

Cite this: *Nanoscale Adv.*, 2021, 3, 5581

Doping-free bandgap tunability in Fe₂O₃ nanostructured films†

Sujit A. Kadam,^a Giang Thi Phan,^{ID a} Duy Van Pham,^{ID *ab} Ranjit A. Patil,^{ID a} Chien-Chih Lai,^a Yan-Ruei Chen,^{ID c} Yung Liou^c and Yuan-Ron Ma^{ID *a}

A tunable bandgap without doping is highly desirable for applications in optoelectronic devices. Herein, we develop a new method which can tune the bandgap without any doping. In the present research, the bandgap of Fe₂O₃ nanostructured films is simply tuned by changing the synthesis temperature. The Fe₂O₃ nanostructured films are synthesized on ITO/glass substrates at temperatures of 1100, 1150, 1200, and 1250 °C using the hot filament metal oxide vapor deposition (HFMOVD) and thermal oxidation techniques. The Fe₂O₃ nanostructured films contain two mixtures of Fe²⁺ and Fe³⁺ cations and two trigonal (α) and cubic (γ) phases. The increase of the Fe²⁺ cations and cubic (γ) phase with the elevated synthesis temperatures lifted the valence band edge, indicating a reduction in the bandgap. The linear bandgap reduction of 0.55 eV without any doping makes the Fe₂O₃ nanostructured films promising materials for applications in bandgap engineering, optoelectronic devices, and energy storage devices.

Received 14th June 2021
Accepted 29th July 2021

DOI: 10.1039/d1na00442e

rsc.li/nanoscale-advances

Introduction

The bandgap is a fundamental parameter when selecting a material for various electronic and optoelectronic device applications as it governs the electronic transport and light-matter interaction properties. In the current research on optoelectronic materials, it is highly desirable to be able to tune the bandgap for developing highly efficient solar cells^{1–3} or enhancing the performance of LEDs and photodetectors.^{4–6} To achieve the desired bandgap for specific applications, doping has been considered as an effective method to tune the bandgap of semiconductors because it can form localized states in the bandgap;^{7–10} however, the intrinsic optical and electrical properties of host materials can be affected.^{10,11} Recently, some promising methods have been developed for tuning the bandgap of semiconducting materials without doping, such as strain engineering,^{12,13} thickness control,^{14,15} and synthesis temperature control.^{11,16} Bandgap tuning without any doping will be a very powerful advantage for applications in bandgap engineering and optoelectronic devices. The bandgap tuning possibility of semiconductors can be classified into three types.

In the first type of bandgap tuning, the highest energy level of the valence band (VB) does not change, while the lowest energy level of the conduction band (CB) is lowered. In the second type of bandgap tuning, the highest energy level of the VB is lifted, while the lowest energy level of the CB remains unchanged. The highest energy level of the VB and lowest energy level of the CB are raised and lowered, respectively, in the third type of bandgap tuning. The schematic illustration in Fig. 1 shows the three types of bandgap tuning. The bandgaps before and after

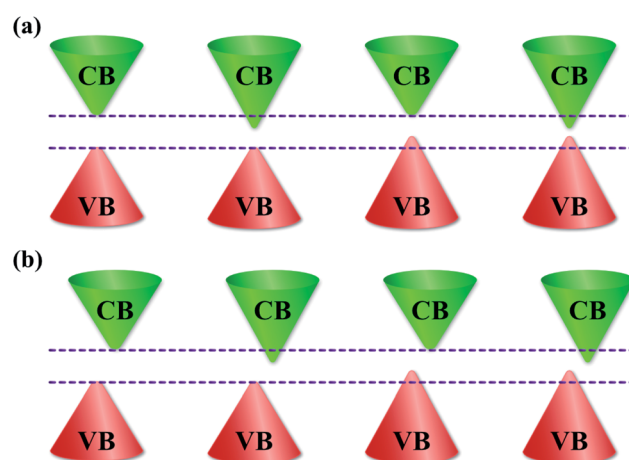


Fig. 1 Types of bandgap tuning. Bandgap before and after tuning for (a) direct and (b) indirect bandgap semiconductors. The bandgap tuning possibility can be classified into three types: (1) the highest levels of the VB remain, while the lowest levels of the CB are lowered; (2) the highest levels of the VB are raised, while the lowest levels of the CB remain unchanged; (3) the highest levels of the VB and the lowest level of the CB are raised and lowered, respectively.

^aDepartment of Physics, National Dong Hwa University, Hualien 97401, Taiwan. E-mail: duypham0611@gmail.com; ronma@mail.ndhu.edu.tw

^bCenter for Condensed Matter Sciences, National Taiwan University, Taipei 10617, Taiwan

^cInstitute of Physics, Academia Sinica, Taipei 11529, Taiwan

† Electronic supplementary information (ESI) available: Top- and side-view FESEM images of the nanostructured films before and after annealing, XRD patterns before annealing, XRD Rietveld analysis of the Fe₂O₃ nanostructured films, TEM analysis of the Fe₂O₃ thin film@1100 °C, etc. See DOI: 10.1039/d1na00442e



tuning of the direct and indirect bandgap semiconductors are shown in Fig. 1a and b, respectively.

Metal oxides with a wide bandgap are excellent candidates for use in optoelectronic devices.^{17,18} Among the numerous metal oxides, iron oxide is an interesting material due to its multiple oxidation states and stability of crystal phases, such as FeO,^{19,20} Fe₃O₄,^{21,22} Fe₇O₉,²³ α -Fe₂O₃,^{24,25} β -Fe₂O₃,^{26,27} and γ -Fe₂O₃.^{28,29} Iron oxide is widely used in a wide range of applications in nanodevices due to its unique properties, such as water splitting,^{24,25,30} supercapacitors,^{31–33} lithium-ion batteries,^{34–36} catalysts,^{37–39} photodetectors,⁴⁰ and gas sensors.^{41–43} In our proposed approach, we synthesized Fe₂O₃ nanostructured films on ITO/glass substrates at various synthesis temperatures. The results show that Fe₂O₃ nanostructured films contain two mixtures of Fe²⁺ and Fe³⁺ cations and two trigonal (α) and cubic (γ) phases. The increase of the Fe²⁺ cations and cubic (γ) phase in Fe₂O₃ nanostructured films with the elevated synthesis temperatures plays a critical role in tuning the bandgap of the Fe₂O₃ nanostructured films. The linear reduction in the bandgap from 1.65 to 1.10 eV without any doping of Fe₂O₃ nanostructured films provides a uniquely controllable route and makes it a promising candidate for applications in bandgap engineering, optoelectronic devices, and energy storage devices.

