



Cite this: *Phys. Chem. Chem. Phys.*, 2021, **23**, 17774

Systematic optimization of a fragment-based force field against experimental pure-liquid properties considering large compound families: application to oxygen and nitrogen compounds†

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The CombiFF approach is a workflow for the automated refinement of force-field parameters against experimental condensed-phase data, considering entire classes of organic molecules constructed using a fragment library *via* combinatorial isomer enumeration. One peculiarity of this approach is that it relies on an electronegativity-equalization scheme to account for induction effects within molecules, with values of the atomic hardness and electronegativity as electrostatic parameters, rather than the partial charges themselves. In a previous article [M. P. Oliveira, M. Andrey, S. R. Rieder, L. Kern, D. F. Hahn, S. Riniker, B. A. C. Horta and P. H. Hünenberger, *J. Chem. Theory. Comput.* 2020, **16**, 7525], CombiFF was introduced and applied to calibrate a GROMOS-compatible united-atom force field for the saturated acyclic (halo-)alkane family. Here, this scheme is employed for the construction of a corresponding force field for saturated acyclic compounds encompassing eight common chemical functional groups involving oxygen and/or nitrogen atoms, namely: ether, aldehyde, ketone, ester, alcohol, carboxylic acid, amine, and amide. Monofunctional as well as homo-polyfunctional compounds are considered. A total of 1712 experimental liquid densities ρ_{liq} and vaporization enthalpies ΔH_{vap} concerning 1175 molecules are used for the calibration (339 molecules) and validation (836 molecules) of the 102 non-bonded interaction parameters of the force field. Using initial parameter values based on the GROMOS 2016H66 parameter set, convergence is reached after five iterations. Given access to one processor per simulated system, this operation only requires a few days of wall-clock computing time. After optimization, the root-mean-square deviations from experiment are 29.9 (22.4) kg m⁻³ for ρ_{liq} and 4.1 (5.5) kJ mol⁻¹ for ΔH_{vap} for the calibration (validation) set. Thus, a very good level of agreement with experiment is achieved in terms of these two properties, although the errors are inhomogeneously distributed across the different chemical functional groups.

Received 6th May 2021,
 Accepted 30th June 2021

DOI: 10.1039/d1cp02001c

rsc.li/pccp

1 Introduction

The ability of classical molecular dynamics (MD) simulation to represent accurately the properties of a given system depends

crucially on the quality of the underlying potential-energy function or force field.^{1–7} One may tentatively distinguish three main strategies in the design of condensed-phase force fields:⁸

(1) In fragment-based force fields^{9–13} (FBFF), the covalent and non-bonded parameters are specified within molecular fragments representative of the relevant chemical functional groups. Molecular topologies are built by assembling these fragments and invoking an assumption of transferability. The non-bonded parameters are calibrated primarily by fitting against experimental thermodynamic data for small organic compounds in the condensed phase.

(2) In hybrid force fields^{14,15} (HYFF), the covalent and van der Waals parameters are selected in a FBFF fashion, but the partial charges are derived based on quantum-mechanics (QM) calculations involving the target molecule, typically *via* electrostatic-potential fitting^{14–28} or electron-density partitioning.^{29–38} The assumption invoked in this case is that of a separability between van der

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† Electronic supplementary information (ESI) available: Includes detailed information concerning: (i) the compounds in the calibration and validation sets of the O + N family; (ii) the experimental data; (iii) the covalent parameters; (iv) the initial values of the non-bonded parameters; (v) the observable-to-parameter ratio for the different atom types; (vi) the comparison with experiment; (vii) the list of outliers; (viii) the comparison with results using the 2016H66 parameter set. Additional material (molecule identifiers, coordinate and topology files, data tables for experimental and simulated values) can be downloaded freely from the internet under ref. 125, where version 1.0 corresponds to the published article (further versions will include revisions and/or expansions of the data set). See DOI: 10.1039/d1cp02001c



Waals coefficients and partial charges. Given a selected charge-derivation scheme, the van der Waals parameters are again calibrated primarily against experimental condensed-phase data for small compounds.

(3) In QM-derived force fields^{39–46} (QDFF), the covalent and non-bonded parameters are determined simultaneously based on QM calculations involving the target molecule. The calculation schemes for the partial charges are the same as in the HYFF case. For the van der Waals coefficients, the derivation typically involves electron-density partitioning as a starting point.^{47–57} The assumption invoked here is that of a compatibility between the non-bonded interaction parameters appropriate for the isolated molecule and for the molecule in the condensed phase.

The HYFF and QDFF schemes are very popular nowadays, in particular because they: (i) benefit from fast QM calculation methods;^{58–65} (ii) promise an exhaustive coverage of the chemical space;^{66,67} (iii) take into account induction effects on the partial charges^{24,33,68} and, possibly, on the van der Waals coefficients;^{44,45,51–53,56,57} (iv) are comparatively easy to automate in terms of topology construction and parameter derivation. However, compared to the FBFF approach, they also present some major shortcomings: (i) the need to specify a reference structure (molecular conformation) and environment (*e.g.* vacuum or continuum solvent) in the QM calculations leads to an implicit dependence of the topological information on configurational information; (ii) the partial charges and van der Waals coefficients are not strictly speaking QM observables, *i.e.* they result from *ad hoc* derivation recipes rather than physics-based rules, and their values may also strongly depend on the choice of a QM level of theory and basis set;^{21–25,69–75} (iii) in practice, some parameters must still be optimized empirically,^{21,26,27,40,46,74,75} *e.g.* van der Waals coefficients in HYFF and van der Waals repulsion coefficients^{38,44,45,56,57,76,77} in QDFF.

In contrast to the HYFF and QDFF approaches, the FBFF scheme fully acknowledges that the non-bonded parameters are truly empirical quantities. They compensate in a mean-field fashion for all sorts of deficiencies in the selected potential-energy function, and they are correlated with a number of associated choices,⁷⁸ including those of the model resolution (*e.g.* united- vs. all-atom), van der Waals combination rules, cutoff distances, and treatment of the long-range interactions. As a result, the connection between these empirical non-bonded parameters and QM-inferred single-molecule properties may in fact be rather weak. However, for a FBFF approach to be useful in practice, it must achieve: (i) a sufficiently broad (even if not exhaustive) coverage of the chemical space; (ii) an appropriate representation of induction effects; (iii) a high degree of automation in the topology construction and parameter optimization. In a recent article,⁸ we introduced a scheme called CombiFF that presents these features.

The goal of CombiFF is the automated refinement of force-field parameters against experimental condensed-phase data, considering entire classes of organic molecules constructed using a fragment library *via* combinatorial isomer

enumeration. The main steps of the scheme are: (i) definition of a molecule family; (ii) combinatorial enumeration of all isomers; (iii) query for experimental data; (iv) automatic construction of the molecular topologies by fragment assembly; (v) iterative refinement of the force-field parameters considering the entire family. This scheme borrows from earlier work on isomer enumeration^{79–82} and topology construction,^{83–95} as well as on automated single-compound force-field optimization approaches such as the POP scheme,^{96,97} the ForceBalance scheme,^{46,98–105} and other related schemes.^{106–111} One key feature of CombiFF is that once the time-consuming task of target-data selection/curation has been performed, the optimization of a force field is entirely automatic and, given access to a sufficient number of processors, only requires a few days of wall-clock computing time. For this reason, CombiFF also represents an ideal framework for assessing the impact of specific functional-form decisions on the accuracy of a force field at an optimal level of parametrization.

As a first application, CombiFF was used in ref. 8 to design a GROMOS-compatible united-atom force field for the saturated acyclic (halo)-alkane family. A total of 749 experimental liquid densities ρ_{liq} and vaporization enthalpies ΔH_{vap} concerning 486 haloalkane molecules were considered for the calibration (228 molecules) and validation (258 molecules) of 32 non-bonded interaction parameters.

An important aspect of this force field is that the atomic partial charges are not specified explicitly within the fragments,¹¹² but determined implicitly using an electronegativity-equalization (EE) scheme,^{8,113} which permits to account for electronic induction effects within molecules. The corresponding atomic parameters, *i.e.* hardness and electronegativity, are expected to factor out these induction effects and to be much less dependent than the partial charges themselves on the covalent environment of the atoms. Note that although the EE scheme serves to generate the atomic partial charges, its function is fundamentally distinct from the electrostatic-potential fitting^{14–28} or electron-density partitioning^{29–38} procedures used in HYFF and QDFF force fields. Whereas the latter procedures aim at deriving partial charges based on a QM calculation, the EE scheme is only used here as a physically-motivated parametric-fitting device for the electrostatic interactions in the liquid, in which the partial charges are solely determined by the condensed-phase properties (no QM input).

For the (halo)-alkane force field,⁸ the parameter calibration resulted in root-mean-square deviations from experiment of 49.8 (27.6) kg m⁻³ for ρ_{liq} and 2.7 (1.8) kJ mol⁻¹ for ΔH_{vap} considering the calibration (validation) set. The values are lower for the validation set, because it contains larger molecules (stronger influence of purely aliphatic interactions). The trends in the optimized parameters along the halogen series and across the compound family were found to be in line with chemical intuition based on considerations related to size, polarizability, softness, electronegativity, induction, and hyperconjugation. This is remarkable considering that the force-field calibration did not involve any QM calculation.



The goal of the present study is to extend the application of CombiFF to the calibration of a GROMOS-compatible united-atom force field for saturated acyclic compounds encompassing eight common chemical functional groups involving oxygen and/or nitrogen atoms, namely ether, aldehyde, ketone, ester, alcohol, carboxylic acid, amine, and amide. Considering compounds of up to ten carbon atoms including up to four occurrences of the same functional group, the corresponding family of molecules is referred to in this article as the O + N family.

The GROMOS-compatible 2016H66 force field¹¹² is used as a starting point for the optimization and as a reference for performance comparison. Considering 62 small organic molecules, this parameter set was calibrated and validated against condensed-phase experimental data not only for ρ_{liq} and ΔH_{vap} , but also for the hydration free energy ΔG_{wat} and the solvation free energy ΔG_{che} in cyclohexane. Note that in addition to 43 compounds of the O + N family, 2016H66 also included 19 molecules representative for thiols, sulfides, disulfides, aromatic compounds, and nucleic-acid bases, which are not considered here.

The present force-field reoptimization using CombiFF relies on a much higher observable-to-parameter ratio. Here, a total of 1712 experimental values for ρ_{liq} and ΔH_{vap} concerning 1175 representative molecules of the O + N family are extracted from nine data sources,^{114–122} and used as target for the calibration and validation of 102 non-bonded interaction parameters.

2 Methodology

The CombiFF workflow for calibrating the parameters of a force field based on experimental data concerning a given compound family is illustrated in Fig. 1. For a detailed description, the reader is referred to our previous article,⁸ where the scheme was introduced (see in particular Appendix A therein). The present section only summarizes its main features, and provides information on its application to the O + N family. For the ease of reference, a few key numbers (symbols and values) relevant to this optimization are summarized in Table 1.

The O + N family of compounds is defined as the union of the 27 subfamilies listed in Table 2. It includes molecules of up to ten carbon atoms representative for eight chemical functional groups: ether, aldehyde, ketone, ester, alcohol, carboxylic acid, amine, and amide. Depending on the subfamily, the given functional group may occur up to four times in the compound. However, molecules combining different types of functional groups are not considered here. The $N_{\text{iso}}^{\text{ot}} = 57\,905$ constitutional isomers of the O + N family were enumerated as canonical SMILES strings^{80,123} using an in-house program (ENU). The numbers N_{iso} of isomers in each of the subfamilies are reported in Table 2.

To collect available experimental data, the SMILES strings must be mapped to various equivalent identifiers that can be found in the different experimental sources. For example, the SMILES string CC(O)C may appear under the alternative identifiers 2-propanol, propan-2-ol, isopropyl alcohol, isopropanol, or CAS 67-63-0. The PubChem database¹²⁴ was used to obtain such alternative identifiers. About one quarter of the isomers

(15 592 molecules) were found in PubChem, and about one fifth of these (3425 molecules) were associated with a CAS registry number. The experimental database (DBS) maintained in our group was queried for ρ_{liq} and ΔH_{vap} values pertaining to all these compounds, giving priority to a match by CAS (when available) over a match by name. The nine data sources accessed were ref. 114–122. Note that the data points from ref. 121 that are marked as “estimated” were discarded. The ΔH_{vap} values from this source for the alcohols and carboxylic acids were excluded as well due to inconsistencies, unless a similar value was found in another source. This resulted in ρ_{liq} and/or ΔH_{vap} values concerning $N_{\text{iso}}^{\text{sim}} = 1175$ compounds, which were distributed into a calibration set of $N_{\text{iso}}^{\text{cal}} = 339$ molecules and a validation set of $N_{\text{iso}}^{\text{val}} = 836$ molecules based on the number of carbon atoms (1–6 vs. 7–10). The structures of these compounds are shown in ESI,† Section S1 (Fig. S1 and S2). The acronyms employed for the individual molecules involve one letter and four digits. The letter is representative of the chemical function (see Table 2). The first digit stands for the number of carbon atoms, with the number ten mapped to the digit zero. Finally, the last three digits form a sequential index, further distinguishing compounds for which the first two symbols are identical.

The GROMOS-compatible molecular topologies of these compounds were generated automatically based on the SMILES strings using an in-house program (TBL), by linking the fragments shown in Fig. 2 *via* bond overlap. A total of 54 fragments are required to construct all the molecules of the O + N family (as well as the saturated acyclic alkanes beyond methane). Note, however, that these fragments are only sufficient to generate molecules containing one or more occurrence of a single type of functional group.

The experimental-data vector \mathbf{X}^{exp} corresponding to the calibration set has the dimension $N_{\text{exp}}^{\text{cal}} = 579$. It encompasses $N_{\rho}^{\text{cal}} = 314$ values for ρ_{liq} and $N_{\Delta H}^{\text{cal}} = 265$ values for ΔH_{vap} , and requires $N_{\text{sim}}^{\text{cal}} = 408$ independent simulations (*i.e.* distinct compounds and P, T -points) for its evaluation. The corresponding vector for the validation set has the dimension $N_{\rho}^{\text{val}} = 1133$, encompasses $N_{\rho}^{\text{val}} = 765$ values for ρ_{liq} and $N_{\Delta H}^{\text{val}} = 368$ values for ΔH_{vap} , and requires $N_{\text{sim}}^{\text{val}} = 997$ independent simulations for its evaluation. The reference experimental values retained for ρ_{liq} and/or ΔH_{vap} , along with the associated P, T -points, are listed in ESI,† Section S2 (Table S1). This material can be downloaded freely from the internet under ref. 125, where the version 1.0 corresponds to the published article (further versions will include revisions and/or expansions of the data set).

