolayer^{45–48} (see the experimental set-up in ESI, Fig. S7†). The circular polarized magneto-PL spectra were recorded under a resonant excitation wavelength of 633 nm (HeNe laser). The selection of the excitation wavelengths is very important (ESI, Fig. S8†). The magnetic field is still applied perpendicularly to the MoS₂ sample c -plane. The circular polarization-resolved PL spectra (obtained at $B = -5$ T, 0 T, and +5 T) of the MoS₂ monolayer on the Si/SiO₂ substrate are shown in Fig. 4a. The left circularly polarized (σ^+) light couples to the K valley, and the right circularly polarized (σ^-) light couples to the K' valley. The A excitonic transition energies are extracted by the Lorentzian line shape fitting. For the MoS₂ monolayer on Si/SiO₂ at $B = 0$ T (Fig. 4a, middle plot), the ($\sigma^+ \rightarrow \sigma^+$) and ($\sigma^- \rightarrow \sigma^-$) PL spectra, where the excited and the detected light with the same chirality, are identical with each other, indicating no energy splitting in these two sets of degenerate-but-inequivalent valleys, required by the time-reversal symmetry. The A excitonic energy is extracted to be 1.82 eV, similar to the previously reported value. In contrast, under a negative magnetic field (−5 T), a small valley splitting can be discerned (Fig. 4a, top plot). The A excitonic transition at the K valley, corresponding to the ($\sigma^+ \rightarrow \sigma^+$) spectrum, shifts to a higher photon energy, whereas that at the K' valley, corresponding to the ($\sigma^- \rightarrow \sigma^-$) spectrum, shifts to a lower photon energy. The valley splitting is calculated and defined as the equation, $\Delta \equiv E(\sigma^+) - E(\sigma^-)$, where $E(\sigma^+)$ and $E(\sigma^-)$ represent the fitting peak energies in ($\sigma^+ \rightarrow \sigma^+$) and ($\sigma^- \rightarrow \sigma^-$) spectra, respectively. For the MoS₂ monolayer, Δ is +1.18 meV at −5 T, in agreement with the previously reported value. The valley splitting Δ changes sign and has a value of −1.05 meV at +5 T (Fig. 4a, bottom plot).

The origin for Zeeman splitting in the MoS₂ monolayer can be understood by studying the contribution of the spin, valley, and

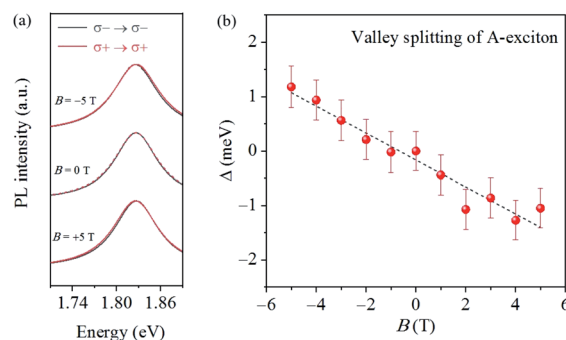


Fig. 4 (a) Normalized polarization-resolved PL spectra under the 0 and ± 5 T magnetic fields. The circular polarized magneto-PL spectra were measured, under resonant excitation wavelength of 633 nm (HeNe laser). The σ^- and σ^+ polarization-resolved spectra almost completely overlap in absence of magnetic field (0 T), and they split inversely under positive and negative magnetic fields (± 5 T). (b) Valley splitting (Δ) *versus* the magnetic field strength. For MoS₂ monolayer on Si/SiO₂ substrate, Δ changes with the magnetic field strength monotonously, with a linear slope of $\sim -0.24 \text{ meV T}^{-1}$.



atomic orbital magnetic moments to the electronic energy levels at K and K' valleys. Given the time-reversal symmetry in the absence of magnetic field (0 T), the energies of electrons at K and K' valleys are intrinsically degenerate. In the presence of a perpendicular magnetic field ($B \perp c$ -plane), this degeneracy is broken, resulting in an energy splitting between the two K and K' valleys due to the Zeeman effect. The detailed description is provided in the ESI, Fig. S9.† To elucidate the field-induced Zeeman splitting, we plot the field-dependent valley Zeeman splitting Δ in Fig. 4b. For the MoS₂ monolayer, Δ increases with the magnetic field at a linear slope of -0.24 ± 0.02 meV T⁻¹. The observed valley splitting can be expressed as $g\mu_B B$, where μ_B means the Bohr magneton, extracting the effective g factor to be 4.14 ± 0.17 , in agreement with reported values previously obtained from the differential reflection spectra.⁴⁷ However, as we can see from Fig. 4b, not only the value of the polarization-resolved photon energy splitting on the scale of milli electron-volt is very small but also the value error is considerable. This magnetic field-induced Zeeman splitting effect is much weaker than the observed Raman scattering rotation under the same measurement conditions. In contrast, Raman spectroscopy provides a more feasible and sensitive probe for exploring the magneto-optical effect induced by external perturbations.