Experimental section

Synthesis of Fe₂O₃ nanostructured films

A larger area of FeO_x nanostructures was synthesized on ITO/glass ($\rho \sim 8 \Omega \text{ sq}^{-1}$) substrates using the HFMOVD technique. The HFMOVD technique is an important deposition method that allows the controllable synthesis of 1D metal-oxide nanostructures with different forms.^{11,16,44–48} By only control of the synthesis temperature, the metals with multiple oxidation states can form the non-stoichiometric metal oxide nanostructures, which contain the phase combination and cation mixture.^{11,16} This leads to a change in the bandgap of the metal oxide nanostructures without any doping with other metals, suggesting that the HFMOVD technique is suitable for bandgap engineering. Clean pure Fe wires (99.9%) with a diameter of ~ 1 mm were fixed between two copper (Cu) electrodes in the chamber. After the pressure of the vacuum chamber was pumped down and maintained at 0.3 Torr, the Fe wires were heated to 1100, 1150, 1200, and 1250 °C for 30 min to sublime hot Fe vapor. The residual oxygen in the chamber reacted with the hot Fe vapor to form iron oxide vapor. The iron oxide vapor immediately condensed to form Fe₃O₄ and metallic Fe nanostructures on the ITO/glass substrate, which was placed at a distance of ~ 4 mm above the Fe wire. The Fe₃O₄ and metallic Fe nanostructures were then annealed and oxidized to produce Fe₂O₃ nanostructured films in a controlled O₂ atmosphere of ~ 500 Torr at a temperature of 500 °C for 4 hours in a quartz tube furnace.

Characterization of the Fe₂O₃ nanostructured films

The morphological features of the Fe₂O₃ nanostructured films were observed using a field emission scanning electron

microscope (FESEM, JEOL JSM-6500F). The crystalline phase of the Fe₂O₃ nanostructured films was confirmed using an X-ray diffractometer (XRD, PANalytical X'Pert PRO). Crystalline analysis of the Fe₂O₃ nanostructured films was performed using a transmission electron microscope (TEM, JEOL JEM-2100) at an accelerating voltage of 200 kV. The simulated electron diffraction patterns were recorded using crystallographic software (Carine Crystallography version 3.1). The electronic structures and quantitative elemental information of the Fe₂O₃ nanostructured films were examined using an X-ray photoelectron spectrometer (XPS, Thermo Scientific Inc., K-alpha) with a microfocus monochromatic Al K α X-ray source. The optical transmittance and absorbance of the Fe₂O₃ nanostructured films were studied using a spectrophotometer (Ocean Optics DH-2000-BAL) with an ITO/glass substrate as a reference.

Results and discussion

The FESEM images in Fig. 2a–d show the side views of the large-area arrays of the Fe₂O₃ nanostructured films synthesized at various temperatures. The Fe₂O₃ nanostructured films are produced by the thermal oxidation of Fe₃O₄ and metallic Fe nanostructures, which are synthesized at temperatures of 1100, 1150, 1200, and 1250 °C. For convenience, the Fe₂O₃ nanostructured films synthesized at the various temperatures are denoted as Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C, respectively. Clearly, the sample@1100 °C is just the thin film, while the samples@1150, 1200, and 1250 °C not only possess thick buffer layers but also nanofibers and nanosheets, which are grown on the buffer layers, as shown in Fig. S1 (ESI[†]). The nanofibers and nanosheets are only formed on the samples@1150, 1200, and 1250 °C. This can be attributed to the presence of metallic iron in these samples before annealing. The existence of metallic iron has been confirmed by the XRD patterns, as shown in Fig. S2 (ESI[†]). During thermal oxidation, metallic iron reacts with the oxygen gas to increase the

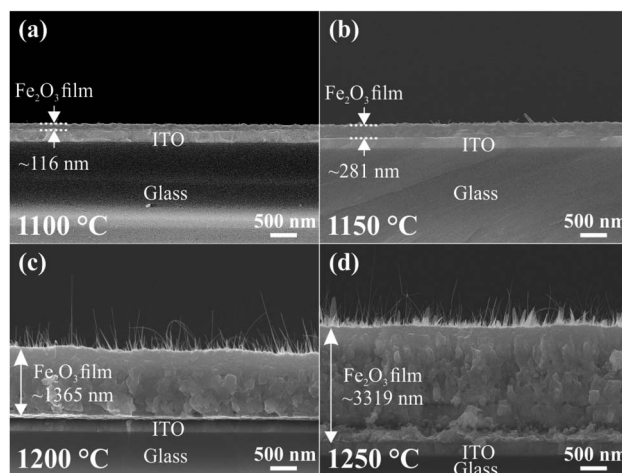


Fig. 2 Surface morphology. (a–d) FESEM images showing the side view of the Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C. The average thickness of the Fe₂O₃ thin film increases from ~ 116 to ~ 3319 nm with the elevated synthesis temperatures.



thickness of the Fe_2O_3 buffer layer as well as forming the Fe_2O_3 nanofibers and sheets. The average thickness of the Fe_2O_3 buffer layers increases from ~ 116 to ~ 3319 nm with elevated synthesis temperatures.

The crystalline phase of the Fe_2O_3 nanostructured films is analyzed by XRD and HRTEM measurements, and the results are shown in Fig. 3. The XRD patterns of ITO coated on the glass substrate and Fe_2O_3 nanostructured films@1100, 1150, 1200, and 1250 °C are displayed in Fig. 3a. For the Fe_2O_3 nanostructured films@1150, 1200, and 1250 °C, all the diffraction peaks reveal that the Fe_2O_3 nanostructured films contain two crystals together, which are the trigonal (α) phase and the cubic (γ) phase. Meanwhile, the observed diffraction pattern of the Fe_2O_3 thin film@1100 °C corresponds to all diffraction peaks of the trigonal (α) phase. The trigonal (α) crystal structure possesses a space group of $R\bar{3}c$ with lattice constants of $a = b = 0.5021$ nm, $c = 1.372$ nm, $\alpha = \beta = 90^\circ$ and $\gamma = 120^\circ$ (ICSD 201097), while the cubic (γ) crystal structure has a space group of $P4_332$ with lattice constants of $a = b = c = 0.8346$ nm and $\alpha = \beta = \gamma = 90^\circ$ (ICSD 87119). According to the Rietveld refinement results, the crystal composition percentage of the cubic (γ) phase increases linearly from 0 to $\sim 23.7\%$ with the elevated synthesis temperatures, as shown in Fig. S3 (ESI[†]). The trigonal (α) and cubic (γ) phases of the Fe_2O_3 nanostructured films have also been confirmed by TEM images. Fig. 3b–j show the single