Most of the state points considered are within 10 K of the standard temperature $T^- = 298.15$ K (70% of the points) and within 0.2 bar of the atmospheric pressure $P^0 = 1$ bar (78% of the points). The values for the remaining points range from 250 to 537 K and from 0.002 to 9.06 bar. The impact of using a (limited) fraction of target values at temperatures differing from T^- as well as the ability of the force field to reproduce experimental data at temperatures differing (reasonably) from T^- have already been investigated for the (halo-)alkane family⁸ (see Section S13 in the ESI,† of this previous article). There, it was shown that: (i) at the calibration level, there is no systematic



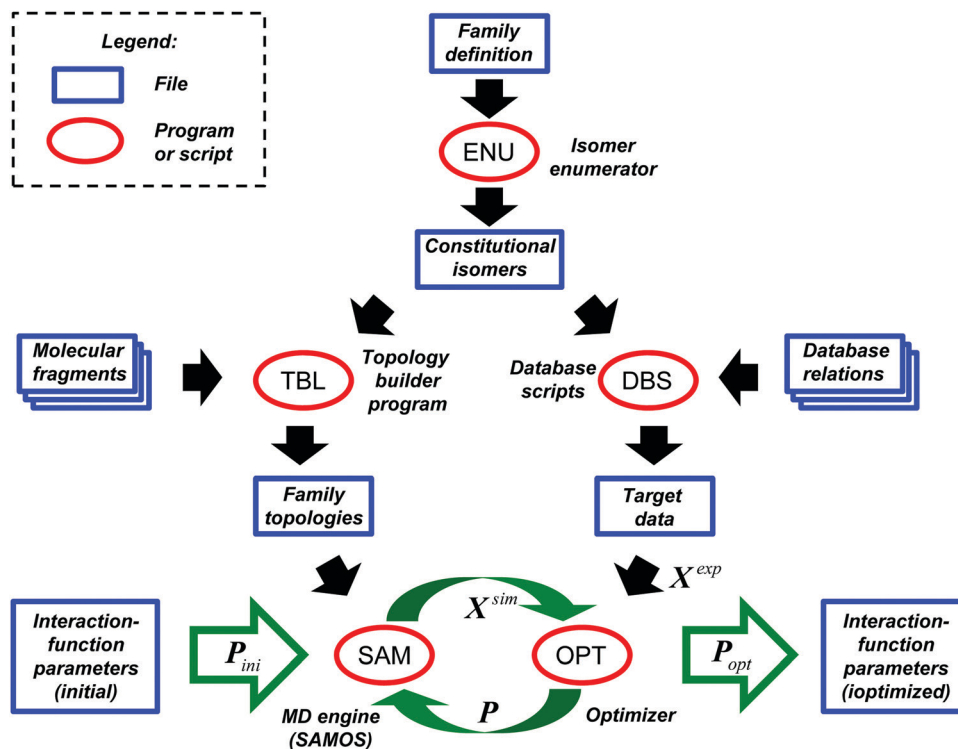


Fig. 1 CombiFF workflow for calibrating the parameters of a force field based on experimental data concerning a given compound family. Based on the definition of the family, the program ENU enumerates all possible constitutional isomers. The program TBL is then used to construct the molecular topologies of the corresponding compounds, and the DBS scripts to extract available target experimental data pertaining to these molecules from an in-house database. By alternating simulations (MD engine SAM) to calculate the vector X^{sim} of simulated observables (as well as its derivative relative to all force-field parameters), and variations of the parameter vector P (optimization script OPT) designed to bring these simulated observables closer to their experimental target vector X^{exp} , the initial parameters P_{ini} are progressively refined into optimal parameters P_{opt} .

correlation between the temperature and the error (calculation vs. experiment) for ρ_{liq} and ΔH_{vap} ; (ii) at the validation level (considering 50 experimental curves for T -dependent properties spanning a range of about 350 K), there is no correlation between the temperature and the error for ρ_{liq} , and only a weak positive correlation for ΔH_{vap} .

The principles of the force-field representation employed here are compatible with those of the GROMOS force field^{126–131} in its 2016H66 variant,¹¹² except for one important difference. The atomic partial charges are determined for each molecule based on an EE scheme.¹¹³ Similarly to our previous work⁸ (see Appendix A.4 therein), charge flows between atoms are only allowed within overall neutral charge groups, and intramolecular Coulombic effects (J -terms in the EE scheme) are only included for first and second covalent neighbors, as the corresponding interatomic distances can be considered (essentially) configuration-independent. The corresponding terms for more remote covalent neighbors are not necessarily negligible, but their inclusion would lead to conformation-dependent charges and require the application of an expensive on-the-fly EE scheme during the simulations. Since the EE scheme is only a parametric-fitting device for the charges, its effective parameters are expected compensate at least in part for possible deficiencies in the representation. The validity of this assumption is ultimately supported by the observation that an excellent fit to the experimental data can be obtained.

The charge groups relevant for the molecules of the O + N family are illustrated in Fig. 3. All the aliphatic (united-)atoms of the molecule (atom types CH0, CH1, CH2 and CH3 in Table 3) that are not explicitly included in one of these charge groups define separate one-particle charge groups with a charge of zero. The charge groups are used in GROMOS for the application of the non-bonded interaction cutoff, which performs a group-based truncation in terms of the centers of geometry of the two charge groups. The Gaussian-cloud interaction accounting for intramolecular Coulombic effects in the EE scheme (J -terms) between first and second covalent neighbors within a charge group relies on effective interatomic distances \bar{r} calculated based on the reference bond lengths and angles of the covalent force-field parameters, along with effective radii that are set to the Lennard-Jones collision diameters σ of the involved (united-)atoms.

The covalent interaction parameters relevant for the O + N family were taken or ported by analogy from the 2016H66 parameter set,¹¹² and kept unaltered. The corresponding information is summarized in ESI,† Section S3 (Table S2 and Fig. S3). Only the non-bonded interaction parameters were subjected to refinement, and solely a subset thereof.

The atomic partial charges are determined indirectly *via* the EE scheme based on two types of atomic parameters, the hardness η and the electronegativity χ . Owing to the use of a geometric-mean combination rule^{132,133} for the Lennard-Jones (LJ) interactions,¹³⁴ the corresponding pairwise coefficients



Table 1 Key numbers (symbols and values) pertaining to the CombiFF force-field calibration and validation applied to the O + N family of compounds. This family is defined as the union of the 27 subfamilies listed in Table 2. The structures of the $N_{\text{iso}}^{\text{sim}}$ representative molecules considered in the simulations are shown in ESI,† Section S1, separately for the calibration set ($N_{\text{iso}}^{\text{cal}}$ molecules with 1–6 carbon atoms, Fig. S1, ESI†) and the validation set ($N_{\text{iso}}^{\text{val}}$ molecules with 7–10 carbon atoms, Fig. S2, ESI†). The experimental data pertaining to the $N_{\text{sim}}^{\text{tot}}$ combinations of compounds and P, T -points can be found in ESI,† Section S2 (Table S1). The parameters associated with the N_{att} atom types (equivalent to EE-types) and the $N_{\text{att}}^{\text{LJ}}$ LJ-types are reported in Tables 3 and 4, respectively. The information concerning the $N_{\text{prm}}^{\text{cov}}$ covalent parameters is summarized in ESI,† Section S3 (Table S2 and Fig. S3). Note that number $N_{\text{prm}}^{\text{cal}}$ of parameters optimized is smaller than the total number $N_{\text{prm}}^{\text{tot}}$ of force-field parameters because only non-bonded parameters are optimized, and solely a subset thereof

Number	Value	Meaning
$N_{\text{iso}}^{\text{tot}}$	57 905	Total number of constitutional isomers in the O + N family
$N_{\text{iso}}^{\text{sim}}$	1175	Total number of isomers with available experimental data ($= N_{\text{iso}}^{\text{cal}} + N_{\text{iso}}^{\text{val}}$)
$N_{\text{exp}}^{\text{tot}}$	1712	Total number of experimental data points ($= N_{\text{exp}}^{\text{cal}} + N_{\text{exp}}^{\text{val}}$)
$N_{\text{sim}}^{\text{tot}}$	1405	Total number of distinct P, T -points in this data ($= N_{\text{sim}}^{\text{cal}} + N_{\text{sim}}^{\text{val}}$)
$N_{\text{iso}}^{\text{cal}}$	339	Compounds included in the calibration set
$N_{\text{exp}}^{\text{cal}}$	579	Experimental data points for the calibration set ($= N_{\rho}^{\text{cal}} + N_{\Delta H}^{\text{cal}}$)
N_{ρ}^{cal}	314	Experimental ρ_{liq} data points for the calibration set
$N_{\Delta H}^{\text{cal}}$	265	Experimental ΔH_{vap} data points for the calibration set
$N_{\text{sim}}^{\text{cal}}$	408	Distinct compounds and P, T -points (<i>i.e.</i> simulations) for the calibration set
$N_{\text{iso}}^{\text{val}}$	836	Compounds included in the validation set
$N_{\text{exp}}^{\text{val}}$	1133	Experimental data points for the validation set ($= N_{\rho}^{\text{val}} + N_{\Delta H}^{\text{val}}$)
N_{ρ}^{val}	765	Experimental ρ_{liq} data points for the validation set
$N_{\Delta H}^{\text{val}}$	368	Experimental ΔH_{vap} data points for the validation set
$N_{\text{sim}}^{\text{val}}$	997	Distinct compounds and P, T -points (<i>i.e.</i> simulations) for the validation set
$N_{\text{att}}^{\text{EE}}, N_{\text{att}}$	47	Number of EE-types (or, equivalently, atom types)
$N_{\text{att}}^{\text{LJ}}$	12	Number of LJ-types
$N_{\text{prm}}^{\text{tot}}$	233	Total number of force-field parameters ($= N_{\text{prm}}^{\text{cov}} + N_{\text{prm}}^{\text{att}}$)
$N_{\text{prm}}^{\text{cov}}$	94	Number of covalent parameters
$N_{\text{prm}}^{\text{mbd}}$	139	Number of non-bonded parameters
$N_{\text{prm}}^{\text{cal}}$	102	Number of parameters that are optimized

are also constructed based on two types of atomic parameters, the collision diameter σ and the well depth ε . Following the GROMOS design principle, the values σ and ε are only used in the combination rule for non-hydrogen-bonding LJ-type pairs (corresponding to the LJ parameters C_6 and $C_{12,1}$ in GROMOS). For hydrogen-bonding LJ-type pairs, GROMOS relies on a modified set of LJ parameters with slightly enhanced repulsion. In this case, alternative values $\bar{\sigma}$ and $\bar{\varepsilon}$ are used instead (corresponding to the LJ parameters C_6 and $C_{12,II}$ in GROMOS). For simplicity, the value of the dispersion coefficient C_6 is kept identical in the two sets. As result, only $\bar{\sigma}$ needs to be specified, while $\bar{\varepsilon}$ can be deduced as $\bar{\varepsilon} = \varepsilon \sigma^6 / \bar{\sigma}^6$. Finally, for third covalent neighbors, yet another pair of values σ^* and ε^* is used in the combination rule. Each atom type of the force field is thus associated with a unique selection for six (non-hydrogen-bonding type) or seven (potentially hydrogen-bonding type) parameters. However, the same σ and ε parameters are often used for different atom types of the same element. As a result, the present force field for the O + N family relies on a number $N_{\text{att}} = 47$ of atom types, which are equivalent to EE-types ($N_{\text{att}}^{\text{EE}} = N_{\text{att}}$), but involves a smaller number $N_{\text{att}}^{\text{LJ}} = 12$ of LJ-types. The final (optimized) values of the EE parameters for the 47 atom types (or EE-types) are reported in Table 3, along with a LJ-type. The latter refers to the entries of Table 4, where the final values of the LJ parameters are reported for the 12 LJ-types. The correspondence between elements, LJ-types, atom-types (EE-types), and

chemical functional groups involving the latter atom types is also illustrated schematically in Fig. 4.

The four aliphatic atom types (CH0, CH1, CH2 and CH3) have no EE parameters, as their charge is always zero. The LJ parameters of these atom types,^{135–137} along with those of the polar hydrogen atom (type HB, zero in GROMOS), as well as all the third-neighbor LJ interaction parameters, were also excluded from the optimization, *i.e.* kept as in the 2016H66 set.¹¹² Note also that, in the absence of parametrization target, the η and χ values of EE-type CH0_N_amd are estimated (no experimental data found for a compound involving this type), and that the $\bar{\sigma}$ value of the LJ-type OR is kept identical to σ (no hydrogen bonding is possible in molecules containing only ether groups). The initial parameter values selected to start the optimization are reported in ESI,† Section S4 (Tables S3 and S4). For the LJ parameters, they were ported from the 2016H66 force field.¹¹² For the EE parameters, they were fitted to best reproduce the atomic partial charges of this force field.

Following from the above choices, the present force field for the O + N family involves $N_{\text{prm}}^{\text{tot}} = 233$ parameters, namely $N_{\text{prm}}^{\text{cov}} = 94$ covalent parameters along with $N_{\text{prm}}^{\text{mbd}} = 139$ non-bonded parameters (2×43 relevant EE-types + 4×7 non-hydrogen-bonding LJ-types + 5×5 potentially hydrogen-bonding LJ-types), among which $N_{\text{prm}}^{\text{cal}} = 102$ are subject to optimization (omitted are 2 EE parameters for CH0_N_amd, 2×12 third-neighbor LJ parameters, 2×5 LJ parameters for CH0 to CH3



Table 2 Compounds of the O + N family. The family is defined as the union of 27 non-overlapping subfamilies of compounds, representative for eight chemical functional groups. The 14 acronyms retained for the different subfamilies (or small groups thereof) are further used in the text, tables, and figures. The one-character variant (Char.; only 11 groups) is used as a first letter in the acronyms of the corresponding molecules. For each subfamily, n stands for the number of carbon atoms, m for the number of occurrences of the functional group in the molecule, N_{iso} for the total number of constitutional isomers, and N_{sim} for the number of these isomers considered in the simulations (*i.e.* for which experimental data could be found). Note that, for simplicity, formaldehyde HCOH and formic acid HCOOH are included in RCOH and RCOOH, respectively. Note also that, the two amino groups in RN_2^* can be of different types (*i.e.* primary, secondary, or tertiary). The structures of the $N_{\text{iso}}^{\text{sim}}$ representative molecules considered in the simulations are shown in ESI,† Section S1, separately for the calibration set ($N_{\text{iso}}^{\text{cal}} = 339$ molecules with 1–6 carbon atoms, Fig. S1, ESI†) and the validation set ($N_{\text{iso}}^{\text{val}} = 836$ molecules with 7–10 carbon atoms, Fig. S2, ESI†)

Function/acronym	Char.	n	m	N_{iso}	N_{sim}	Subfamily description
Ethers						
ROR	O	1–10	1	817	85	C ₁ –C ₁₀ mono-ethers
	O	1–10	2	2544	46	C ₁ –C ₁₀ di-ethers (including acetals and ketals)
	O	1–10	3	4936	17	C ₁ –C ₁₀ tri-ethers (including acetals and ketals)
	O	1–10	4	6614	6	C ₁ –C ₁₀ tetra-ethers (including acetals and ketals)
Aldehydes						
RCOH	A	1	1	1	1	Formaldehyde
	A	2–10	1	372	35	C ₂ –C ₁₀ mono-aldehydes
	A	2–10	2	551	4	C ₂ –C ₁₀ di-aldehydes
Ketones						
RCOR	K	1–10	1	335	85	C ₁ –C ₁₀ mono-ketones
	K	1–10	2	463	18	C ₁ –C ₁₀ di-ketones
	K	1–10	3	379	2	C ₁ –C ₁₀ tri-ketones
Esters						
HCOOR	F	1–10	1	372	16	C ₁ –C ₁₀ mono-esters (only formates)
HCOOR	F	1–10	2	550	2	C ₁ –C ₁₀ di-esters (only formates)
RCOOR	E	1–10	1	662	146	C ₁ –C ₁₀ mono-esters (without formates)
RCOOR	E	1–10	2	1364	49	C ₁ –C ₁₀ di-esters (without formates)
Alcohols						
ROH	L	1–10	1	879	280	C ₁ –C ₁₀ mono-ols
	L	1–10	2	3670	101	C ₁ –C ₁₀ di-ols
	L	1–10	3	11 249	6	C ₁ –C ₁₀ tri-ols
Carboxylic acids						
RCOOH	C	1	1	1	1	Formic acid
	C	2–10	1	372	51	C ₂ –C ₁₀ mono-carboxylic acids
	C	3–10	2	550	6	C ₃ –C ₁₀ di-carboxylic acids
Amines						
RNH ₂	M	1–10	1	879	52	C ₁ –C ₁₀ mono-primary-amines
RNHR	N	1–10	1	817	42	C ₁ –C ₁₀ mono-secondary-amines
RNR ₂	R	1–10	1	420	43	C ₁ –C ₁₀ mono-tertiary-amines
RN ₂ *	N	1–10	2	17 665	51	C ₁ –C ₁₀ di-amines
Amides						
RCONH ₂	D	1–10	1	372	9	C ₁ –C ₁₀ mono-primary-amides
RCONHR	D	1–10	1	662	7	C ₁ –C ₁₀ mono-secondary-amides
RCONR ₂	D	1–10	1	409	14	C ₁ –C ₁₀ mono-tertiary-amides
Total	—	—	—	57 905	1175	Total over the 27 subfamilies

and HB, and 1 LJ parameter for the potentially hydrogen-bonding OR). The optimization of these parameters against $N_{\text{exp}}^{\text{cal}} = 579$ experimental data points in the calibration set involves an observable-to-parameter ratio of 5.7 (up to 16.8 when also considering the $N_{\text{exp}}^{\text{val}} = 1133$ data points in the validation set). This ratio is further analyzed for each of the EE- and LJ-types separately in ESI,† Section S5 (Tables S5 and S6). A favorable observable-to-parameter ratio is observed in most cases, although three EE-types (CH3_O_ol, H_CO_acd, and CH1_N_amd) only occur in a single representative molecule, and one (CH0_N_amd) is not represented at all.