Conclusions

In summary, we have successfully conducted polarization-resolved micro-zone Raman and PL spectroscopy investigation on the non-magnetic MoS₂ monolayer in the presence of magnetic fields. The relationship between the A-mode and the polarization angle between the excited and scattered lights can be efficiently modulated by the perpendicular magnetic field. A gigantic linearly polarized inelastically scattered light rotation ($\sim 125^\circ$) as well as the modest Zeeman splitting (~ 1.2 meV) were observed. The experimental observation results and analysis reveal the mechanism of the magneto-Raman effect in MoS₂, which is related to the behavior of electron motion in the presence of a perpendicular magnetic field, originating from the symmetry breaking induced by the external field, which dominates the electron motion during the scattering process. Such a tremendous magneto-optical effect at a monolayer scale and the underlying technological potentials for novel magneto-optical devices should arouse a new ramification of the research upon the inelastic scattering light. Our findings demonstrate the ultrathin vdWs materials as a unique platform for investigating numerous ground-breaking magneto-optical effects, which are sensitive to external perturbations.

Methods

Material preparation

Monolayer MoS₂ was synthesized on 285 nm Si/SiO₂ substrates via the CVD method inside a two-zone tube furnace using S (99.999%) and MoO₃ (99.99%) powders as the solid-state precursors, PTAS organics as the seeding promoter by catalyzing reduced metal oxides onto substrate surfaces, and high-purity Ar as the transport gas.

Sample characterization

The detailed morphology of MoS₂ was investigated by AFM (Cypher S, Asylum Research). The room-temperature polarization-dependent Raman spectra of MoS₂ were collected on a commercial Raman confocal system (Alfa300R, WITec) excited by a 2.33 eV (532 nm) laser and dispersed by a spectrometer equipped with a diffraction grating of 1800 grooves per mm. The excitation beam was focused onto the sample surface by a 100 \times objective lens. The emitted and scattered lights were collected by the same lens. The PL and Raman spectroscopy mapping images were captured by the same system.

Magneto-Raman measurement

The magneto-Raman spectra were captured by a home-built optical measurement system in the Faraday geometry based on a closed-cycle helium cryostat (attoDRY2100, attocube). The perpendicular magnetic fields of ± 5 T were generated and applied by the superconducting solenoidal magnet around the vacuum cryostat. The 2.33 eV excitation laser spot of 3 μ m diameter was focused at normal incidence onto the sample surface by a 50 \times objective lens. In magneto-Raman measurements, the optical signals were dispersed by a spectrometer (Andor SR500) equipped with a diffraction grating of 1200 groove/mm and a slit of 10 μ m, enabling the wavelength resolution < 0.1 nm with a 25 μ m pixel charge-coupled device (CCD) detector. In consideration of the thermal conduction and the stability of the superconducting magnets, all the magneto-optical spectra in this study were taken at 270 K, slightly lower than room temperature.

Conflicts of interest

The authors declare no competing financial interest.

Acknowledgements

This work was supported by the National Natural Science Foundation of China (No. 12004182, 51522206, 11774173, 11574151, 51790492, 61521004, 11804160 and 61874003), Jiangsu Province Science Foundation for Youth (No. BK20200481 and BK20170821), and the Fundamental Research Funds for the Central Universities (No. 30920021152, 30915011203, 30918011334 and 30919011403).

References

- 1 H. Yu, X. Cui, X. Xu and W. Yao, *Natl. Sci. Rev.*, 2015, **2**, 57–70.
- 2 Q. H. Wang, K. Kalantar-Zadeh, A. Kis, J. N. Coleman and M. S. Strano, *Nat. Nanotechnol.*, 2012, **7**, 699.
- 3 X. Xu, W. Yao, D. Xiao and T. F. Heinz, *Nat. Phys.*, 2014, **10**, 343.
- 4 G. Aivazian, Z. Gong, A. M. Jones, R.-L. Chu, J. Yan, D. G. Mandrus, C. Zhang, D. Cobden, W. Yao and X. Xu, *Nat. Phys.*, 2015, **11**, 148–152.
- 5 S. Sahoo, A. P. S. Gaur, M. Ahmadi, M. J. F. Guinel and R. S. Katiyar, *J. Phys. Chem. C*, 2013, **117**, 9042–9047.