nanosheet and nanowire, high-resolution TEM (HRTEM) images, the selected area electron diffraction (SAED) patterns, and the simulated electron diffraction (SED) patterns of the Fe_2O_3 nanostructured films@1250 °C. Fig. 3b–d show the low- and high-resolution TEM images, and the SAED patterns of the 2D Fe_2O_3 nanosheets. The only interplanar atomic spacing of 0.25 nm is found in the HRTEM in Fig. 3c, which is indexed to the (110) plane in the trigonal (α) phase. The SED patterns with a [001] zone axis in Fig. 3k are consistent with the SAED pattern in Fig. 3d, indicating that the 2D Fe_2O_3 nanosheets possess only the trigonal (α) phase. Fig. 3e–g show the low- and high-resolution TEM images, and the SAED patterns of the 1D Fe_2O_3 nanofibers. The HRTEM in Fig. 3f reveals that the 1D Fe_2O_3 nanofibers contain both trigonal (α) and cubic (γ) phases. The interplanar atomic spacings of 0.25 and 0.37 nm correspond to the (110) and (012) planes in the trigonal (α) phase, while the interplanar atomic spacings of 0.25 and 0.29 nm are consistent with the (113) and (022) planes in the cubic (γ) phase. Fig. 3h and i show the SED patterns for the trigonal (α) and cubic (γ) phases with the [221] and [211] zone axes, respectively. When the two SED patterns are completely overlapped, as shown in Fig. 3j, the superposed SED patterns match well with the SAED pattern in Fig. 3g, confirming that the 1D Fe_2O_3 nanofibers consist of the two trigonal (α) and cubic (γ) phases. Similar to 2D Fe_2O_3 nanosheets, the 2D Fe_2O_3 thin

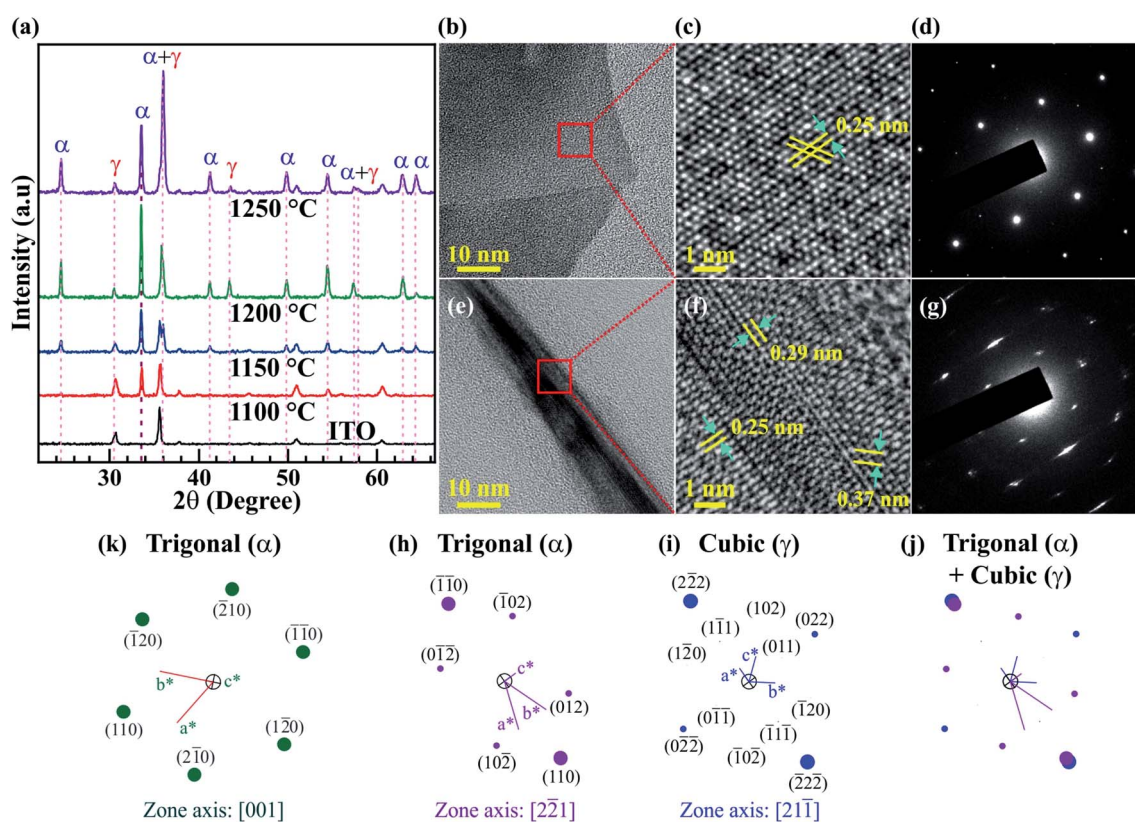


Fig. 3 Crystal structure determination. (a) XRD spectra of the ITO/glass substrates and the Fe_2O_3 nanostructured films@1100, 1150, 1200, and 1250 °C. Fe_2O_3 nanostructured films contain two crystal structures of the trigonal (α) phase and cubic (γ) phase. (b–g) Single nanosheet, single nanowire, HRTEM, and SAED patterns of the Fe_2O_3 nanostructured films@1250 °C. (h–k) Simulated electron diffraction patterns of trigonal (α) and cubic (γ) phases with various zone axes. TEM results confirm the existence of the cubic (γ) phase in nanofibers.



film@1100 °C contains only the trigonal (α) phase with an interplanar atomic spacing of 0.25 nm, which corresponds to the (110) plane. The low- and high-resolution TEM images and the SAED patterns of the 2D Fe₂O₃ thin film@1100 °C are shown in Fig. S4 (ESI[†]). The XRD and TEM results indicate that the cubic (γ) phase is only formed at temperatures higher than 1100 °C.