The search for optimal parameters was performed as in our previous work⁸ (see Appendix A.7 therein), by minimizing an objective function $Q(\mathbf{P}; \mathbf{X}^{\text{exp}})$ of the parameter vector \mathbf{P} accounting for the deviation of the simulated-data vector $\mathbf{X}^{\text{sim}}(\mathbf{P})$ relative to the experimental-data vector \mathbf{X}^{exp} . This function is defined as

$$Q(\mathbf{P}; \mathbf{X}^{\text{exp}}) = W^{-1} \sum_{n=1}^{N_n} s_n^{-1} \sum_{m=1}^{N_m} w_{nm} |X_{nm}^{\text{sim}}(\mathbf{P}) - X_{nm}^{\text{exp}}| \quad (1)$$

with $W = \sum_{n=1}^{N_n} \sum_{m=1}^{N_m} w_{nm}$,



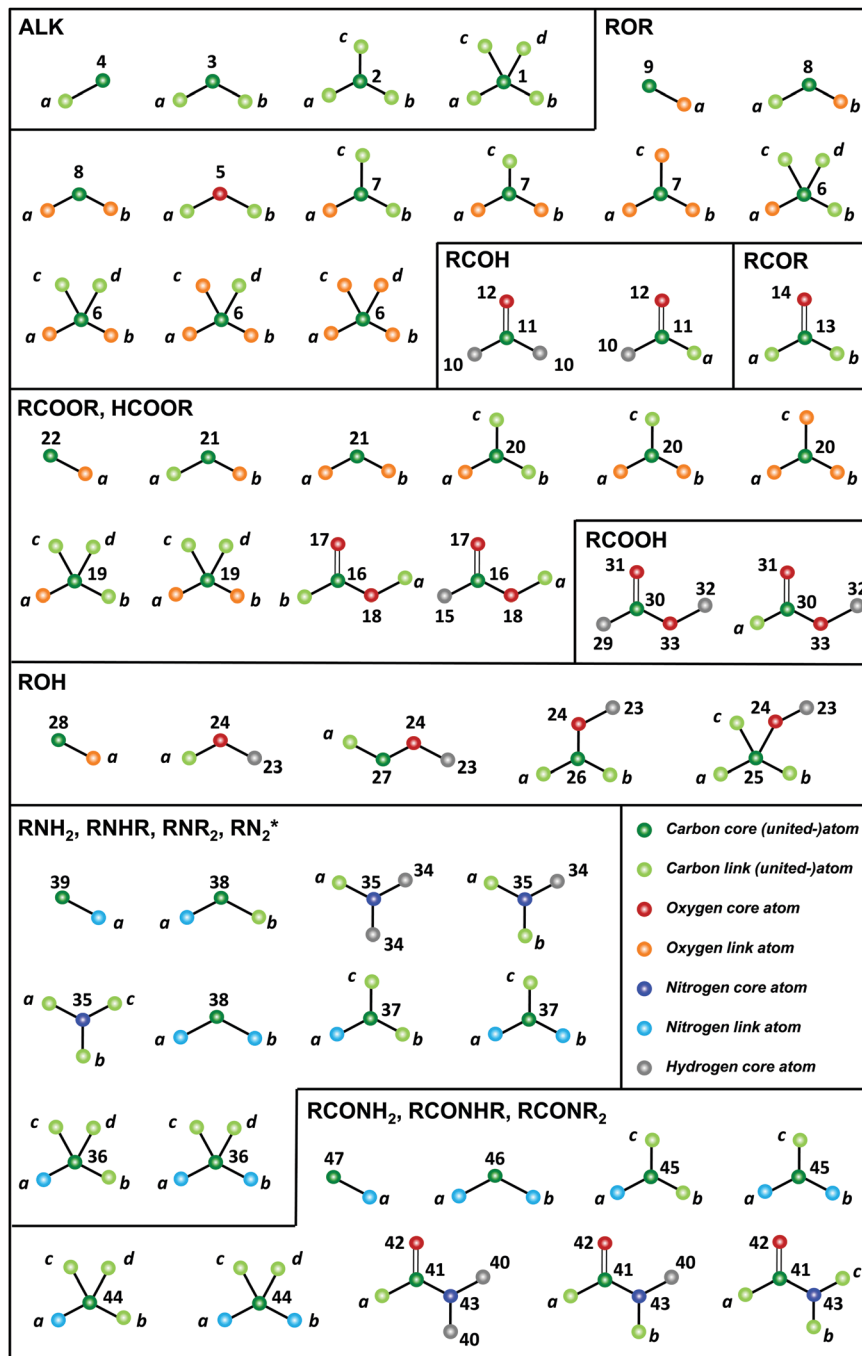


Fig. 2 Molecular-topology fragments required for the representation of the O + N family. The 54 fragments represented are required for the construction of the saturated acyclic alkanes beyond methane (ALK) plus the molecules of the O + N family. Note that these fragments are only sufficient to generate molecules containing one or more occurrence of a single type of functional group. The acronyms used for the different molecule groups are explicated in Table 2. Note that, for simplicity, formaldehyde HCOH and formic acid HCOOH are included in RCOH and RCOOH, respectively. The atoms of a fragment that can be linked (link atoms) are labeled with lower-case letters. The other atoms within the fragment (core atoms) are labeled with their atom types, referring to the numbering of Table 3. A linkage between two fragments is performed by overlapping a core-link bond of the two fragments.

where the index n corresponds to the N_n observable types and the index m to the N_m molecules in the family. The s_n coefficients eliminate the dependence on a unit system and adjust the relative weights of different observables in terms of perceived (*i.e.* subjective) extent of “badness”. They are set here to

20 kg m^{-3} for the ρ_{liq} observables and 1 kJ mol^{-1} for the ΔH_{vap} observables. The unitless coefficients w_{nm} can be used to weigh differently the contributions of specific observable/molecule combinations. For simplicity, all the combinations included (also considering observables at multiple state points) are



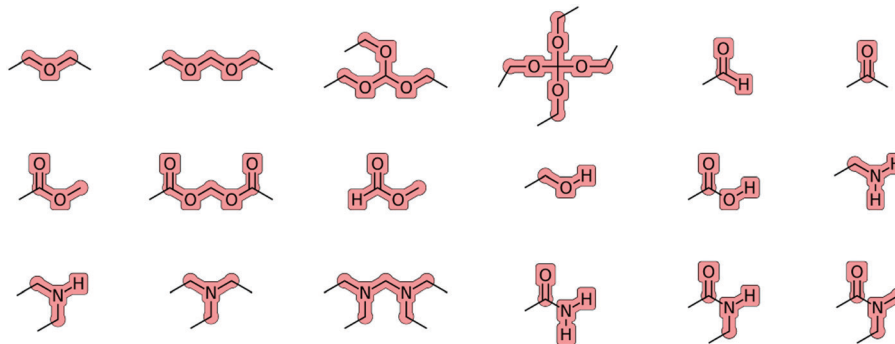


Fig. 3 Charge groups relevant for the compounds of the O + N family. Charge flows in the EE scheme are only permitted between atoms belonging to the same overall neutral charge group. All the aliphatic (united-)atoms of the molecule (atom types CH0, CH1, CH2 and CH3 in Table 3) that are not explicitly included in one of these charge groups define separate one-particle charge groups with a charge of zero.

associated here with the same weight $w_{nm} = 1$, while all the omitted ones (absence of experimental data) have $w_{nm} = 0$. The prefactor W ensures that the overall magnitude of Q is not affected by the number of data points included.

The optimization is performed by assuming that $\mathbf{X}^{\text{sim}}(\mathbf{P})$ is approximately linear in parameter changes within a small trust region around a reference point \mathbf{P}^0 in parameter space, *i.e.* using the local first-order approximation $\tilde{Q}(\mathbf{P}; \mathbf{P}^0, \mathbf{X}^{\text{exp}})$ to $Q(\mathbf{P}; \mathbf{X}^{\text{exp}})$ defined by

$$\tilde{Q}(\mathbf{P}; \mathbf{P}^0, \mathbf{X}^{\text{exp}}) = W^{-1} \sum_{n=1}^{N_n} s_n^{-1} \sum_{m=1}^{N_m} w_{nm} |X_{nm}^{\text{sim}}(\mathbf{P}^0) + S_{nm}(\mathbf{P}^0) \cdot (\mathbf{P} - \mathbf{P}^0) - X_{nm}^{\text{exp}}|, \quad (2)$$

where $\mathbf{S}(\mathbf{P}^0)$ is the sensitivity matrix defined by the variations of the different molecule/observable combinations with respect to variations of the N_k parameters around the point \mathbf{P}^0 , *i.e.*

$$S_{nm,k}(\mathbf{P}^0) = \left(\frac{\partial X_{nm}^{\text{sim}}(\mathbf{P})}{\partial P_k} \right)_{\mathbf{P}=\mathbf{P}^0}. \quad (3)$$

This matrix is calculated next to the observables themselves during the MD simulations at \mathbf{P}^0 using appropriate statistical-mechanical expressions.^{8,96,97,99,100,138} Based on eqn (2), it is possible to determine a point \mathbf{P}^* in parameter space that will minimize \tilde{Q} within a specified trust region around \mathbf{P}^0 . This region is defined here in terms of maximal allowed relative changes in each of the parameters over an iteration, set to 5% for all parameters optimized. Since the function is convex, the point \mathbf{P}^* is unique. Once determined, it is selected as a new point \mathbf{P}^0 to carry out simulations in view of a next iteration. Note that convexity only applies to the local approximation \tilde{Q} of eqn (2), but not to the objective function Q of eqn (1) itself, which will generally present many local minima in parameter space.

In practice, the optimization algorithm is carried out by an optimizer script (OPT), and involves the following steps: (1) select an initial guess for \mathbf{P}_i^0 at iteration $i = 0$; (2) perform $N_{\text{sim}}^{\text{cal}}$ simulations to get the simulated-data vector \mathbf{X}^{sim} along with the sensitivity matrix \mathbf{S} at \mathbf{P}_i^0 ; (3) calculate the real value $Q_i^{\text{real}} = Q(\mathbf{P}_i^0; \mathbf{X}^{\text{exp}})$ of the objective function at this point in

parameter space; (4) minimize $\tilde{Q}(\mathbf{P}; \mathbf{P}_i^0, \mathbf{X}^{\text{exp}})$ in eqn (2) with respect to \mathbf{P} starting from \mathbf{P}_i^0 and remaining within the trust radius, leading to \mathbf{P}_i^* ; (5) calculate the predicted value $Q_{i+1}^{\text{pred}} = \tilde{Q}(\mathbf{P}_i^*; \mathbf{P}_i^0, \mathbf{X}^{\text{exp}})$ of the objective function; (6) set \mathbf{P}_{i+1}^0 to \mathbf{P}_i^* , increment i , and iterate to step (2) until convergence.

Step (2) is the expensive part of the calculation. In contrast, step (4) is inexpensive, and carried out in practice using a simplex minimization. Considering steps (3) and (5), from $i = 1$ onward, the real value Q_i^{real} at iteration i can be compared to the predicted value Q_i^{pred} from the previous iteration, giving a hint about the accuracy of the linearization in eqn (2) within the imposed trust radius. If the algorithm terminates at iteration i_{max} , the final parameter set is $\mathbf{P}_{i_{\text{max}}}^0$ with the value $Q_{i_{\text{max}}}^{\text{real}}$ of the objective function. Although $\mathbf{P}_{i_{\text{max}}}^*$ and $Q_{i_{\text{max}+1}}^{\text{pred}}$ are available, it is preferable to stop at a force field with an explicitly calculated objective function. In the present application to the O + N family, a total of $i_{\text{max}} = 5$ iterations were performed to reach convergence. Given the setup adopted in the simulations and access to 408 processors (3 GHz Intel Xeon), *i.e.* one for each of the $N_{\text{sim}}^{\text{cal}}$ simulations to be carried out at each iteration, the full optimization required about three days of wall-clock computing time. The non-bonded parameters of the final force field are reported in Tables 3 and 4, along with ESI,[†] Table S2 for the covalent terms.