- 6 M. Thripuranthaka, R. V. Kashid, C. Sekhar Rout and D. J. Late, *Appl. Phys. Lett.*, 2014, **104**, 081911.
- 7 J. Zhao, N. Li, H. Yu, Z. Wei, M. Liao, P. Chen, S. Wang, D. Shi, Q. Sun and G. Zhang, *Adv. Mater.*, 2017, **29**, 1702076.
- 8 Y. Wang, C. Cong, C. Qiu and T. Yu, *Small*, 2013, **9**, 2857–2861.
- 9 B. Chakraborty, A. Bera, D. V. S. Muthu, S. Bhowmick, U. V. Waghmare and A. K. Sood, *Phys. Rev. B*, 2012, **85**, 161403.
- 10 J. Ji, A. Zhang, J. Fan, Y. Li, X. Wang, J. Zhang, E. W. Plummer and Q. Zhang, *Proc. Natl. Acad. Sci. U. S. A.*, 2016, **113**, 2349.
- 11 L. Du, Z. Jia, Q. Zhang, A. Zhang, T. Zhang, R. He, R. Yang, D. Shi, Y. Yao, J. Xiang, G. Zhang and Q. Zhang, *2D Mater.*, 2018, **5**, 035028.
- 12 K.-z. Du, X.-z. Wang, Y. Liu, P. Hu, M. I. B. Utama, C. K. Gan, Q. Xiong and C. Kloc, *ACS Nano*, 2015, **10**, 1738–1743.
- 13 X. Wang, K. Du, Y. Y. F. Liu, P. Hu, J. Zhang, Q. Zhang, M. H. S. Owen, X. Lu, C. K. Gan and P. Sengupta, *2D Mater.*, 2016, **3**, 031009.
- 14 J.-U. Lee, S. Lee, J. H. Ryoo, S. Kang, T. Y. Kim, P. Kim, C.-H. Park, J.-G. Park and H. Cheong, *Nano Lett.*, 2016, **16**, 7433–7438.
- 15 B. Huang, G. Clark, E. Navarro-Moratalla, D. R. Klein, R. Cheng, K. L. Seyler, D. Zhong, E. Schmidgall, M. A. McGuire and D. H. Cobden, *Nature*, 2017, **546**, 270.
- 16 C. Gong, L. Li, Z. Li, H. Ji, A. Stern, Y. Xia, T. Cao, W. Bao, C. Wang and Y. Wang, *Nature*, 2017, **546**, 265.
- 17 Y. Deng, Y. Yu, Y. Song, J. Zhang, N. Z. Wang, Z. Sun, Y. Yi, Y. Z. Wu, S. Wu and J. Zhu, *Nature*, 2018, **563**, 94.
- 18 A. Oseroff and P. S. Pershan, *Phys. Rev. Lett.*, 1968, **21**, 1593–1596.
- 19 K. Y. Choi, S. A. Zvyagin, G. Cao and P. Lemmens, *Phys. Rev. B*, 2004, **69**, 104421.
- 20 J. Dillon Jr, H. Kamimura and J. Remeika, *J. Phys. Chem. Solids*, 1966, **27**, 1531–1549.
- 21 D. L. Peterson, A. Petrou, M. Dutta, A. K. Ramdas and S. Rodriguez, *Solid State Commun.*, 1982, **43**, 667–669.
- 22 D. L. Peterson, D. U. Bartholomew, U. Debska, A. K. Ramdas and S. Rodriguez, *Phys. Rev. B*, 1985, **32**, 323–340.
- 23 V. B. Podobedov, D. B. Romero, A. Weber, J. P. Rice, R. Schreekala, M. Rajeswari, R. Ramesh, T. Venkatesan and H. D. Drew, *Appl. Phys. Lett.*, 1998, **73**, 3217–3219.
- 24 V. A. Amelichev, B. Güttler, O. Y. Gorbenko, A. R. Kaul, A. A. Bosak and A. Y. Ganin, *Phys. Rev. B*, 2001, **63**, 104430.
- 25 M. N. Iliev and M. V. Abrashev, *J. Raman Spectrosc.*, 2001, **32**, 805–811.
- 26 K. F. Mak, C. Lee, J. Hone, J. Shan and T. F. Heinz, *Phys. Rev. Lett.*, 2010, **105**, 136805.
- 27 A. Splendiani, L. Sun, Y. Zhang, T. Li, J. Kim, C.-Y. Chim, G. Galli and F. Wang, *Nano Lett.*, 2010, **10**, 1271–1275.
- 28 T. Cao, G. Wang, W. Han, H. Ye, C. Zhu, J. Shi, Q. Niu, P. Tan, E. Wang, B. Liu and J. Feng, *Nat. Commun.*, 2012, **3**, 887.