The chemical compositions and electronic structures of the Fe₂O₃ nanostructured films were further identified by X-ray photoelectron spectroscopy (XPS). The XPS spectra for Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C are shown in Fig. 4a–e. All the XPS spectra of the Fe-2p core-level in Fig. 4a–d are decomposed further into four peaks at binding energies of ~709.67, ~722.90, ~710.57, and ~724.14 eV, assigned to Fe²⁺ 2p_{3/2}, Fe²⁺ 2p_{1/2}, Fe³⁺ 2p_{3/2}, and Fe³⁺ 2p_{1/2}, respectively. In addition to the double peaks of Fe 2p_{3/2} and Fe 2p_{1/2}, the XPS spectra show the presence of two satellite (S₁ and S₂) and two surface (S_{p1} and S_{p2}) peaks. The binding energies of the Fe-2p core-level agreed well with the previous reports.^{49–51} Obviously, the Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C contain two Fe²⁺ and Fe³⁺ cations. According to quantitative XPS results, the amount of Fe²⁺ cations increases linearly from ~6.11 to ~13.5% with the elevated synthesis temperatures, as shown in Fig. 4e. In addition, the increase of Fe²⁺ cations is also confirmed using the energy-dispersive X-ray spectroscopy (EDS) line profile, as shown in Fig. S5 (ESI[†]). The Fe/O ratio increases linearly with increasing distance, indicating that the closer we get to the interface between the Fe₂O₃ film and ITO substrate,

the higher the amount of Fe²⁺ cations. The combination of the trigonal (α) and cubic (γ) phases and the two mixtures of the Fe²⁺ and Fe³⁺ cations may cause impacts on the bandgap of the Fe₂O₃ nanostructured films. Fig. 4f shows the photoluminescence (PL) spectra of the Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C. All the PL spectra were recorded in the wavelength range from 300 to 1000 nm. However, no PL peaks are observed in this region, indicating that the Fe₂O₃ nanostructured films have an indirect bandgap.

Optical spectroscopy was used to investigate the bandgap reduction of the Fe₂O₃ nanostructured films. Fig. 5 shows the photographs, transmittance, and absorbance spectra, and the bandgap of the Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C on the ITO/glass substrates. The photographs in Fig. 5a show the color of Fe₂O₃ nanostructured films, which changes from orange to deep red with elevated synthesis temperatures. The change in color of Fe₂O₃ nanostructured films is expected to produce a variety of absorbance and transmittance spectra. Fig. 5b and c show the transmittance and absorbance spectra for the Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C in the wavelength range from 250 to 1000 nm. Clearly, the minimum transmittance for each spectrum is less than 1%, occurs at a wavelength of ~315 nm, and extends to wavelengths of ~364, ~415, ~513, and ~610 nm for the nanostructured films@1100, 1150, 1200, and 1250 °C, respectively. As a result, the maximum absorbance for each spectrum of the Fe₂O₃ nanostructured films is increasingly broadened and the absorbance edge shifts towards longer

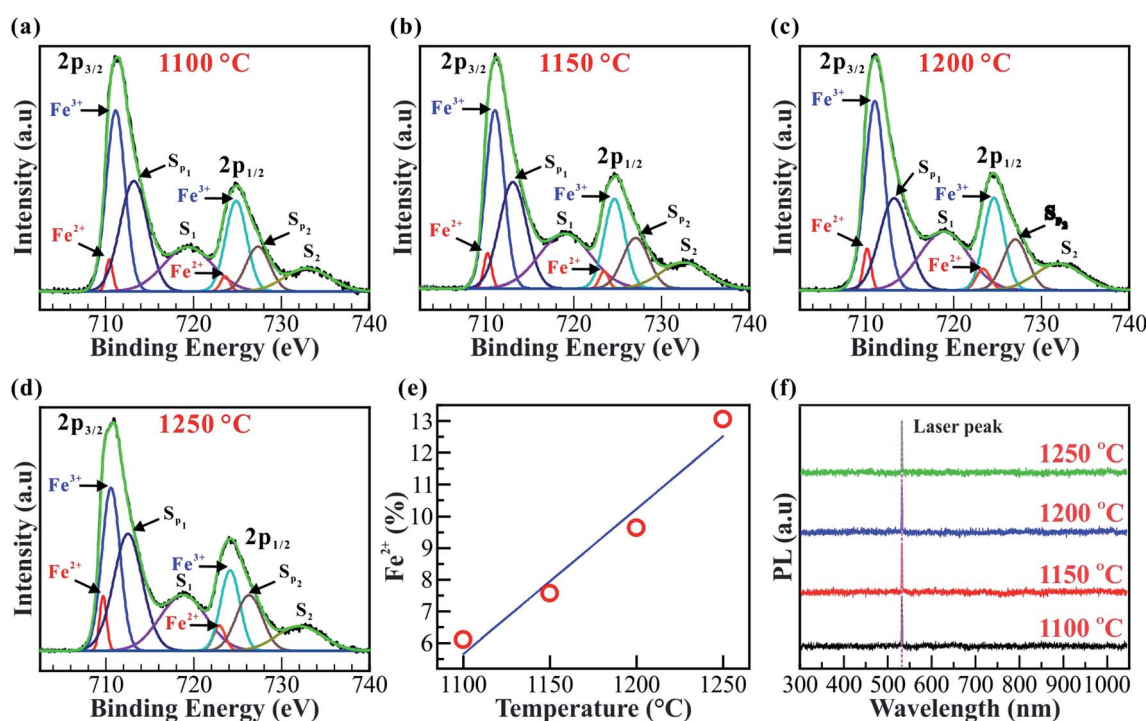


Fig. 4 Chemical and photoluminescence (PL) analysis. (a–d) XPS spectra of the Fe₂O₃ nanostructured films. All the XPS spectra consist of Fe²⁺ 2p_{3/2}, Fe³⁺ 2p_{3/2}, Fe²⁺ 2p_{1/2}, Fe³⁺ 2p_{1/2}, surface and satellite peaks. (e) The linear graphs show the increase of the Fe²⁺ cations from ~6.11 to ~13.5% with the elevation of synthesis temperature, respectively. (f) PL spectra of the Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C. No PL peaks are observed in the range from 300 to 1000 nm, indicating that the Fe₂O₃ nanostructured films have an indirect bandgap.