The simulations were performed using an in-house GROMOS-compatible simulation engine in C++ called SAMOS (SAM). The pure-liquid MD simulations were carried out under periodic boundary conditions based on cubic computational boxes containing 512 molecules, and in the isothermal-isobaric ensemble at the reference pressures P and temperatures T listed in ESI,[†] Table S1. The temperature was maintained close to its reference value using a Nosé–Hoover thermostat¹³⁹ with a coupling time of 0.1 ps. The pressure was maintained close to its reference value using an Andersen barostat¹⁴⁰ with a coupling time of 0.5 ps. The isothermal compressibility was set to $4.575 \times 10^{-4} \text{ kJ}^{-1} \text{ mol nm}^3$ for compatibility with the standard GROMOS setup.¹¹² The ideal-gas simulations relied on stochastic dynamics^{141–145} (SD). They were also carried out under periodic boundary conditions based on cubic computational boxes



Table 3 Atom types (or, equivalently, EE-types) of the proposed GROMOS-compatible force field for the O + N family, along with the final (optimized) values of the associated EE parameters. The $N_{\text{att}} = 47$ atom types (or, equivalently, $N_{\text{att}}^{\text{EE}} = 47$ EE-types) are listed along with their usage, the associated LJ-type (referring to the entries of Table 4), and the optimized values of the EE parameters, *i.e.* the hardness η and the electronegativity χ . The eight parameters corresponding to the aliphatic (united-)atom types CH0 to CH3 (not involved in the EE scheme, zero charge) need not be specified. The η and χ values of CH0_N_amd (between parentheses) are estimated by the average values over the three other CH*n*_N_amd types in the absence of parametrization target (no experimental data found for a compound involving this type). The initial values of the EE parameters (used to start the optimization) can be found in ESI,† Table S3

Idx	Atom type (EE-type)	LJ-type	η [e^{-1} V]	χ [V]	Usage
Aliphatic carbon (united-)atoms					
1	CH0	CH0	—	—	CH ₀ carbon atom (methanetriyl group)
2	CH1	CH1	—	—	CH ₁ carbon united-atom (methanetriyl group)
3	CH2	CH2	—	—	CH ₂ carbon united-atom (methylene group)
4	CH3	CH3	—	—	CH ₃ carbon united-atom (methyl group)
Ether					
5	O_eth	OR	35.447	37.782	Ether oxygen atom
6	CH0_O_eth	CH0	2.814	21.953	Alkoxyated CH ₀ atom
7	CH1_O_eth	CH1	6.568	19.993	Alkoxyated CH ₁ united-atom
8	CH2_O_eth	CH2	13.981	16.921	Alkoxyated CH ₂ united-atom
9	CH3_O_eth	CH3	22.115	14.233	Alkoxyated CH ₃ united-atom
Aldehyde					
10	H_CO_ald	HC	17.156	21.368	Aldehyde hydrogen atom
11	C_ald	C=O	15.981	16.008	Aldehyde carbonyl carbon atom
12	O_ald	O=C	8.821	23.815	Aldehyde carbonyl oxygen atom
Ketone					
13	C_ket	C=O	36.256	15.724	Ketone carbonyl carbon atom
14	O_ket	O=C	17.326	38.611	Ketone carbonyl oxygen atom
Ester					
15	H_CO_est	HC	16.648	24.832	Formate ester hydrogen atom
16	C_est	C=O	21.158	11.713	Ester carbonyl carbon atom
17	O_est	O=C	12.815	28.520	Ester carbonyl oxygen atom
18	O_R_est	OR	25.554	29.741	Ester acylated oxygen atom
19	CH0_O_est	CH0	9.033	19.163	Ester oxygen-linked CH ₀ atom
20	CH1_O_est	CH1	12.401	18.349	Ester oxygen-linked CH ₁ united-atom
21	CH2_O_est	CH2	36.940	13.503	Ester oxygen-linked CH ₂ united-atom
22	CH3_O_est	CH3	19.577	16.726	Ester oxygen-linked CH ₃ united-atom
Alcohol					
23	H_ol	HB	35.794	17.095	Hydroxyl hydrogen atom
24	O_ol	OH	31.057	46.647	Hydroxyl oxygen atom
25	CH0_O_ol	CH0	36.145	19.963	Hydroxylated CH ₀ atom
26	CH1_O_ol	CH1	32.298	18.392	Hydroxylated CH ₁ united-atom
27	CH2_O_ol	CH2	30.423	17.200	Hydroxylated CH ₂ united-atom
28	CH3_O_ol	CH3	29.995	17.135	Hydroxylated CH ₃ united-atom
Carboxylic acid					
29	H_CO_acd	HC	28.828	12.561	Formic acid hydrogen atom
30	C_acd	C=O	29.109	14.269	Carboxylic acid carbonyl carbon atom
31	O_acd	O=C	36.286	45.338	Carboxylic acid carbonyl oxygen atom
32	H_O_acd	HB	31.782	14.733	Carboxylic acid hydroxyl hydrogen atom
33	O_H_acd	OH	41.291	36.891	Carboxylic acid hydroxyl oxygen atom
Amine					
34	H_N_amn	HB	44.559	7.467	Amine hydrogen atom
35	N_amn	N_amn	39.441	47.379	Amine nitrogen atom
36	CH0_N_amn	CH0	33.307	13.667	Aminated CH ₀ atom
37	CH1_N_amn	CH1	31.777	15.046	Aminated CH ₁ united-atom
38	CH2_N_amn	CH2	36.009	12.981	Aminated CH ₂ united-atom
39	CH3_N_amn	CH3	31.313	13.650	Aminated CH ₃ united-atom
Amide					
40	H_N_amd	HB	29.386	14.179	Amide nitrogen-linked hydrogen atom
41	C_amd	C=O	30.320	14.492	Amide carbonyl carbon atom
42	O_amd	O=C	28.316	42.678	Amide carbonyl oxygen atom
43	N_amd	N_amd	30.230	37.064	Amide acylated nitrogen atom
44	CH0_N_amd	CH0	(35.585)	(15.614)	Amide nitrogen-linked CH ₀ atom (estimated)
45	CH1_N_amd	CH1	35.544	19.373	Amide nitrogen-linked CH ₁ united-atom
46	CH2_N_amd	CH2	36.828	13.164	Amide nitrogen-linked CH ₂ united-atom
47	CH3_N_amd	CH3	34.383	14.305	Amide nitrogen-linked CH ₃ united-atom

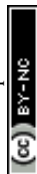


Table 4 LJ-types of the proposed GROMOS-compatible force field for the O + N family, along with the final (optimized) values of the associated LJ parameters. The $N_{\text{att}}^{\text{LJ}} = 12$ LJ-types are listed along with their usage and the optimized values of the LJ parameters, *i.e.* the collision diameter σ and the well depth ϵ . These LJ-types are invoked in the specification of the $N_{\text{att}} = 47$ atom types of Table 3. The LJ interaction parameters corresponding to a non-hydrogen-bonding LJ-type pair are obtained by applying a geometric-mean combination rule^{132,133} to the values of σ and ϵ associated with the LJ-types of the two involved atoms. For a hydrogen-bonding LJ-type pair, the alternative values $\bar{\sigma}$ and $\bar{\epsilon}$ associated with the LJ-types of the two involved atoms are used instead. The value of C_6 is common, so that only $\bar{\sigma}$ is specified in the table, while $\bar{\epsilon}$ can be deduced as $\bar{\epsilon} = \epsilon\sigma^6/\bar{\sigma}^6$ (0.359, 0.614, 0.619, 0.619, and 0.504 kJ mol⁻¹ for OR, O=C, OH, N_amn, and N_amd, respectively). The hydrogen-bonding LJ-type pairs are those involving a potential hydrogen-bond donor (OH, N_amn, or N_amd) and a potential hydrogen-bond acceptor (OR, O=C, OH, N_amn, or N_amd). Finally, for third covalent neighbors, the alternative values σ^* and ϵ^* are to be used instead. Note that the ten LJ parameters σ and ϵ corresponding to the aliphatic (united-)atom types CH0 to CH3 and to the polar hydrogen atom HB were not subject to optimization, but taken directly from the 2016H66 parameter set.¹¹² The same applies to all the σ^* and ϵ^* parameters. The $\bar{\sigma}$ value of OR (between parentheses) is kept identical to σ in the absence of parametrization target (no hydrogen bonding is possible in molecules containing only ether groups). The initial values of the LJ parameters (used to start the optimization) can be found in ESI,[†] Table S4

LJ type	σ [nm]	$\bar{\sigma}$	σ^*	ϵ [kJ mol ⁻¹]	ϵ^*	Usage
Carbon						
CH0	0.664	—	0.336	0.007	0.406	CH ₀ carbon atom (methanetetryl group)
CH1	0.502	—	0.330	0.095	0.567	CH ₁ carbon united-atom (methanetriyl group)
CH2	0.407	—	0.316	0.411	1.176	CH ₂ carbon united-atom (methylene group)
CH3	0.375	—	0.309	0.867	1.946	CH ₃ carbon united-atom (methyl group)
C=O	0.345	—	0.336	0.326	0.406	Carbonyl carbon atom
Oxygen						
OR	0.301	(0.301)	0.287	0.359	1.011	Ether oxygen atom
O=C	0.296	0.313	0.262	0.857	1.725	Carbonyl oxygen atom
OH	0.287	0.298	0.287	0.776	1.011	Hydroxyl oxygen atom
Nitrogen						
N_amn	0.312	0.308	0.298	0.572	0.877	Amine nitrogen atom
N_amd	0.320	0.312	0.298	0.429	0.877	Amide nitrogen atom
Hydrogen						
HC	0.223	—	0.223	0.119	0.119	Carbonyl-linked hydrogen atom
HB	0.000	—	0.000	0.000	0.000	Oxygen- or nitrogen-linked hydrogen atom

containing 512 molecules but with all intermolecular interactions turned off, and in the canonical ensemble at the same temperatures as the corresponding pure-liquid simulations. The friction coefficient was set to 2 ps⁻¹.

Newton's equations of motion (MD) or Langevin's equations of motion (SD) were integrated using the leap-frog scheme^{141,146} with a timestep of 2 fs. Constraints on all bond lengths were enforced by application of the SHAKE procedure¹⁴⁷ with a relative geometric tolerance of 10⁻⁵. The non-bonded interactions were calculated using a twin-range scheme¹⁴⁸ based on charge-group distances with short- and long-range cutoff radii set to 0.8 and 1.4 nm, respectively, and an update frequency of 5 timesteps for the short-range pairlist and intermediate-range interactions. In the pure-liquid simulations, the mean effect of the omitted electrostatic interactions beyond the long-range cutoff distance was reintroduced using a reaction-field correction.^{149,150} The corresponding permittivities are listed in ESI,[†] Table S1.

Additional details about the simulation protocols can be found in ref. 8. The protocol applied here is essentially the same, except for the use of a doubled friction coefficient in SD and a SHAKE tolerance reduced by a factor of ten. These modifications were found necessary for a better temperature control in the case of molecules with explicit hydrogen atoms. The GROMOS-compatible molecular topologies and equilibrated liquid configurations for the $N_{\text{sim}}^{\text{tot}} = 1175$ molecules considered

here can be downloaded freely from the internet under ref. 125 (version 1.0).

3 Results and discussion

The evolution of the objective function Q against the iteration number i is illustrated in Fig. 5. The graph shows both the real value Q_i^{real} at iteration i calculated according to eqn (1) and its predicted value Q_i^{pred} calculated based on the simulations at iteration $i - 1$ according to eqn (2). After two iterations, the predicted and real values agree very well, indicating that the linear approximation is accurate within the chosen trust region for the parameter variations. During the first three iterations, the objective function drops sharply. It is essentially converged after four iterations, and the fifth one brings no further improvement. The final force field corresponds to iteration $i = 5$ and is associated with a value of 1.69 for Q .

The evolution of the $N_{\text{prm}}^{\text{cal}} = 102$ non-bonded interaction parameters subject to calibration against the iteration number i is shown in Fig. 6 and 7 for the LJ and EE parameters, respectively. The largest parameter changes typically occur within the first three iterations, where the decrease of Q is most pronounced. However, the magnitudes of these changes still remain limited. This is not unexpected considering that the initial values selected to start the optimization (ESI,[†] Section S4) are ported from the 2016H66



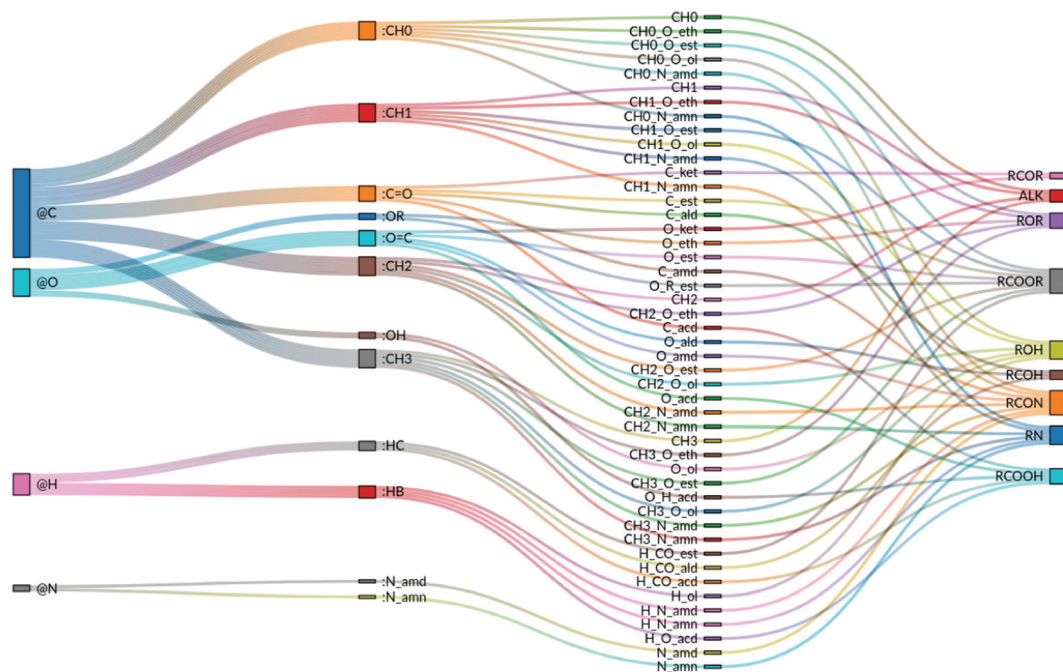


Fig. 4 Correspondence between the 4 elements, the 12 LJ-types, the 47 atom-types (EE-types) and 9 chemical functional groups involving the latter atom types in the proposed GROMOS-compatible force field for the O + N family. The first column refers to the elements, the second to the LJ-types (see Table 4), the third to the atom- or EE-types (see Table 3), and the fourth to the chemical functional groups (see Table 2; note that ALK was added for alkanes and that RCOOR, RCON and RN group all the esters, amides and amines, respectively, of Table 2).

parameter set,¹¹² and thus already expected to perform reasonably well.

The level of agreement between the optimized force field and experiment in terms of ρ_{liq} and ΔH_{vap} is illustrated in Fig. 8 for both the calibration (top panels) and validation (bottom panels) sets. The corresponding numerical values can be found in ESI,[†] Section S6 (Table S.7). The statistics per compound

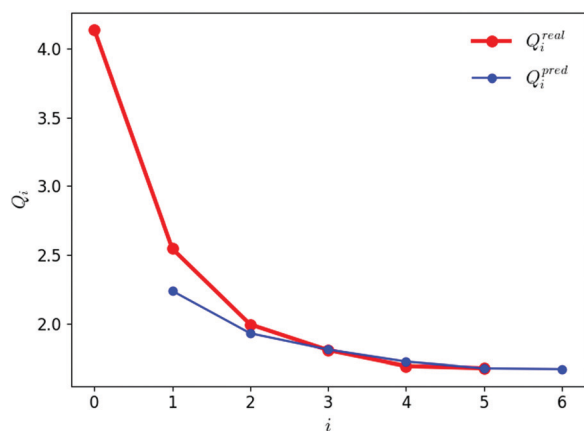


Fig. 5 Evolution of the predicted and real values of the objective function against the iteration number during the force-field parameter optimization. The real value Q_i^{real} at iteration i is calculated according to eqn (1). The predicted value Q_i^{pred} is calculated according to eqn 2 based on the simulations at iteration $i - 1$. The first simulations at $i = 0$ using the initial parameter set leads to a first real value Q_0^{real} and a first predicted value Q_0^{pred} . The last simulations at $i = 5$ using the final (optimized) parameter set leads to the final real value Q_5^{real} (and a predicted value Q_5^{pred} , which is discarded).

types is provided in Table 5 and illustrated graphically in Fig. 9 for both the calibration (top panels) and validation (bottom panels) sets, distinguishing between molecules presenting one, two, three, or four occurrences of the specific functional group. In this statistics, two classes of compounds are also considered separately, namely the non-hydrogen-bonding ones (NHB), including ethers, aldehydes, ketones, esters, tertiary amines, and tertiary amides, and the hydrogen-bonding ones (HBD), including alcohols, carboxylic acids, the other amines, and the other amides.