- 29 H. Zeng, J. Dai, W. Yao, D. Xiao and X. Cui, *Nat. Nanotechnol.*, 2012, **7**, 490–493.
- 30 K. F. Mak, K. He, J. Shan and T. F. Heinz, *Nat. Nanotechnol.*, 2012, **7**, 494–498.
- 31 T. F. Jaramillo, K. P. Jørgensen, J. Bonde, J. H. Nielsen, S. Hørch and I. Chorkendorff, *Science*, 2007, **317**, 100.
- 32 Y. Wan, Z. Zhang, X. Xu, Z. Zhang, P. Li, X. Fang, K. Zhang, K. Yuan, K. Liu, G. Ran, Y. Li, Y. Ye and L. Dai, *Nano Energy*, 2018, **51**, 786–792.
- 33 B. Radisavljevic, A. Radenovic, J. Brivio, V. Giacometti and A. Kis, *Nat. Nanotechnol.*, 2011, **6**, 147–150.
- 34 M. Zhao, Y. Ye, Y. Han, Y. Xia, H. Zhu, S. Wang, Y. Wang, D. A. Muller and X. Zhang, *Nat. Nanotechnol.*, 2016, **11**, 954.
- 35 H. S. Lee, S.-W. Min, Y.-G. Chang, M. K. Park, T. Nam, H. Kim, J. H. Kim, S. Ryu and S. Im, *Nano Lett.*, 2012, **12**, 3695–3700.
- 36 O. Lopez-Sanchez, D. Lembke, M. Kayci, A. Radenovic and A. Kis, *Nat. Nanotechnol.*, 2013, **8**, 497.
- 37 X. Wang, P. Wang, J. Wang, W. Hu, X. Zhou, N. Guo, H. Huang, S. Sun, H. Shen, T. Lin, M. Tang, L. Liao, A. Jiang, J. Sun, X. Meng, X. Chen, W. Lu and J. Chu, *Adv. Mater.*, 2015, **27**, 6575–6581.
- 38 Y.-H. Lee, X.-Q. Zhang, W. Zhang, M.-T. Chang, C.-T. Lin, K.-D. Chang, Y.-C. Yu, J. T.-W. Wang, C.-S. Chang, L.-J. Li and T.-W. Lin, *Adv. Mater.*, 2012, **24**, 2320–2325.
- 39 J. Lee, S. Pak, P. Giraud, Y.-W. Lee, Y. Cho, J. Hong, A. R. Jang, H.-S. Chung, W.-K. Hong, H. Y. Jeong, H. S. Shin, L. G. Occhipinti, S. M. Morris, S. Cha, J. I. Sohn and J. M. Kim, *Adv. Mater.*, 2017, **29**, 1702206.
- 40 A. J. Marsden, M. Phillips and N. R. Wilson, *Nanotechnology*, 2013, **24**, 255704.
- 41 C. Lee, Q. Li, W. Kalb, X.-Z. Liu, H. Berger, R. W. Carpick and J. Hone, *Science*, 2010, **328**, 76.
- 42 S.-L. Li, H. Miyazaki, H. Song, H. Kuramochi, S. Nakaharai and K. Tsukagoshi, *ACS Nano*, 2012, **6**, 7381–7388.
- 43 B. Huang, J. Cenker, X. Zhang, E. L. Ray, T. Song, T. Taniguchi, K. Watanabe, M. A. McGuire, D. Xiao and X. Xu, *Nat. Nanotechnol.*, 2020, **15**, 212–216.
- 44 J. A. Stamper and B. H. Ripin, *Phys. Rev. Lett.*, 1975, **34**, 138–141.
- 45 A. Srivastava, M. Sidler, A. V. Allain, D. S. Lembke, A. Kis and A. Imamoglu, *Nat. Phys.*, 2015, **11**, 141.
- 46 G. Aivazian, Z. Gong, A. M. Jones, R.-L. Chu, J. Yan, D. G. Mandrus, C. Zhang, D. Cobden, W. Yao and X. Xu, *Nat. Phys.*, 2015, **11**, 148.
- 47 A. V. Stier, K. M. McCreary, B. T. Jonker, J. Kono and S. A. Crooker, *Nat. Commun.*, 2016, **7**, 10643.
- 48 Y. Wang, L. Deng, Q. Wei, Y. Wan, Z. Liu, X. Lu, Y. Li, L. Bi, L. Zhang, H. Lu, H. Chen, P. Zhou, L. Zhang, Y. Cheng, X. Zhao, Y. Ye, W. Huang, S. J. Pennycook, K. P. Loh and B. Peng, *Nano Lett.*, 2020, **20**, 2129–2136.