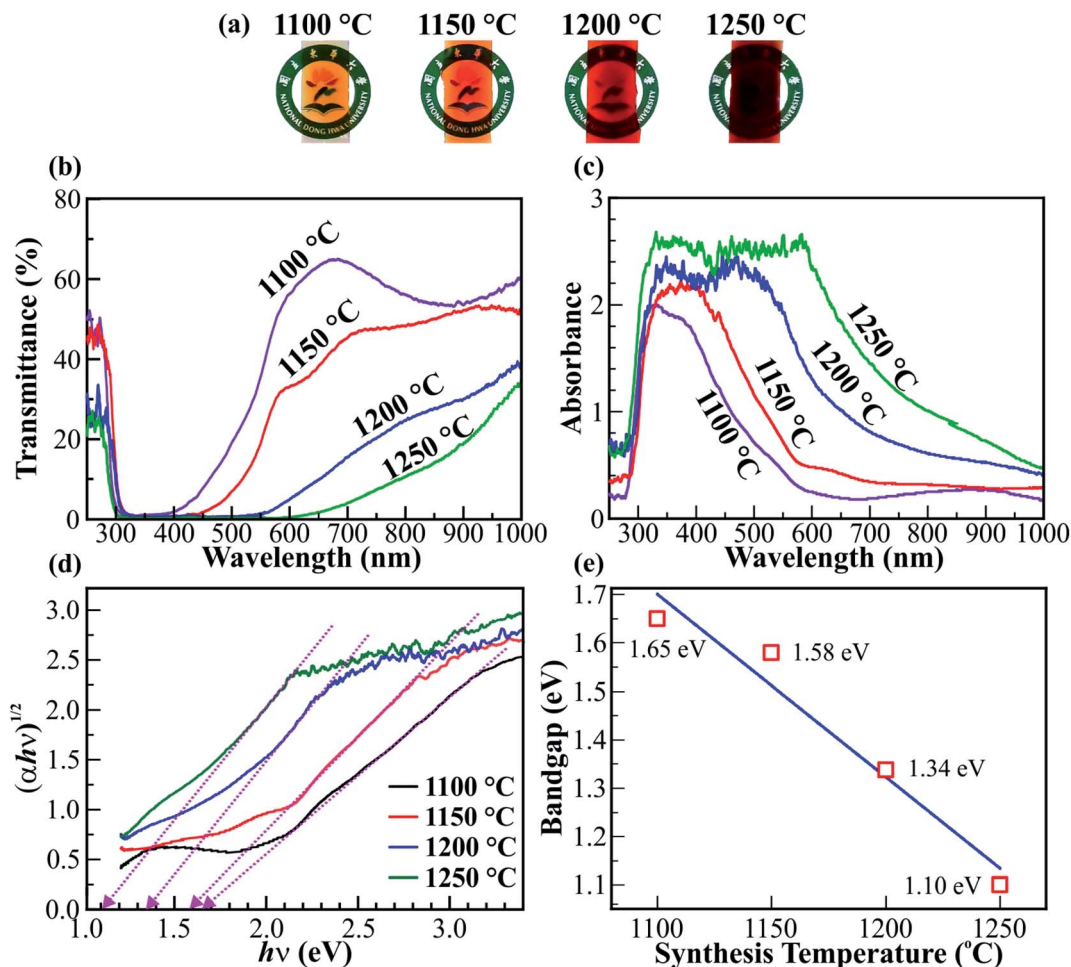


Fig. 5 Bandgap properties. (a) Photographs of the Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C on ITO/glass substrates. The color of the Fe₂O₃ nanostructured films changes from orange to deep red with the elevation of synthesis temperature. (b) The transmittance and (c) absorbance spectra of the Fe₂O₃ nanostructured films synthesized at various temperatures. (d) The linear graphs of $(\alpha h\nu)^{1/2}$ vs. $h\nu$ for determining the E_g of the Fe₂O₃ nanostructured films. (e) The graph shows the linear bandgap reduction of the Fe₂O₃ nanostructured films with the elevation of synthesis temperature.

wavelengths with the elevated synthesis temperatures. The red-shift of the absorbance edge is strongly related to the narrowing of the bandgap in the Fe₂O₃ nanostructured films. The indirect bandgap (E_g) of the Fe₂O₃ nanostructured films can be determined using the Tauc equation $\alpha h\nu = A(h\nu - E_g)^2$,⁵² where α is the absorption coefficient, h is Planck's constant, ν is the frequency of the incident light, and A is a constant. The absorbance spectra in Fig. 5c were transformed into the Tauc plot, as shown in Fig. 5d. Using the extrapolation method, the value of the bandgap (E_g) is determined as the intersection of the fitting dotted straight lines and $h\nu$ axis. Therefore, the bandgaps of Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C are 1.65, 1.58, 1.34, and 1.10 eV, respectively. The graph in Fig. 5e shows the temperature dependence of the bandgap of Fe₂O₃ nanostructured films. A linear fit between the bandgap values versus synthesis temperature indicates that the bandgap (E_g) of the Fe₂O₃ nanostructured films reduces linearly with the elevated synthesis temperatures. The linear bandgap reduction of 0.55 eV is due to the gain in the Fe²⁺ cations, the cubic (γ) phase in the Fe₂O₃ nanostructured films, and the thickness of

the Fe₂O₃ nanostructured films. However, the influence of thickness on the bandgap reduction of the Fe₂O₃ nanostructured films is only ~ 0.11 eV, as shown in Fig. S6–S9 (ESI†). The influence of the ITO/glass substrate on the optical properties of the Fe₂O₃ nanostructured films is also negligible (see Fig. S10, ESI†). The phase combination and cation mixture create structural disorder, leading to the formation of localized states. To determine the structural disorder of the Fe₂O₃ nanostructured films, the Urbach energy (E_u) was used for investigation purposes (see Fig. S11, ESI†). The linear increase of the E_u from 0.288 to 0.594 eV indicates the increase of localized states, leading to a significant bandgap reduction of the Fe₂O₃ nanostructured films. In addition, the bandgap reduction due to gain in metal cations and phase mixing has been demonstrated in previous studies.^{11,16} The bandgap reduction of 0.55 eV of the Fe₂O₃ nanostructured films is larger than that of Sb–BiVO₄ films,⁵³ Bi₂S₃ colloidal nanocrystals,⁵⁴ W₁₈O₄₉ nanorods,¹¹ polydopamine/Ag₃PO₄/graphene oxide,⁵⁵ and crystalline Sb₂S₃.⁵⁶