Considering all compounds of the calibration set, the optimized force field has an overall root-mean-square deviation (RMSD) relative to experiment of 29.9 kg m^{-3} for ρ_{liq} and 4.1 kJ mol^{-1} for ΔH_{vap} . Thus, given the choice of the scaling factors s_n in eqn (1), namely 20 kg m^{-3} for ρ_{liq} and 1 kJ mol^{-1} for ΔH_{vap} , the final value of 1.69 for is dominated by the ΔH_{vap} errors. A similar and somewhat stronger bias is also observed in the validation set, with final RMSD values of 22.4 kg m^{-3} for ρ_{liq} and 5.5 kJ mol^{-1} for ΔH_{vap} . This observation contrasts with the results of the previous application of CombIFF to the saturated haloalkanes,⁸ where the contributions of the two types of observables were of comparable magnitudes, with RMSD values of $49.8 (27.6) \text{ kg m}^{-3}$ for ρ_{liq} and $2.7 (1.8) \text{ kJ mol}^{-1}$ for ΔH_{vap} over the calibration (validation) set. For the O + N family, one also observes a tendential underestimation of the density and overestimation of the vaporization enthalpy, with average deviations (AVED) relative to experiment of $-6.9 (-9.1) \text{ kg m}^{-3}$ for ρ_{liq} and $+0.5 (+2.6) \text{ kJ mol}^{-1}$ for ΔH_{vap} over the calibration (validation) set. The corresponding values for the saturated haloalkanes⁸ were comparatively smaller in



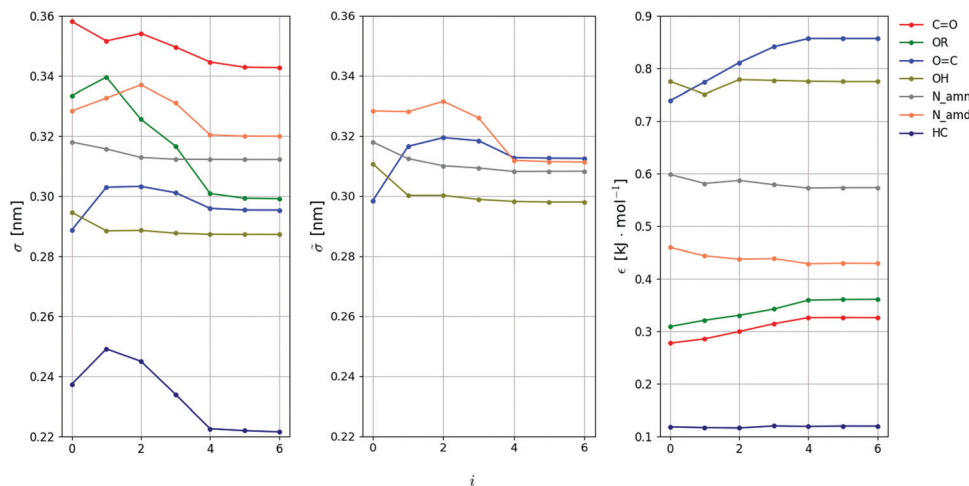


Fig. 6 Evolution of the LJ interaction parameters against the iteration number during the force-field parameter optimization. The value of each parameter is reported at iteration i . The 18 parameters considered are the collision diameter σ and well depth ϵ for 7 non-hydrogen-bonding LJ-types, along with the collision diameter $\tilde{\sigma}$ for 4 (out of the 5) potentially hydrogen-bonding LJ-types, see Table 4. Since the value of C_6 is common to the two sets, only $\tilde{\sigma}$ is displayed, while $\tilde{\epsilon}$ can be deduced as $\tilde{\epsilon} = \epsilon\sigma^6/\tilde{\sigma}^6$. The $\tilde{\sigma}$ value of the LJ-type OR (not shown) is kept identical to σ in the absence of parametrization target (no hydrogen bonding is possible in molecules containing only ether groups). The final force-field parameters are those corresponding to iteration $i = 5$ (the values at $i = 6$ correspond to proposed changes for a next iteration, and are discarded).

magnitude, namely +1.2 (+5.3) kg m⁻³ for ρ_{liq} and 0.0 (0.0) kJ mol⁻¹ for ΔH_{vap} .

In both calibration and validation sets, the errors are not distributed homogeneously across the different chemical groups. The ρ_{liq} values are nearly systematically underestimated, but the monocarboxylic acids, the amines (except primary), and the secondary amides present positive deviations instead. The underestimation is particularly pronounced for the dicarboxylic acids in the validation set (AVED value of -98.4 kg m⁻³). Similarly, the ΔH_{vap} values are nearly systematically overestimated, but a few compound groups in the calibration set present slightly negative deviations, all below 1.0 kJ mol⁻¹ in magnitude except for the diethers, the monocarboxylic acids, and the primary and secondary amides (down to -2.4 kJ mol⁻¹). The overestimation is particularly pronounced for the formic diesters and the triketones (AVED values of +15.5 and +20.2 kJ mol⁻¹, respectively). For both ρ_{liq} and ΔH_{vap} , the deviations relative to experiment within a given group of molecules are also seen to nearly systematically increase upon increasing the number of occurrences of the functional group in the molecule.

The most prominent outliers of both calibration and validation sets, namely the 58 molecules (62 simulations) with deviations larger than 80.0 kg m⁻³ for ρ_{liq} and/or larger than 8.0 kJ mol⁻¹ for ΔH_{vap} , are discussed in ESI,[†] Section S7 (Table S8 and Fig. S4). Among these 58 molecules, 40 encompass two functional groups (two notable exceptions being methanol and trimethylamine). Methanol was also used in the calibration of 2016H66. Similarly to the results shown here, ρ_{liq} and ΔH_{vap} are overestimated, but 2016H66 leads to significantly lower errors in terms of ρ_{liq} .

Finally, a comparison between the present force field and the original 2016H66 parameter set¹¹² is presented in ESI,[†] Section S8 (Tables S9–S11). After optimization of the non-bonded interaction parameters, the overall agreement with experiment is enhanced, although not very pronouncedly.

Moreover, it is also non-systematic across the different groups of molecules. In particular, the accuracy enhancement is most significant for polyfunctional molecules compared to monofunctional ones, although these molecules still present comparatively large deviations after the parameter refinement (see above).

4 Conclusions

In this article, the CombiFF scheme⁸ (Fig. 1) was applied to the calibration of a GROMOS-compatible united-atom force field for the O + N family of compounds (Table 2). This force field relies on 47 atom types (equivalent to EE-types) distributed over 12 LJ-types. The aliphatic (united-)atom EE- and LJ-types are common with the previously optimized (halo-)alkane force field.⁸ The calibration of the 102 non-bonded interaction parameters was performed here against 579 experimental liquid densities ρ_{liq} and vaporization enthalpies ΔH_{vap} concerning 339 small molecules, which represents an observable-to-parameter ratio of 5.7. A collection of 1133 additional ρ_{liq} and ΔH_{vap} values concerning 836 other molecules was used for subsequent validation, leading to an effective observable-to-parameter ratio of 16.8. The calibration of the force field required five iterations, corresponding to about three days of wall-clock computing time using 408 processors.

The optimized (final) force-field parameters (reported in Tables 3 and 4, along with ESI,[†] Table S2) lead to RMSD (AVED) values of 29.9 (-6.9) kg m⁻³ for ρ_{liq} and 4.1 (0.5) kJ mol⁻¹ for ΔH_{vap} over the calibration set. The corresponding values over the validation set are 22.4 (-9.1) kg m⁻³ and 5.5 (2.6) kJ mol⁻¹, respectively. Considering the two sets together (1175 molecules, 1712 data points), the relevant RMSD values are 24.8 kg m⁻³ for ρ_{liq} and 4.9 kJ mol⁻¹ for ΔH_{vap} . This shows that a very good level



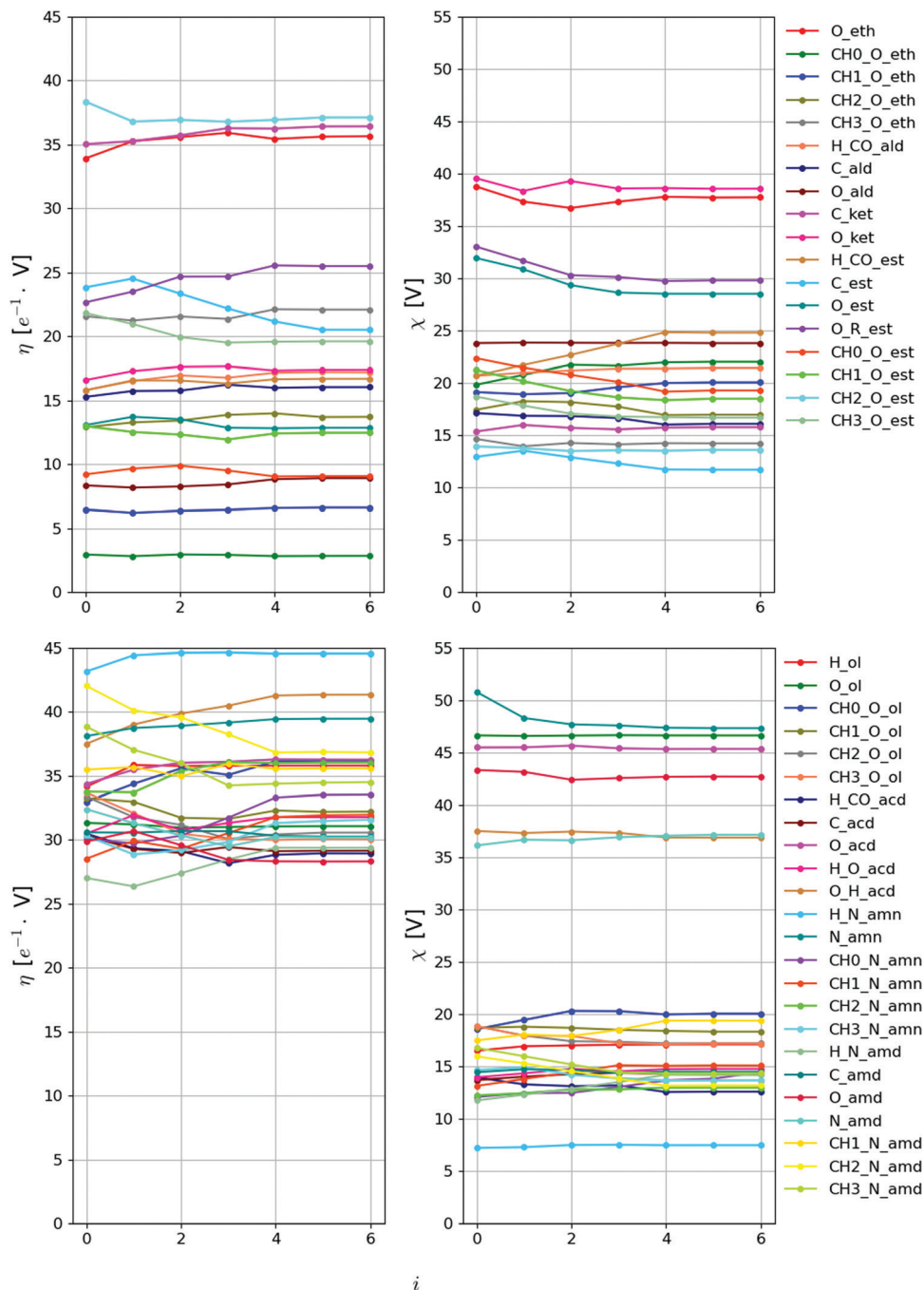


Fig. 7 Evolution of the EE interaction parameters against the iteration number during the force-field parameter optimization. The value of each parameter is reported at iteration i . The 84 parameters considered are the hardness η and electronegativity χ for the 42 EE-types subject to optimization, see Table 3. The final force-field parameters are those corresponding to iteration $i = 5$ (the values at $i = 6$ correspond to proposed changes for a next iteration, and are discarded).

of agreement with experiment can be achieved for the O + N family using a simple united-atom force field, and that the overall observable-to-parameter ratio of 16.8 is sufficient to define an appropriate set of parameters. However, appropriateness does not imply uniqueness. As observed previously,⁸ the solution reached upon optimization is essentially unique for the LJ interaction parameters, but a larger extent of degeneracy is observed for the EE parameters. This degeneracy leads to a

significant variability in these parameters, whereas the atomic partial charges themselves are not significantly affected.¹⁵¹

A more detailed analysis of the errors reveals three main trends: (i) the residual errors are biased towards larger discrepancies in terms of ΔH_{vap} relative to ρ_{liq} (compared to our previous work on haloalkanes,⁸ where the errors were balanced between the two observables); (ii) the ρ_{liq} values are nearly systematically underestimated and the ΔH_{vap} values are