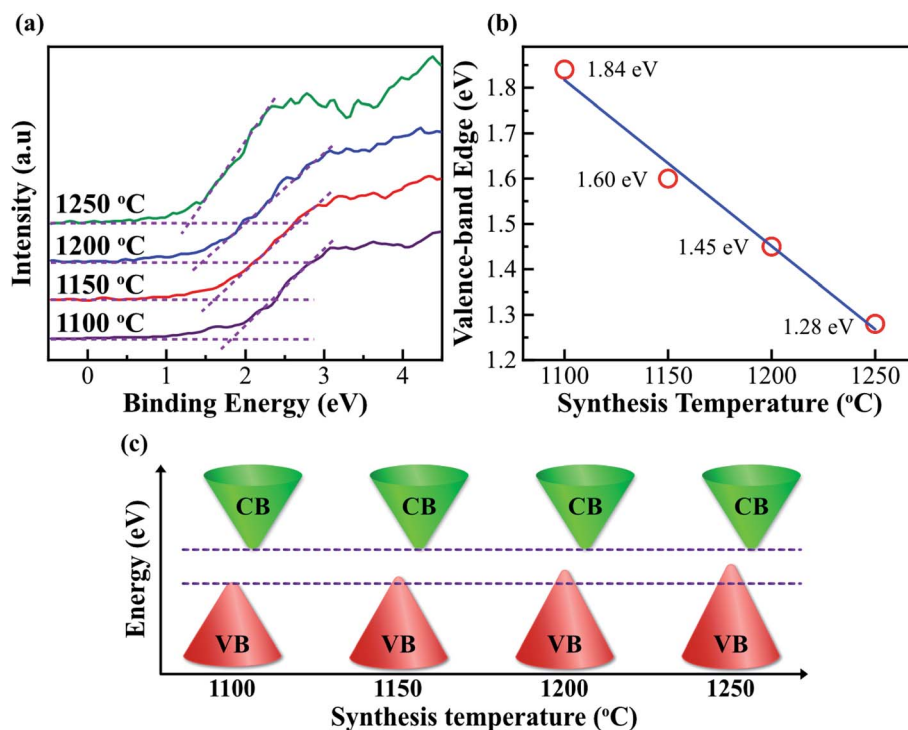


Fig. 6 Valence-band-edge lifting. (a) The high-resolution valence-band (VB) XPS spectra of the Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C. The binding energies of the VB edges are defined at the intersections of the dashed lines. (b) The graph shows the linear reduction of the VB edge with the elevation of synthesis temperature, indicating that the VB edge is lifted at high synthesis temperatures. (c) A schematic diagram shows the bandgap reduction of the Fe₂O₃ nanostructured films with the elevation of synthesis temperature. The characteristics of the bandgap reduction in the Fe₂O₃ nanostructured films are consistent with the second type of bandgap tuning, as shown in Fig. 1b.

In order to further understand the reduction of bandgap in Fe₂O₃ nanostructured films, the valence band (VB) is also investigated. Fig. 6a shows the VB XPS spectra of the Fe₂O₃ film-nanofibers. The VB edge is defined as the intersection of the two dashed lines.¹¹ The binding energies of the VB edge for Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C are located at 1.84, 1.60, 1.45, and 1.28 eV, respectively. Obviously, the binding energy of the VB edge reduces linearly with the elevated synthesis temperature, as shown in Fig. 6b. This means that the VBs of the Fe₂O₃ nanostructured films are lifted at the elevated synthesis temperature. The VB edge lift of 0.56 eV is completely matched with the reduction of bandgap in the Fe₂O₃ nanostructured films, as mentioned above. The highest levels of the VB are lifted, while the lowest levels of the CB are unchanged. As a result, the bandgap of Fe₂O₃ nanostructured films reduces with the elevated synthesis temperatures, as shown in the schematic illustration in Fig. 6c. The bandgap characteristics of the Fe₂O₃ nanostructured films are consistent with the second type of bandgap tuning, as shown in Fig. 1b. The reduction in the bandgap of Fe₂O₃ nanostructured films is only due to the gain in the Fe²⁺ cations and the cubic (γ) phase composition. This suggests that the bandgap (E_g) of the Fe₂O₃ nanostructured films can be tuned by controlling only the synthesis temperature. The characteristics of bandgap tuning without any doping make the Fe₂O₃ nanostructured films promising candidates for applications in bandgap engineering, optoelectronic devices, and energy storage devices.

The electrochemical charge storage properties of Fe₂O₃ nanostructured films@1100, 1150, 1200, and 1250 °C are studied *via* cyclic voltammetry (CV) and electrochemical impedance spectroscopy (EIS) in 1 M KOH electrolyte. The CV and EIS results are shown in Fig. S12 (ESI[†]). The 1D nanorods produced at 1100 °C show the lowest specific capacitance value of $\sim 82 \text{ F g}^{-1}$ at a scan rate of 50 mV s^{-1} , while those produced at 1250 °C show the highest specific capacitance of $\sim 352 \text{ F g}^{-1}$ at a low scan rate of 5 mV s^{-1} . These specific capacitance values are higher than those of Fe₂O₃ thin film,³³ Fe₂O₃ nanorods,^{33,57} Fe₂O₃ nanosheets,⁵⁷ Fe₂O₃ nanoparticles,⁵⁷ and nanoporous Fe₂O₃/CNT.⁵⁸ This result makes Fe₂O₃ film nanofibers a potential supercapacitor electrode material candidate.

Conclusion

In this study, Fe₃O₄ and metallic Fe nanostructures were synthesized on ITO/glass substrates using the HFMOVD technique at temperatures of 1100, 1150, 1200, and 1250 °C. The Fe₃O₄ and metallic Fe nanostructures were then fully oxidized using the thermal oxidation technique, yielding Fe₂O₃ nanostructured films at 500 °C for 4 hours in a ~ 500 Torr O₂ atmosphere. The formation of nanosheets and nanofibers can be attributed to the presence of metallic iron in the samples that were synthesized at temperatures of 1150, 1200, and 1250 °C. The Fe₂O₃ nanostructured films comprise two combinations of the trigonal (α) and cubic (γ) phases, and two mixtures of the



Fe²⁺ and Fe³⁺ cations. The crystal composition of the cubic (γ) phase and the amount of the Fe²⁺ cations increase linearly from 0 to ~23.7% and ~6.11 to ~13.5% with the elevated synthesis temperature, respectively, implying that the valence band will gain more electrons and be lifted at higher synthesis temperatures. Therefore, the phase combination and cation mixture play a critical role in tuning the bandgap of the Fe₂O₃ nanostructured films. The linear bandgap reduction of 0.55 eV is only controlled by the synthesis temperature without any doping, indicating the potential of the Fe₂O₃ nanostructured films for applications in bandgap engineering, optoelectronic devices, and energy storage devices.

Author contributions

Sujit A. Kadam: conceptualization, methodology, validation, investigation, data curation, writing – original draft. Giang Thi Phan: formal analysis, investigation, methodology, validation, writing – original draft. Duy Van Pham: data curation, methodology, validation, formal analysis, investigation, visualization, funding acquisition, writing – original draft, writing – review & editing. Ranjit A Patil: investigation, resources. Chien-Chih Lai: resources, software, funding acquisition. Yan-Ruei Chen: investigation. Yung Liou: investigation, resources. Yuan-Ron Ma: investigation, resources, validation, writing – review & editing, project administration, funding acquisition, supervision.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

The authors would like to thank the Ministry of Science and Technology (MOST), Taiwan for their financial support of this research under Contract No. MOST 108-2112-M-259-007-MY3, MOST 109-2112-M-259-008, and MOST 110-2811-M-002-651.

References

- 1 K. Gao, Y. Kan, X. Chen, F. Liu, B. Kan, L. Nian, X. Wan, Y. Chen, X. Peng, T. P. Russell, Y. Cao and A. K. Y. Jen, *Adv. Mater.*, 2020, **32**, 1906129.
- 2 R. Prasanna, A. Gold-Parker, T. Leijtens, B. Conings, A. Babayigit, H.-G. Boyen, M. F. Toney and M. D. McGehee, *J. Am. Chem. Soc.*, 2017, **139**, 11117–11124.
- 3 T. Todorov, T. Gershon, O. Gunawan, Y. S. Lee, C. Sturdevant, L.-Y. Chang and S. Guha, *Adv. Energy Mater.*, 2015, **5**, 1500799.
- 4 S.-W. Wang, H. Medina, K.-B. Hong, C.-C. Wu, Y. Qu, A. Manikandan, T.-Y. Su, P.-T. Lee, Z.-Q. Huang, Z. Wang, F.-C. Chuang, H.-C. Kuo and Y.-L. Chueh, *ACS Nano*, 2017, **11**, 8768–8776.
- 5 S. Wang, Y. Guo, D. Feng, L. Chen, Y. Fang, H. Shen and Z. Du, *J. Mater. Chem. C*, 2017, **5**, 4724–4730.
- 6 S. Engmann, A. J. Barito, E. G. Bittle, N. C. Giebink, L. J. Richter and D. J. Gundlach, *Nat. Commun.*, 2019, **10**, 227.
- 7 Y. Kim, H. Bark, B. Kang and C. Lee, *ACS Appl. Mater. Interfaces*, 2019, **11**, 12613–12621.
- 8 B. Li, T. Xing, M. Zhong, L. Huang, N. Lei, J. Zhang, J. Li and Z. Wei, *Nat. Commun.*, 2017, **8**, 1958.
- 9 S. Feng, Z. Lin, X. Gan, R. Lv and M. Terrones, *Nanoscale Horiz.*, 2017, **2**, 72–80.
- 10 G. M. Lohar, S. T. Jadhav, M. V. Takale, R. A. Patil, Y. R. Ma, M. C. Rath and V. J. Fulari, *J. Colloid Interface Sci.*, 2015, **458**, 136–146.
- 11 G. T. Phan, D. V. Pham, R. A. Patil, C.-C. Lai, W.-C. Yeh, Y. Liou and Y.-R. Ma, *Applied Materials Today*, 2019, **15**, 605–613.
- 12 Y. Lu and J. Guo, *Nano Res.*, 2010, **3**, 189–199.
- 13 Y. Li, T. Wang, M. Wu, T. Cao, Y. Chen, R. Sankar, R. K. Ulaganathan, F. Chou, C. Wetzel, C.-Y. Xu, S. G. Louie and S.-F. Shi, *2D Mater.*, 2018, **5**, 021002.
- 14 B. Deng, V. Tran, Y. Xie, H. Jiang, C. Li, Q. Guo, X. Wang, H. Tian, S. J. Koester, H. Wang, J. J. Cha, Q. Xia, L. Yang and F. Xia, *Nat. Commun.*, 2017, **8**, 14474.
- 15 Y. L. Huang, Y. Chen, W. Zhang, S. Y. Quek, C.-H. Chen, L.-J. Li, W.-T. Hsu, W.-H. Chang, Y. J. Zheng, W. Chen and A. T. S. Wee, *Nat. Commun.*, 2015, **6**, 6298.
- 16 D. V. Pham, R. A. Patil, J.-H. Lin, C.-C. Lai, Y. Liou and Y.-R. Ma, *Nanoscale*, 2016, **8**, 5559–5566.
- 17 B. S. Chang, A. Martin, B. Thomas, A. Li, R. W. Dorn, J. Gong, A. J. Rossini and M. M. Thuo, *ACS Mater. Lett.*, 2020, **2**, 1211–1217.
- 18 W. S. Choi, M. F. Chisholm, D. J. Singh, T. Choi, G. E. Jellison and H. N. Lee, *Nat. Commun.*, 2012, **3**, 689.
- 19 H. Zeuthen, W. Kudernatsch, L. R. Merte, L. K. Ono, L. Lammich, F. Besenbacher and S. Wendt, *ACS Nano*, 2015, **9**, 573–583.
- 20 R. Paul, R. G. Reifengerger, T. S. Fisher and D. Y. Zemlyanov, *Chem. Mater.*, 2015, **27**, 5915–5924.
- 21 J. Muro-Cruces, A. G. Roca, A. López-Ortega, E. Fantechi, D. del-Pozo-Bueno, S. Estradé, F. Peiró, B. Sepúlveda, F. Pineider, C. Sangregorio and J. Nogues, *ACS Nano*, 2019, **13**, 7716–7728.
- 22 J. Zhao, Y. Niu, B. Ren, H. Chen, S. Zhang, J. Jin and Y. Zhang, *Chem. Eng. J.*, 2018, **347**, 574–584.
- 23 R. Sinmyo, E. Bykova, S. V. Ovsyannikov, C. McCammon, I. Kuppenko, L. Ismailova and L. Dubrovinsky, *Sci. Rep.*, 2016, **6**, 32852.
- 24 A. G. Tamirat, J. Rick, A. A. Dubale, W.-N. Su and B.-J. Hwang, *Nanoscale Horiz.*, 2016, **1**, 243–267.
- 25 Y. Makimizu, J. Yoo, M. Poornajar, N. T. Nguyen, H.-J. Ahn, I. Hwang, S. Kment and P. Schmuki, *J. Mater. Chem. A*, 2020, **8**, 1315–1325.
- 26 J. D. Emery, C. M. Schlepütz, P. Guo, S. C. Riha, R. P. H. Chang and A. B. F. Martinson, *ACS Appl. Mater. Interfaces*, 2014, **6**, 21894–21900.
- 27 Y. Zhang, N. Zhang, T. Wang, H. Huang, Y. Chen, Z. Li and Z. Zou, *Appl. Catal., B*, 2019, **245**, 410–419.
- 28 K. Liang, L. S. Hui and A. Turak, *Nanoscale*, 2019, **11**, 9076–9084.