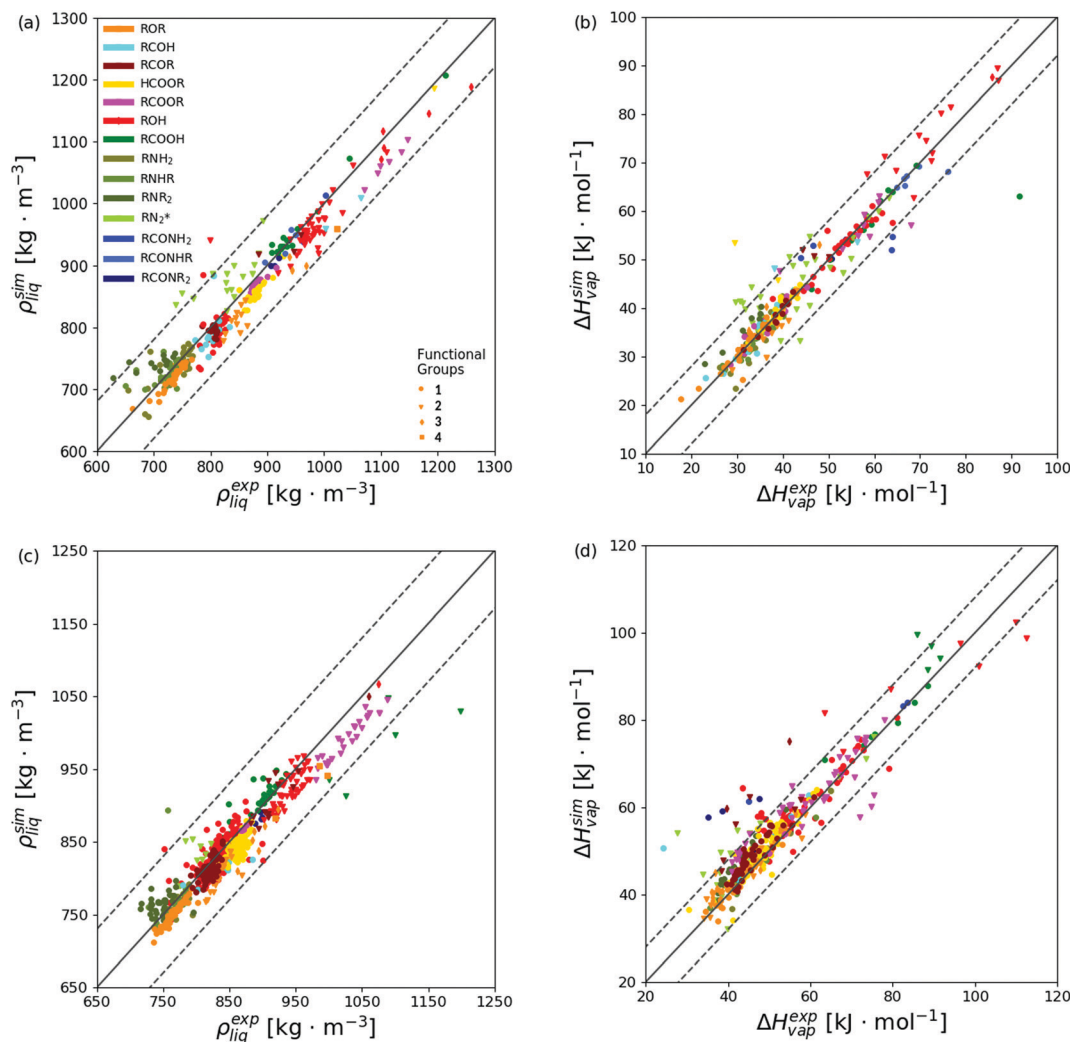


Fig. 8 Simulated versus experimental properties based on the optimized force field. The results are reported for the calibration set (a and b) and the validation set (c and d), considering the pure-liquid density ρ_{liq} and vaporization enthalpy ΔH_{vap} . The colors are selected according to the 14 groups of molecules defined in Table 2. The symbols are selected according to the number of occurrences of the functional group in a molecule. The diagonal solid lines indicate perfect agreement, and the ranges between the two dashed lines indicate agreement within $\pm 80.0 \text{ kg m}^{-3}$ for ρ_{liq} or $\pm 8.0 \text{ kJ mol}^{-1}$ for ΔH_{vap} (scaling factors s_n used in eqn (1) multiplied by 4 and 8, respectively). The corresponding numerical values can be found in ESI,† Section S6, the statistics per groups of molecules in Table 5 and Fig. 9, and the information about the main outliers (points outside the ranges defined by the dashed lines in any of the four panels) in ESI,† Section S7.

nearly systematically overestimated relative to experiment; (iii) the deviations nearly systematically increase with the number of occurrences of the functional group in a molecule. These observations may result in part from the following three factors.

First, the ΔH_{vap} comparison is likely to be affected by intrinsically larger experimental and computational errors compared to the ρ_{liq} comparison. Experimentally, the ΔH_{vap} values result from more complicated measurements (temperature-dependence of the vapor pressure or calorimetry at the boiling point for ΔH_{vap} vs. simple volumetric measurement for ρ_{liq}), and are potentially also affected by more significant ambiguities related to the measurement pressure (isochoric measurement at the vapor pressure of the liquid vs. isobaric measurement in the presence of an inert gas) and the possible

need for and/or application of real-gas corrections. Computationally, the ΔH_{vap} values probe a less local and more collective property (energetics for ΔH_{vap} vs. packing for ρ_{liq}), and are potentially also affected by more significant calculation ambiguities (influence of the treatment of long-range interactions, neglect of the difference in polarization between liquid-phase and gas-phase molecules when using a non-polarizable force field).

Second, there appears to be conflicting requirements imposed by the two types of observables. If a given parameter variation enhances the tightness of molecular packing, it will generally also increase the strength of attractive intermolecular interactions. As a result, one generally expects the changes in ρ_{liq} and ΔH_{vap} induced by a given parameter variation to be positively correlated. In the optimized force field proposed here for the O+N family, a compromise is reached and the residual



Table 5 Statistics concerning the discrepancies between simulated and experimental properties based on the optimized force field. The results are reported for the calibration set (left) and the validation set (right). For selected groups of molecules (Table 2) and numbers m of occurrences of the functional group in the molecule, the number N_{ρ}^{cal} of experimental ρ_{liq} values and the number $N_{\Delta H}^{\text{cal}}$ of experimental ΔH_{vap} values are reported, along with the root-mean-square deviation (RMSD) and the average deviation (AVED) between simulation and experiment for both properties. The last three lines refer to non-hydrogen-bonding (NHB), hydrogen-bonding (HBD), and the entire set (All) of molecules. Considering calibration and validation sets together (1175 molecules, 1712 data points), the overall RMSD values are 24.8 kg m^{-3} for ρ_{liq} and 4.9 kJ mol^{-1} for ΔH_{vap} . The corresponding AVED values are -8.4 kg m^{-3} and 1.7 kJ mol^{-1} , respectively, and the corresponding mean unsigned deviations are 18.1 kg m^{-3} and 3.1 kJ mol^{-1} , respectively. This data is illustrated graphically in Fig. 9

Code	m	Calibration						Validation					
		N_{ρ}^{cal}	$\rho_{\text{liq}} [\text{kg m}^{-3}]$		$N_{\Delta H}^{\text{cal}}$	$\Delta H_{\text{vap}} [\text{kJ mol}^{-1}]$		N_{ρ}^{val}	$\rho_{\text{liq}} [\text{kg m}^{-3}]$		$N_{\Delta H}^{\text{val}}$	$\Delta H_{\text{vap}} [\text{kJ mol}^{-1}]$	
			RMSD	AVED		RMSD	AVED		RMSD	AVED		RMSD	AVED
ROR	1	28	15.7	-13.8	23	1.8	0.5	56	15.4	-14.6	19	1.5	0.1
	2	11	37.8	-34.8	12	2.7	-1.8	28	30.4	-28.2	17	2.1	0.1
	3	3	51.5	-48.5	4	3.5	3.0	6	39.0	-36.9	11	3.7	2.8
	4	1	64.2	-64.2	2	0.6	0.6	3	42.1	-40.7	3	1.5	1.2
	1-4	43	28.4	-22.8	41	2.3	0.1	93	24.0	-21.0	50	2.3	0.8
RCOH	1	16	27.3	-9.4	13	1.7	0.6	21	21.9	-16.1	8	9.6	5.2
	2	2	50.1	-49.8	2	7.3	3.8	0	—	—	0	—	—
	1-2	18	30.7	-13.9	15	3.1	1.0	21	21.9	-16.1	8	9.6	5.2
RCOR	1	11	15.2	-8.4	11	0.9	-0.1	73	12.8	-7.4	49	3.3	2.8
	2	3	24.8	-0.5	3	5.1	4.1	11	18.0	-5.2	6	10.5	9.2
	3	0	—	—	0	—	—	1	9.6	-9.6	2	20.2	20.2
	1-3	14	17.7	-6.7	14	2.5	0.8	85	13.5	-7.1	57	5.9	4.1
HCOOR	1	12	12.2	-11.3	10	1.6	0.5	4	2.6	-2.4	3	0.8	0.3
	2	1	7.5	-7.5	2	17.6	15.5	0	—	—	0	—	—
	1-2	13	11.9	-11.0	12	7.3	3.0	4	2.6	-2.4	3	0.8	0.3
RCOOR	1	20	24.1	-22.0	19	1.1	0.5	115	21.1	-18.8	70	3.2	1.9
	2	6	45.4	-45.3	9	4.9	-0.6	27	37.5	-36.5	38	5.9	2.0
	1-2	26	30.4	-27.3	28	2.9	0.2	142	25.1	-22.2	108	4.3	1.9
ROH	1	33	28.4	-14.0	32	2.1	-0.2	247	16.1	-3.1	55	5.1	2.3
	2	45	34.5	-17.9	12	5.2	2.9	52	21.6	-15.6	10	8.7	0.9
	3	5	38.4	-27.0	1	1.9	1.9	1	8.2	-8.2	0	—	—
	1-3	83	32.5	-16.9	45	3.2	0.7	300	17.2	-5.3	65	5.8	2.0
RCOOH	1	17	12.0	6.3	13	8.0	-2.4	35	17.1	11.7	7	3.1	1.1
	2	0	—	—	0	—	—	5	107.6	-98.4	4	8.1	6.8
	1-2	17	12.0	6.3	13	8.0	-2.4	40	41.3	-2.1	11	5.5	3.2
RNH ₂	1	34	20.2	-2.0	22	2.1	-0.3	17	14.1	-5.9	9	2.2	0.3
RNHR	1	25	22.0	5.4	19	1.4	-0.5	16	36.2	7.8	10	1.8	0.0
RNR ₂	1	13	46.1	39.4	12	3.6	3.2	31	25.6	18.4	17	4.9	4.7
RN ₂ *	2	16	57.8	48.1	31	6.1	2.1	9	36.2	28.0	19	8.9	5.1
	1-2	88	34.9	15.3	84	4.1	1.0	73	27.8	11.6	55	6.0	3.3
RCONH ₂	1	4	15.6	-8.0	4	8.7	-2.0	1	2.2	-2.2	3	9.5	5.9
RCONHR	1	4	10.7	3.4	7	3.1	-1.4	0	—	—	0	—	—
RCONR ₂	1	4	8.7	-6.1	2	0.5	-0.1	6	11.6	-9.0	8	12.5	9.6
	1	12	12.0	-3.6	13	5.4	-1.4	7	10.8	-8.1	11	11.8	8.6
NHB	—	131	29.0	-12.9	124	3.5	0.9	382	22.3	-14.5	251	5.2	2.7
HBD	—	183	30.6	-2.6	141	4.6	0.2	383	22.4	-3.6	117	6.1	2.4
All	—	314	29.9	-6.9	265	4.1	0.5	765	22.4	-9.1	368	5.5	2.6

negative errors for ρ_{liq} and positive errors for ΔH_{vap} cannot be relieved by further parameter variations. This situation may arise from inaccurate experimental data (incompatibility between ρ_{liq} and ΔH_{vap} observables) and/or from a lack of flexibility of the force-field functional form (*e.g.* limitations of the LJ combination rules or of the EE charge scheme, absence of explicit electronic polarizability).

Third, the lower force-field accuracy for polyfunctional compounds may result from a combination of multiple factors:

(i) insufficient number of polyfunctional compounds in the calibration set (*e.g.* dicarboxylic acids and triketones are only included in the validation set); (ii) inaccurate representation of the conformational properties of the molecules (torsional-dihedral and third-neighbor parameters ported rather crudely from the 2016H66 force field without reoptimization, see ESI,[†] Fig. S3); (iii) inaccurate representation of the charge transfers within the molecules (charge transfers in the EE scheme are only allowed here within charge groups).



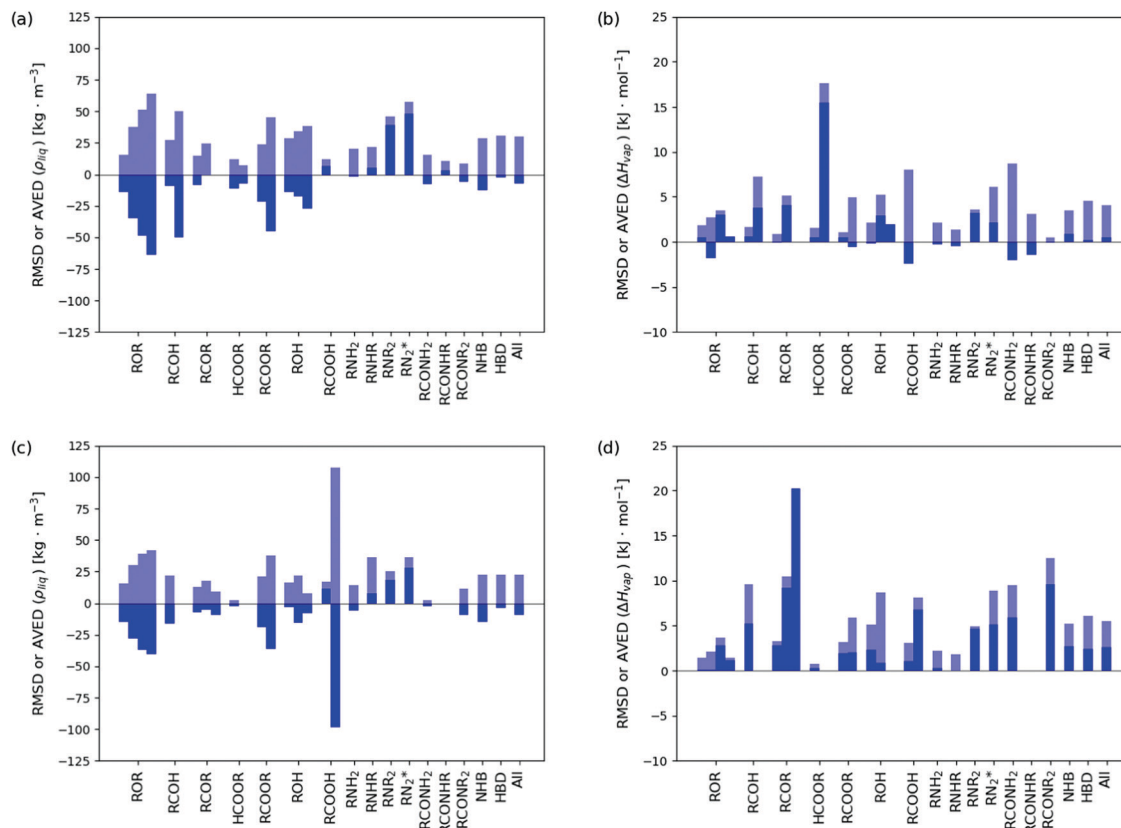


Fig. 9 Statistics concerning the discrepancies between simulated and experimental properties based on the optimized force field. The results are reported for the calibration set (a and b) and the validation set (c and d). The values of the root-mean square deviation (RMSD, transparent bar) and average deviation (AVED, solid bar) in ρ_{liq} (a and c) and ΔH_{vap} (b and d) for selected groups of molecules (Table 2) are compared. The successive bars in each group correspond to compounds containing 1, 2, 3 or 4 occurrences of the functional group. The last three bars refer to non-hydrogen-bonding (NHB), hydrogen-bonding (HBD), and the entire set (All) of molecules. The corresponding numerical values can be found in Table 5.

Compared to the 2016H66 parameter set,¹¹² the calibration/validation of which involved 43 monofunctional molecules of the O + N family, the present reoptimization leads to a general improvement in terms both ρ_{liq} and ΔH_{vap} , especially for polyfunctional molecules. However, this improvement is neither extremely pronounced nor entirely systematic for the monofunctional compounds. This suggests that the 2016H66 parametrization was already close to optimal for these 43 compounds.

A key component of the selected force-field representation is that the atomic partial charges are generated using an EE scheme, which permits to take into account induction effects. Thus, the atomic partial charges can change according to the chemical environment within the considered molecules. For example, primary, secondary and tertiary alcohols have different charge sets in the present force field. In contrast, in the 2016H66 set,¹¹² the charges were fixed for a given functional group, irrespective of its chemical environment. This feature is expected to enhance significantly the model flexibility, especially when multiple functional groups are present within the same molecule. However, the requirement to define small neutral charge groups and the corresponding definitions adopted in this work still represent a limitation on the representation of inductive effects. For example, in compounds involving a carbonyl group, the C_{α} carbon (united-atom (*i.e.* directly attached

to the carbonyl group) has a charge of zero (which is not the case for a H_{α}). This is consistent with the charge-group choices made in the 2016H66 set,¹¹² but clearly at odds with chemical intuition.

Generally speaking, a number of the force-field representation choices made here will have to be addressed again in future work, and the parameter optimization repeated accordingly. This will include in particular a reconsideration of: (i) the parameters that were not reoptimized here, *i.e.* the bond-stretching and bond-angle bending parameters, the torsional-dihedral and third-neighbor interaction parameters, and the non-bonded interaction parameters of the aliphatic types; (ii) the choice of the EE- and LJ-type sets, *e.g.* by introducing distinct LJ-types for alcohol *vs.* acid as well as ether *vs.* ester oxygen atoms; (iii) the choice of model resolution, *i.e.* united- *vs.* all-atom; (iv) the choice of a combination rule, *e.g.* geometric-mean *vs.* others, and its possible (partial) by-passing; (v) the restriction of EE charge flows to small neutral charge groups (see above). The latter restriction could for example be alleviated by using a more general EE scheme⁸ (see Conclusion section therein) involving damped charge transfers throughout the entire molecule, applied with a smooth atom-based truncation of the non-bonded interactions in the simulations.¹⁵² Efforts are currently in progress along these different lines.

A possible further development of CombiFF would be the design of a polarizable force field of the fluctuating-charge



type,^{153–163} where an on-the-fly EE scheme would incorporate the effect of the configuration-dependent (*i.e.* local and instantaneous) electric potential on the atomic partial charges of each molecule during the MD simulation (now also including the *J*-terms for intramolecular Coulombic effects beyond first- and second-neighbors). A particularly appealing feature of this development is that it would not require any additional force-field parameters, but a mere CombiFF recalibration of the existing ones under application of the fluctuating-charge scheme.

Work is also in progress to expand the CombiFF calibration/validation to other chemical families, to polyfunctional molecules mixing different types of functional groups, and to the consideration of further thermodynamic, transport, and dielectric properties of the liquid (as well as properties concerning the gas and solid phases). In particular, for calibration, it might be of interest to consider vapor pressures P_{vap} in addition to (or instead of) vaporization enthalpies ΔH_{vap} as a target for probing the intermolecular energetics. The quantity P_{vap} is more readily available experimentally and easier to measure (thus likely affected by smaller uncertainties). The price to pay is that its calculation is more difficult than that of ΔH_{vap} , as it corresponds to a free-energy (rather than energy) calculation. In terms of validation, it will be essential to assess the accuracy of the CombiFF force fields not only in terms of the optimization targets ρ_{liq} and ΔH_{vap} , but also in terms of other properties. Such an assessment considering both the previously reported (halo-)alkane force field⁸ and the present O + N force field has already been performed in terms of nine additional properties (surface-tension coefficient γ , isothermal compressibility κ_{T} , isobaric thermal expansion coefficient α_{p} , isobaric heat capacity c_{p} , static relative dielectric permittivity ϵ , self-diffusion coefficient D , shear viscosity η , hydration free energy ΔG_{wat} and free energy of solvation ΔG_{che} in cyclohexane). The calculated values of these additional properties show reasonable to good agreement with experiment, except for c_{p} , D and η , where larger discrepancies are observed, in large part related to the classical treatment of the vibrations and the use of united atoms. These results will be reported in a forthcoming article.

Conflicts of interest

There are no conflicts to declare.

Note added after first publication

This article replaces the version published on 19th July 2021, which contained various errors.

Acknowledgements

Financial support by the Swiss National Science Foundation (Grant No. 200021-175944) is gratefully acknowledged. The authors are also grateful to Bill Acree, Sadra Kashef Ol Gheta, Salomé Rieder, Alžbeta Kubincová, and Bruno Horta for insightful discussions and useful suggestions.

References

- 1 T. A. Halgren, *Curr. Opin. Struct. Biol.*, 1995, **5**, 205.
- 2 P. H. Hünenberger and W. F. van Gunsteren, Empirical classical interaction functions for molecular simulations, in *Computer simulation of biomolecular systems, theoretical and experimental applications*, ed. W. F. van Gunsteren, P. K. Weiner, A. J. Wilkinson, Kluwer/Escom Science Publishers, Dordrecht, The Netherlands, 1997, vol. 3, pp. 3–82.
- 3 P. H. Hünenberger and W. F. van Gunsteren, Empirical classical force fields for molecular systems, in *Lecture notes in Chemistry*, ed. A. F. Sax, Springer Verlag, Berlin, Germany, 1999, pp. 177–214.
- 4 D. A. MacKerell Jr, *J. Comput. Chem.*, 2004, **25**, 1584.
- 5 L. Monticelli and D. P. Tieleman, *Methods Mol. Biol.*, 2013, **924**, 197.
- 6 P. S. Nerenberg and T. Head-Gordon, *Curr. Opin. Struct. Biol.*, 2018, **49**, 129.
- 7 S. Riniker, *J. Chem. Inf. Model.*, 2018, **58**, 565.
- 8 M. P. Oliveira, M. Andrey, S. R. Rieder, L. Kern, D. F. Hahn, S. Riniker, B. A. C. Horta and P. H. Hünenberger, *J. Chem. Theory Comput.*, 2020, **16**, 7525.
- 9 W. L. Jorgensen, *J. Am. Chem. Soc.*, 1981, **103**, 335.
- 10 P. K. Weiner and P. A. Kollman, *J. Comput. Chem.*, 1981, **2**, 287.
- 11 B. R. Brooks, R. E. Bruccoleri, B. D. Olafson, D. J. States, S. Swaminathan and M. Karplus, *J. Comput. Chem.*, 1983, **4**, 187.
- 12 J. Hermans, H. J. C. Berendsen, W. F. van Gunsteren and J. P. M. Postma, *Biopolymers*, 1984, **23**, 1513.
- 13 W. F. van Gunsteren and H. J. C. Berendsen, *Groningen molecular simulation (GROMOS) library manual*, BIOMOS, Groningen, The Netherlands, 1987.
- 14 C. I. Bayly, P. Cieplak, W. D. Cornell and P. A. Kollman, *J. Phys. Chem.*, 1993, **97**, 10269.
- 15 Y. Duan, C. Wu, S. Chowdhury, M. C. Lee, G. Xiong, W. Yang Zhang, R. Cieplak, P. Luo, R. Lee, T. Caldwell, J. Wang and J. Kollman, *J. Comput. Chem.*, 2003, **24**, 1999.
- 16 C. M. Breneman and K. B. Wilberg, *J. Comput. Chem.*, 1990, **11**, 361.
- 17 D. E. Williams, *Rev. Comput. Chem.*, 1991, **2**, 219.
- 18 R. H. Henchman and J. W. Essex, *J. Comput. Chem.*, 1999, **20**, 483.
- 19 A. Jakalian, D. B. Jack and C. I. Bayly, *J. Comput. Chem.*, 2002, **23**, 1623.
- 20 C. Chipot, *J. Comput. Chem.*, 2003, **24**, 409.
- 21 M. Udier-Blagović, P. Morales de Tirado, S. A. Pearlman and W. L. Jorgensen, *J. Comput. Chem.*, 2004, **25**, 1322.
- 22 A. Stachowicz, A. Styrz and J. Korchowiec, *J. Mol. Model.*, 2011, **17**, 2217.
- 23 A. Stachowicz and J. Korchowiec, *Struct. Chem.*, 2012, **23**, 1449.
- 24 A. Ahmed and S. I. Sandler, *J. Chem. Theory Comput.*, 2013, **9**, 2774.
- 25 J. P. M. Jämbeck, F. Mocci, A. P. Lyubartsev and A. Laaksonen, *J. Comput. Chem.*, 2013, **34**, 187.



- 26 D. S. Cerutti, J. E. Rice, W. C. Swope and D. A. Case, *J. Phys. Chem. B*, 2013, **117**, 2328.
- 27 D. S. Cerutti, W. C. Swope, J. E. Rice and D. A. Case, *J. Chem. Theory Comput.*, 2014, **10**, 4515.
- 28 C. M. Ionescu, D. Sehnal, F. L. Falginella, P. Pant, L. Pravda, T. Bouchal, S. V. Vařeková, S. Geidl and J. Koča, *J. Cheminform.*, 2015, **7**, 50.
- 29 F. L. Hirshfeld, *Theor. Chim. Acta*, 1977, **44**, 129.
- 30 P. Bultinck, C. Van Alsenoy, P. W. Ayers and R. Carbó-Dorca, *J. Chem. Phys.*, 2007, **126**, 144111.
- 31 O. Yakovenko, A. A. Oliferenko, V. G. Bdzhola, V. A. Palyulin and N. S. Zefirov, *J. Comput. Chem.*, 2008, **29**, 1332.
- 32 T. A. Manz and D. S. Sholl, *J. Chem. Theory Comput.*, 2010, **6**, 2455.
- 33 N. M. Garrido, M. Jorge, A. J. Queimada, J. R. B. Gomes, I. G. Economou and E. A. Macedo, *Phys. Chem. Chem. Phys.*, 2011, **13**, 17384.
- 34 T. A. Manz and D. S. Sholl, *J. Chem. Theory Comput.*, 2012, **8**, 2844.
- 35 L. P. Lee, D. J. Cole, C.-K. Skylaris, W. L. Jorgensen and M. C. Payne, *J. Chem. Theory Comput.*, 2013, **9**, 2981.
- 36 L. P. Lee, N. Gabaldon Limas, D. J. Cole, M. C. Payne, C.-K. Skylaris and T. A. Manz, *J. Chem. Theory Comput.*, 2014, **10**, 5377.
- 37 T. Verstraelen, S. Vandenbrande, F. Heidar-Zadeh, L. Vanduyfhuys, V. Van Speybroeck, M. Waroquier and P. W. Ayers, *J. Chem. Theory Comput.*, 2016, **12**, 3894.
- 38 A. Peérez de la Luz, J. A. Aguilar-Pineda, J. G. Méndez-Bermúdez and J. Alejandro, *J. Chem. Theory Comput.*, 2018, **14**, 5949.
- 39 D. Reith and K. N. Kirschner, *Comput. Phys. Commun.*, 2011, **182**, 2184.
- 40 S. Grimme, *J. Chem. Theory Comput.*, 2014, **10**, 4497.
- 41 G. Prampolini, M. Campetella, N. De Mitri, P. R. Livotto and I. Cacelli, *J. Chem. Theory Comput.*, 2016, **12**, 5525.
- 42 J.-P. Piquemal and K. D. Jordan, *J. Chem. Phys.*, 2017, **147**, 161401.
- 43 P. Xu, E. B. Guidez, C. Bertoni and M. S. Gordon, *J. Chem. Phys.*, 2018, **148**, 090901.
- 44 J. T. Horton, A. E. A. Allen, L. S. Dodda and D. J. Cole, *J. Chem. Inf. Model.*, 2019, **59**, 1366.
- 45 A. E. A. Allen, M. J. Robertson, M. C. Payne and D. J. Cole, *ACS Omega*, 2019, **4**, 14537.
- 46 S. M. Kantonen, H. S. Muddana, M. Schauerperl, N. M. Henriksen and L. Wang, *J. Chem. Theory Comput.*, 2020, **16**, 1115.
- 47 X. Chu and A. Dalgarno, *J. Chem. Phys.*, 2004, **121**, 4083.
- 48 A. Olasz, K. Vanommeslaeghe, A. Krishtal, T. Veszprémi, C. Van Alsenoy and P. Geerlings, *J. Chem. Phys.*, 2007, **127**, 224105.
- 49 A. J. Stone and A. J. Misquitta, *Int. Rev. Phys. Chem.*, 2007, **26**, 193.
- 50 A. J. Stone, *Science*, 2008, **321**, 787.
- 51 A. Tkatchenko and M. Scheffler, *Phys. Rev. Lett.*, 2009, **102**, 073005.
- 52 A. Tkatchenko, R. A. DiStasio Jr, R. Car and M. Scheffler, *Phys. Rev. Lett.*, 2012, **108**, 236402.
- 53 V. V. Gobre and A. Tkatchenko, *Nat. Commun.*, 2013, **4**, 2341.
- 54 T. Bučko, S. Lebègue, J. Hafner and J. G. Ángyán, *J. Chem. Theory Comput.*, 2013, **9**, 4293.
- 55 T. Bučko, S. Lebègue, J. G. Ángyán and J. Hafner, *J. Chem. Phys.*, 2014, **141**, 034114.
- 56 D. J. Cole, J. Z. Vilseck, J. Tirado-Rives, M. C. Payne and W. L. Jorgensen, *J. Chem. Theory Comput.*, 2016, **12**, 2312.
- 57 M. Mohebifar, E. R. Johnson and C. N. Rowley, *J. Chem. Theory Comput.*, 2017, **13**, 6146.
- 58 C.-K. Skylaris, P. D. Haynes, A. A. Mostofi and M. C. Payne, *J. Chem. Phys.*, 2005, **122**, 084119.
- 59 I. S. Ufimtsev, N. Luehr and T. J. Martinez, *J. Phys. Chem. Lett.*, 2011, **2**, 1789.
- 60 D. R. Bowler and T. Miyazaki, *Rep. Prog. Phys.*, 2012, **75**, 036503.
- 61 J. Dziedzic, S. J. Fox, T. Fox, C. S. Tautermann and C.-K. Skylaris, *Int. J. Quantum Chem.*, 2013, **113**, 771.
- 62 K. Wilkinson and C.-K. Skylaris, *J. Comput. Chem.*, 2013, **34**, 2446.
- 63 K. A. Wilkinson, N. D. M. Hine and C.-K. Skylaris, *J. Chem. Theory Comput.*, 2014, **10**, 4782.
- 64 G. Lever, D. J. Cole, R. Lonsdale, K. E. Ranaghan, D. J. Wales, A. J. Mulholland, C.-K. Skylaris and M. C. Payne, *J. Phys. Chem. Lett.*, 2014, **5**, 3614.
- 65 C. A. Renison, K. D. Fernandes and K. J. Naidoo, *J. Comput. Chem.*, 2015, **36**, 1410.
- 66 L. Ruddigkeit, R. van Deursen, L. C. Blum and J.-L. Reymond, *J. Chem. Inf. Model.*, 2012, **52**, 2864.
- 67 J.-L. Reymond, *Acc. Chem. Res.*, 2015, **48**, 722.
- 68 W. L. Jorgensen, K. P. Jensen and A. N. Alexandrova, *J. Chem. Theory Comput.*, 2007, **3**, 1987.
- 69 A. J. Stone, *J. Chem. Theory Comput.*, 2005, **1**, 1128.
- 70 D. L. Mobley, E. Dumont, J. D. Chodera and K. A. Dill, *J. Phys. Chem. B*, 2007, **111**, 2242.
- 71 P. Bultinck, P. W. Ayers, S. Fias, K. Tiels and C. Van Alsenoy, *Chem. Phys. Lett.*, 2007, **444**, 205.
- 72 O. Beckstein and B. I. Iorga, *J. Comput.-Aided Mol. Des.*, 2012, **26**, 635.
- 73 O. Beckstein, A. Fourrier and B. I. Iorga, *J. Comput.-Aided Mol. Des.*, 2014, **28**, 265.
- 74 J. Z. Vilseck, J. Tirado-Rives and W. L. Jorgensen, *J. Chem. Theory Comput.*, 2014, **10**, 2802.
- 75 L. S. Dodda, J. Z. Vilseck, K. J. Cutrona and W. L. Jorgensen, *J. Chem. Theory Comput.*, 2015, **11**, 4273.
- 76 E. Boulanger, L. Huang, C. Rupakheti, A. D. MacKerell Jr. and B. Roux, *J. Chem. Theory Comput.*, 2018, **14**, 3121.
- 77 K. M. Visscher and D. P. Geerke, *J. Chem. Theory Comput.*, 2019, **15**, 1875.
- 78 Y. M. H. Gonçalves, C. Senac, P. F. J. Fuchs, P. H. Hünenberger and B. A. C. Horta, *J. Chem. Theory Comput.*, 2019, **15**, 1806.
- 79 J. R. Ullmann, *J. ACM*, 1976, **23**, 31.
- 80 D. Weininger, A. Weininger and J. L. Weininger, *J. Chem. Inf. Comput. Sci.*, 1989, **29**, 97.