- 29 J. Huang, M. Guo, H. Ke, C. Zong, B. Ren, G. Liu, H. Shen, Y. Ma, X. Wang, H. Zhang, Z. Deng, H. Chen and Z. Zhang, *Adv. Mater.*, 2015, **27**, 5049–5056.
- 30 G. Segev, H. Dotan, K. D. Malviya, A. Kay, M. T. Mayer, M. Grätzel and A. Rothschild, *Adv. Energy Mater.*, 2016, **6**, 1500817.
- 31 L.-L. Tian, M.-J. Zhang, C. Wu, Y. Wei, J.-X. Zheng, L.-P. Lin, J. Lu, K. Amine, Q.-C. Zhuang and F. Pan, *ACS Appl. Mater. Interfaces*, 2015, **7**, 26284–26290.
- 32 J. Zhao, Z. Li, X. Yuan, Z. Yang, M. Zhang, A. Meng and Q. Li, *Adv. Energy Mater.*, 2018, **8**, 1702787.
- 33 Y. Li, J. Xu, T. Feng, Q. Yao, J. Xie and H. Xia, *Adv. Funct. Mater.*, 2017, **27**, 1606728.
- 34 T. Jiang, F. Bu, X. Feng, I. Shakir, G. Hao and Y. Xu, *ACS Nano*, 2017, **11**, 5140–5147.
- 35 X. Zhu, Y. Zhu, S. Murali, M. D. Stoller and R. S. Ruoff, *ACS Nano*, 2011, **5**, 3333–3338.
- 36 K. Cao, L. Jiao, H. Liu, Y. Liu, Y. Wang, Z. Guo and H. Yuan, *Adv. Energy Mater.*, 2015, **5**, 1401421.
- 37 J. C. Espinosa, C. Catalá, S. Navalón, B. Ferrer, M. Álvaro and H. García, *Appl. Catal., B*, 2018, **226**, 242–251.
- 38 S. H. Yoo, D. Jang, H.-I. Joh and S. Lee, *J. Mater. Chem. A*, 2017, **5**, 748–755.
- 39 K. Zhao, H. Tang, B. Qiao, L. Li and J. Wang, *ACS Catal.*, 2015, **5**, 3528–3539.
- 40 C. Yin, C. Gong, J. Chu, X. Wang, C. Yan, S. Qian, Y. Wang, G. Rao, H. Wang, Y. Liu, X. Wang, J. Wang, W. Hu, C. Li and J. Xiong, *Adv. Mater.*, 2020, **32**, 2002237.
- 41 L. Wan, H. Song, J. Ma, Y. Ren, X. Cheng, J. Su, Q. Yue and Y. Deng, *ACS Appl. Mater. Interfaces*, 2018, **10**, 13028–13039.
- 42 Z.-F. Dou, C.-Y. Cao, Q. Wang, J. Qu, Y. Yu and W.-G. Song, *ACS Appl. Mater. Interfaces*, 2012, **4**, 5698–5703.
- 43 W. Geng, S. Ge, X. He, S. Zhang, J. Gu, X. Lai, H. Wang and Q. Zhang, *ACS Appl. Mater. Interfaces*, 2018, **10**, 13702–13711.
- 44 D. V. Pham, R. A. Patil, C.-C. Yang, W.-C. Yeh, Y. Liou and Y.-R. Ma, *Nano Energy*, 2018, **47**, 105–114.
- 45 R. A. Patil, R. S. Devan, J.-H. Lin, Y.-R. Ma, P. S. Patil and Y. Liou, *Sol. Energy Mater. Sol. Cells*, 2013, **112**, 91–96.
- 46 J.-H. Lin, R. A. Patil, M.-A. Wu, L.-G. Yu, K.-D. Liu, W.-T. Gao, R. S. Devan, C.-H. Ho, Y. Liou and Y.-R. Ma, *J. Mater. Chem. C*, 2014, **2**, 8667–8672.
- 47 R. S. Devan, Y.-R. Ma, M. A. More, R. T. Khare, V. V. Antad, R. A. Patil, V. P. Thakare, R. S. Dhayal and L. Schmidt-Mende, *RSC Adv.*, 2016, **6**, 98722–98729.
- 48 G. T. Phan, D. V. Pham, R. A. Patil, C.-H. Tsai, C.-C. Lai, W.-C. Yeh, Y. Liou and Y.-R. Ma, *Sol. Energy Mater. Sol. Cells*, 2021, **231**, 111306.
- 49 W. Zhang, Y. Wu, J. Qi, M. Chen and R. Cao, *Adv. Energy Mater.*, 2017, **7**, 1602547.
- 50 Y. Wu, M. Chen, Y. Han, H. Luo, X. Su, M.-T. Zhang, X. Lin, J. Sun, L. Wang, L. Deng, W. Zhang and R. Cao, *Angew. Chem., Int. Ed.*, 2015, **54**, 4870–4875.
- 51 H. Wang, F. Yin, B. Chen and G. Li, *J. Mater. Chem. A*, 2015, **3**, 16168–16176.
- 52 T. Wang, B. Daiber, J. M. Frost, S. A. Mann, E. C. Garnett, A. Walsh and B. Ehrler, *Energy Environ. Sci.*, 2017, **10**, 509–515.
- 53 A. Loiudice, J. Ma, W. S. Drisdell, T. M. Mattox, J. K. Cooper, T. Thao, C. Giannini, J. Yano, L. W. Wang, I. D. Sharp and R. Buonsanti, *Adv. Mater.*, 2015, **27**, 6733–6740.
- 54 M. Bernechea, Y. Cao and G. Konstantatos, *J. Mater. Chem. A*, 2015, **3**, 20642–20648.
- 55 X. Xie, C. Mao, X. Liu, L. Tan, Z. Cui, X. Yang, S. Zhu, Z. Li, X. Yuan, Y. Zheng, K. W. K. Yeung, P. K. Chu and S. Wu, *ACS Cent. Sci.*, 2018, **4**, 724–738.
- 56 W. Dong, H. Liu, J. K. Behera, L. Lu, R. J. H. Ng, K. V. Sreekanth, X. Zhou, J. K. W. Yang and R. E. Simpson, *Adv. Funct. Mater.*, 2018, **29**, 1806181.
- 57 Y. Li, Q. Li, L. Cao, X. Cui, Y. Yang, P. Xiao and Y. Zhang, *Electrochim. Acta*, 2015, **178**, 171–178.
- 58 C.-H. Xu, P.-Y. Shen, Y.-F. Chiu, P.-W. Yeh, C.-C. Chen, L.-C. Chen, C.-C. Hsu, I. C. Cheng and J.-Z. Chen, *J. Alloys Compd.*, 2016, **676**, 469–473.