- 81 R. Grund, *Konstruktion molekularer Graphen mit gegebenen Hybridisierungen und überlappungsfreien Fragmenten*, Lehrstuhl II für Mathematik der Universität Bayreuth, 1994.
- 82 N. Schneider, R. A. Sayle and G. A. Landrum, *J. Chem. Inf. Model.*, 2015, **55**, 2111.
- 83 J. Wang, W. Wang, P. A. Kollman and D. A. Case, *J. Mol. Graphics Modell.*, 2006, **25**, 247.
- 84 K. Vanommeslaeghe and A. D. MacKerell Jr, *J. Chem. Inf. Model.*, 2012, **52**, 3144.
- 85 K. Vanommeslaeghe, E. P. Raman and A. D. MacKerell Jr, *J. Chem. Inf. Model.*, 2012, **52**, 3155.
- 86 A. K. Malde, L. Zuo, M. Breeze, M. Stroet, D. Poger, P. C. Nair, C. Oostenbrink and A. E. Mark, *J. Chem. Theory Comput.*, 2011, **7**, 4026.
- 87 K. B. Koziara, M. Stroet, A. K. Malde and A. E. Mark, *J. Comput.-Aided Mol. Des.*, 2014, **28**, 221.
- 88 M. Stroet, B. Caron, K. M. Visscher, D. P. Geerke, A. K. Malde and A. E. Mark, *J. Chem. Theory Comput.*, 2018, **14**, 5834.
- 89 E. Krieger, G. Koraimann and G. Vriend, *Proteins: Struct., Funct., Genet.*, 2002, **47**, 393.
- 90 A. W. Schüttelkopf and D. M. F. van Aalten, *Acta Crystallogr.*, 2004, **D60**, 1355.
- 91 A. A. S. T. Ribeiro, B. A. C. Horta and R. B. de Alencastro, *J. Braz. Chem. Soc.*, 2008, **19**, 1433.
- 92 V. Zoete, M. A. Cuendet, A. Grosdidier and O. Michielin, *J. Comput. Chem.*, 2011, **32**, 2359.
- 93 C. Margreitter, D. Petrov and B. Zagrovic, *Nucleic Acids Res.*, 2013, **41**, W422.
- 94 S. Jo, X. Cheng, S. M. Islam, L. Huang, H. Rui, A. Zhu, H. S. Lee, Y. Qi, W. Han, K. Vanommeslaeghe, A. D. MacKerell Jr, B. Roux and W. Im, *Adv. Protein Chem. Struct. Biol.*, 2014, **96**, 235.
- 95 Y. Pevzner, E. Frugier, V. Schalk, A. Caffisch and H. L. Woodcock, *J. Chem. Inf. Model.*, 2014, **54**, 2612.
- 96 M. Di Pierro and R. Elber, *J. Chem. Theory Comput.*, 2013, **9**, 3311.
- 97 M. Di Pierro, M. L. Mugnai and R. Elber, *J. Phys. Chem. B*, 2015, **119**, 836.
- 98 L.-P. Wang, J. Chen and T. van Voorhis, *J. Chem. Theory Comput.*, 2013, **9**, 452.
- 99 L.-P. Wang, T. Head-Gordon, J. W. Ponder, P. Ren, J. D. Chodera, P. K. Eastman, T. J. Martinez and V. S. Pande, *J. Phys. Chem. B*, 2013, **117**, 9956.
- 100 L.-P. Wang, T. J. Martinez and V. S. Pande, *J. Phys. Chem. Lett.*, 2014, **5**, 1885.
- 101 R. Qi, L.-P. Wang, Q. Wang, V. S. Pande and P. Ren, *J. Chem. Phys.*, 2015, **143**, 014504.
- 102 M. L. Laury, L.-P. Wang, V. S. Pande, T. Head-Gordon and J. W. Ponder, *J. Phys. Chem. B*, 2015, **119**, 9423.
- 103 K. A. McKiernan, L.-P. Wang and V. S. Pande, *J. Chem. Theory Comput.*, 2016, **12**, 5960.
- 104 A. D. Wade, L.-P. Wang and D. J. Huggins, *J. Chem. Inf. Model.*, 2018, **58**, 1766.
- 105 Y. Qiu, P. S. Nerenberg, T. Head-Gordon and L.-P. Wang, *J. Phys. Chem. B*, 2019, **123**, 7061.
- 106 J. Yin, A. T. Fenley, N. M. Henriksen and M. K. Gilson, *J. Phys. Chem. B*, 2015, **119**, 10145.
- 107 J. Yin, N. M. Henriksen, H. S. Muddana and M. K. Gilson, *J. Chem. Theory Comput.*, 2018, **14**, 3621.
- 108 L. N. Naden and M. R. Shirts, *J. Chem. Theory Comput.*, 2015, **12**, 1806.
- 109 M. Stroet, K. B. Koziara, A. K. Malde and A. E. Mark, *J. Chem. Theory Comput.*, 2017, **13**, 6201.
- 110 R. A. Messerly, S. M. Razavi and M. R. Shirts, *J. Chem. Theory Comput.*, 2018, **14**, 3144.
- 111 R. A. Messerly, M. S. Barhaghi, J. J. Potoff and M. R. Shirts, *J. Chem. Eng. Data*, 2019, **64**, 3701.
- 112 B. A. C. Horta, P. T. Merz, P. Fuchs, J. Dolenc, S. Riniker and P. H. Hünenberger, *J. Chem. Theory Comput.*, 2016, **12**, 3825.
- 113 T. Verstraelen, V. Van Speybroeck and M. Waroquier, *J. Chem. Phys.*, 2009, **131**, 044127.
- 114 W. Acree Jr and J. S. Chickos, *J. Phys. Chem. Ref. Data*, 2016, **45**, 033101.
- 115 M. Frenkel, X. Hong, R. C. Wilhoit and K. R. Hall, in *Thermodynamic properties of organic compounds and their mixtures. Densities of alcohols*, ed. K. R. Hall and K. N. Marsh, Landolt-Börnstein Series, Springer-Verlag, Berlin/Heidelberg, Deutschland, 2000, vol. IV/8G.
- 116 M. Frenkel, X. Hong, R. C. Wilhoit and K. R. Hall, in *Thermodynamic properties of organic compounds and their mixtures. Densities of esters and ethers*, ed. K. R. Hall and K. N. Marsh, Landolt-Börnstein Series, Springer-Verlag, Berlin/Heidelberg, Deutschland, 2001, vol. IV/8H.
- 117 M. Frenkel, X. Hong, Q. Dong, X. Yan and R. D. Chirico, in *Thermodynamic properties of organic compounds and their mixtures. Densities of phenols, aldehydes, ketones, carboxylic acids, amines, nitriles, and nitrohydrocarbons*, ed. K. R. Hall and K. N. Marsh, Landolt-Börnstein Series, Springer-Verlag/Springer-Verlag, Berlin/Heidelberg, Deutschland, 2002, vol. IV/8I.
- 118 M. Frenkel, R. D. Chirico, V. Diky, Q. Dong, K. N. Marsh, J. H. Dymond, W. A. Wakeham, S. E. Stein, E. Königsberger and A. R. H. Goodwin, *Pure Appl. Chem.*, 2006, **78**, 541.
- 119 C. Wohlfahrt, *Static dielectric constants of pure liquids and binary liquid mixtures*, Springer, Berlin, Germany, 2008, vol. IV/17.
- 120 J. R. Rumble, *CRC Handbook of Chemistry and Physics*, CRC Press/Taylor and Francis, Boca Raton, USA, 98th edn, 2018.
- 121 C. L. Yaws, *Thermophysical properties of chemicals and hydrocarbons*, Gulf Professional Publishing (Elsevier), Oxford, UK, 2nd edn, 2014.
- 122 Springer Nature, Springer Materials database. Available at: <https://materials.springer.com>, 2018.
- 123 D. Weininger, *J. Chem. Inf. Comput. Sci.*, 1988, **28**, 31.
- 124 S. Kim, J. Chen, T. Cheng, A. Gindulyte, J. He, S. He, Q. Li, B. A. Shoemaker, P. A. Thiessen, B. Yu, L. Zaslavsky, J. Zhang and E. E. Bolton, *Nucleic Acids Res.*, 2019, **47**, D1102.
- 125 P. H. Hünenberger, CombiFF Data Collection in the ETHZ Research Collection (tar-file CombiFF_saturated_O_



- and_N_compounds, version 1.0 corresponds to the published article), 2021, DOI: 10.3929/ethz-b-000445271.
- 126 W. F. van Gunsteren, S. R. Billeter, A. A. Eising, P. H. Hünenberger, P. Krüger, A. E. Mark, W. R. P. Scott and I. G. Tironi, *Biomolecular simulation: The GROMOS96 manual and user guide*, Verlag der Fachvereine, Zürich, Switzerland, 1996.
- 127 W. F. van Gunsteren, X. Daura and A. E. Mark, GROMOS force field, in *Encyclopedia of computational chemistry*, ed. P. Schleyer, John Wiley & Sons, Chichester, UK, 1998, vol. 2, pp. 1211–1216.
- 128 W. R. P. Scott, P. H. Hünenberger, I. G. Tironi, A. E. Mark, S. R. Billeter, J. Fennen, A. E. Torda, T. Huber, P. Krüger and W. F. van Gunsteren, *J. Phys. Chem. A*, 1999, **103**, 3596.
- 129 C. Oostenbrink, A. Villa, A. E. Mark and W. F. van Gunsteren, *J. Comput. Chem.*, 2004, **25**, 1656.
- 130 M. Christen, P. H. Hünenberger, D. Bakowies, R. Baron, R. Bürgi, D. P. Geerke, T. N. Heinz, M. A. Kastenholz, V. Kräutler, C. Oostenbrink, C. Peter, D. Trzesniak and W. F. van Gunsteren, *J. Comput. Chem.*, 2005, **26**, 1719.
- 131 W. F. van Gunsteren, The GROMOS software for biomolecular simulation. Available at: <http://www.gromos.net>, 05/05/2011.
- 132 A. T. Hagler, E. Huler and S. Lifson, *J. Am. Chem. Soc.*, 1974, **96**, 5319.
- 133 S. Lifson, A. T. Hagler and P. Dauber, *J. Am. Chem. Soc.*, 1979, **101**, 5111.
- 134 J. E. Lennard-Jones, *Physica*, 1937, **4**, 941.
- 135 X. Daura, A. E. Mark and W. F. van Gunsteren, *J. Comput. Chem.*, 1998, **19**, 535.
- 136 L. D. Schuler and W. F. van Gunsteren, *Mol. Simul.*, 2000, **25**, 301.
- 137 L. D. Schuler, X. Daura and W. F. van Gunsteren, *J. Comput. Chem.*, 2001, **22**, 1205.
- 138 E. Bourasseau, M. Haboudou, A. Boutin, A. H. Fuchs and P. Ungerer, *J. Chem. Phys.*, 2003, **118**, 3020.
- 139 W. G. Hoover, *Phys. Rev. A: At., Mol., Opt. Phys.*, 1985, **31**, 1695.
- 140 H. C. Andersen, *J. Chem. Phys.*, 1980, **72**, 2384.
- 141 W. F. van Gunsteren, H. J. C. Berendsen and J. A. C. Rullmann, *Mol. Phys.*, 1981, **44**, 69.
- 142 E. Guàrdia and J. A. Padró, *J. Chem. Phys.*, 1985, **83**, 1917.
- 143 W. F. van Gunsteren and H. J. C. Berendsen, *Mol. Simul.*, 1988, **1**, 173.
- 144 S. Yun-yu, W. Lu and W. F. van Gunsteren, *Mol. Simul.*, 1988, **1**, 369.
- 145 W. F. van Gunsteren, Molecular dynamics and stochastic dynamics simulation: a primer, in *Computer simulation of biomolecular systems, theoretical and experimental applications*, ed. W. F. van Gunsteren, P. K. Weiner, A. J. Wilkinson, ESCOM Science Publishers, B.V., Leiden, The Netherlands, 1993, vol. 2, pp. 3–36.
- 146 R. W. Hockney, *Methods Comput. Phys.*, 1970, **9**, 135.
- 147 J.-P. Ryckaert, G. Ciccotti and H. J. C. Berendsen, *J. Comput. Phys.*, 1977, **23**, 327.
- 148 H. J. C. Berendsen, W. F. van Gunsteren, H. R. J. Zwinderman and R. G. Geurtsen, *Ann. N. Y. Acad. Sci.*, 1986, **482**, 269.
- 149 J. A. Barker and R. O. Watts, *Mol. Phys.*, 1973, **26**, 789.
- 150 I. G. Tironi, R. Sperb, P. E. Smith and W. F. van Gunsteren, *J. Chem. Phys.*, 1995, **102**, 5451.
- 151 T. Verstraelen, P. Bultinck, V. Van Speybroeck, P. W. Ayers, D. Van Neck and M. Waroquier, *J. Chem. Theory Comput.*, 2011, **7**, 1750.
- 152 A. Kubincová, P. H. Hünenberger and M. Krishnan, *J. Chem. Phys.*, 2020, **152**, 104713.
- 153 A. K. Rappé and W. A. Goddard, *J. Phys. Chem.*, 1991, **95**, 3358.
- 154 B. A. Wells, C. De Bruin-Dickason and A. L. Chaffee, *J. Phys. Chem. C*, 2015, **119**, 456.
- 155 S. W. Rick, S. J. Stuart and B. J. Berne, *J. Chem. Phys.*, 1994, **101**, 6141.
- 156 S. W. Rick and B. J. Berne, *J. Am. Chem. Soc.*, 1996, **118**, 672.
- 157 S. W. Rick, *J. Chem. Phys.*, 2001, **114**, 2276.
- 158 J. L. Banks, G. A. Kaminski, R. H. Zhou, D. T. Mainz, B. J. Berne and R. A. Friesner, *J. Chem. Phys.*, 1999, **110**, 741.
- 159 H. A. Stern, G. A. Kaminski, J. L. Banks, R. Zhou, B. J. Berne and R. A. Friesner, *J. Phys. Chem. B*, 1999, **103**, 4730.
- 160 S. Patel and C. L. Brooks III, *J. Comput. Chem.*, 2004, **25**, 1.
- 161 S. Patel, A. D. MacKerell Jr and C. L. Brooks III, *J. Comput. Chem.*, 2004, **25**, 1504.
- 162 S. Patel and C. L. Brooks III, *Mol. Simul.*, 2006, **32**, 231.
- 163 Z.-Z. Yang, J.-J. Wang and D.-X. Zhao, *J. Comput. Chem.*, 2014, **35**, 1690.

