

Cite this: *J. Mater. Chem. A*, 2020, **8**, 9850Unexpected trends in the enhanced Ce³⁺ surface concentration in ceria–zirconia catalyst materials†Weizi Yuan,^a Qing Ma,^b Yangang Liang,^c Chengjun Sun,^d K. V. L. V. Narayanachari,^a Michael J. Bedzyk,^a Ichiro Takeuchi^c and Sossina M. Haile^{*,a}

Despite the immense importance of ceria–zirconia solid solutions in heterogeneous catalysis, and the growing consensus that catalytic activity correlates with the concentration of reduced Ce³⁺ species and accompanying oxygen vacancies, the extent of reduction at the surfaces of these materials, where catalysis occurs, is unknown. Using angle-resolved X-ray Absorption Near Edge Spectroscopy (XANES), we quantify under technologically relevant conditions the Ce³⁺ concentration in the surface (2–3 nm) and bulk regions of ceria–zirconia films grown on single crystal yttria-stabilized zirconia, YSZ (001). In all circumstances, we observe substantial Ce³⁺ enrichment at the surface relative to the bulk. Surprisingly, the degree of enhancement is highest in the absence of Zr. This behavior stands in direct contrast to that of the bulk in which the Ce³⁺ concentration monotonically increases with increasing Zr content. These results suggest that while Zr enhances the oxygen storage capacity in ceria, undoped ceria may have higher surface catalytic activity. They further urge caution in the use of bulk properties as surrogate descriptors for surface characteristics and hence catalytic activity.

Received 9th March 2020

Accepted 24th April 2020

DOI: 10.1039/d0ta02762f

rsc.li/materials-a

Introduction

Ceria (CeO_{2-δ}) and ceria–zirconia (Ce_{1-x}Zr_xO_{2-δ}) solid solutions are important catalytic materials, either serving directly as catalysts themselves, or serving as supports for metal nanoparticle catalysts.^{1–4} The ease with which Ce undergoes changes in oxidation state is believed to account for the buffering capacity of ceria, which in turn is associated with its high catalytic activity.⁵ Introduction of zirconium increases the reducibility of the material,^{2,6–9} as has been quantified in recent measurements which reveal the relationship between bulk oxidation state, bulk vacancy concentration and environmental conditions.^{10–12} This behavior is quite distinct from that of rare-earth doping of ceria, which largely leaves the bulk thermodynamic properties unchanged from those of undoped ceria.¹³ Arguably, however, the redox properties of the *surfaces* are likely to be more important to the catalytic characteristics than those of the bulk. Here, quantitative data are scarce, particularly with respect to *operando* (experimental) observations.

Computational studies of the ceria–zirconia system suggest enhanced surface vacancy concentrations relative to the bulk.^{4,5,14–20} In particular, Balducci *et al.* have predicted that oxygen vacancy segregation to the surface increases with Zr content and is more pronounced in the sequence (310) > (110) > (111), in accord with the sequence of stability of these surfaces.¹⁵ In a series of studies of CeO₂ and Ce_{0.75}Zr_{0.25}O₂, Yang *et al.*^{17–19} found (110) and (111) surface vacancy formation energies in all cases to be lower than those of the bulk, but in contrast to Balducci,¹⁵ they observed no trend in this differential with chemistry. Experimentally, direct detection of oxygen vacancies is challenging, particularly in comparison to the detection of Ce³⁺. For reasons of electroneutrality, the concentrations of these two types of species are expected to be correlated²¹ (although other charge compensation modes cannot be entirely ruled out), and thus most experimental work, including the present study, focuses on the detection of reduced trivalent cerium. A number of experimental investigations have shown that the surfaces of undoped^{22–26} or rare-earth doped²⁷ ceria are more reduced than the bulk, in agreement with computational results. The sole study focused on ceria–zirconia appears to be the work of Zhao *et al.*²⁸ These authors inferred the Ce³⁺ concentration on CeO₂ and Ce_{0.5}Zr_{0.5}O₂ from microkinetic analysis of the measured rate of CO oxidation over reduced powders of the oxides. The results indicated that the vacancy concentration on ceria is almost twice that on Ce_{0.5}Zr_{0.5}O₂, suggesting that Zr suppresses the vacancy formation on the oxide surface, in contrast to its vacancy-enhancing influence on

^aDepartment of Materials Science and Engineering, Northwestern University, Illinois 60208, USA. E-mail: sossina.haile@northwestern.edu

^bDND-CAT, Northwestern Synchrotron Research Center at the Advanced Photon Source, Argonne, Illinois 60439, USA

^cDepartment of Materials Science and Engineering, University of Maryland, College Park, Maryland 20742, USA

^dAdvanced Photon Source, Argonne National Laboratory, Argonne, Illinois 60439, USA

† Electronic supplementary information (ESI) available. See DOI: 10.1039/d0ta02762f

the bulk. Beyond this work, limited to the comparison of two specific compositions and without control of surface termination, no other studies have been directed towards the critically important ceria–zirconia system. The present work addresses this gap. Using angle-resolved X-ray Absorption Near Edge Spectroscopy (XANES), we probe both the surface and bulk regions of thin-film ceria–zirconia solid solutions. We show that the surfaces are indeed far more reduced than the bulk, but that contrary to the bulk (and in agreement with the proposal by Zhao for $\text{Ce}_{0.5}\text{Zr}_{0.5}\text{O}_2$), the extent of reduction on the surface decreases with increasing Zr content.

We study here (001)-oriented $\text{Ce}_{1-x}\text{Zr}_x\text{O}_{2-\delta}$ films (hereafter CZO $xx\%$, where xx is the Zr content), obtained by growth on yttria-stabilized zirconia (YSZ) of the same orientation. The (001) surface is selected because of the high activity and high surface vacancy concentration of this face relative to the more stable (110) and (111) faces.^{3,29–32} XANES is selected as the method of characterization because it benefits from facile compatibility with environmental chambers for *in situ* experiments and because the X-ray penetration depth can be readily varied by changing the angle of incidence, permitting access to both near-surface and bulk regions of a material. Specifically, by fixing the angle of incidence of the incoming X-rays to a value below the critical angle at which total reflection occurs, sensitivity is limited to the top 2–3 nanometers of the material, while use of larger X-ray incident angles provides access to the entire film. Furthermore, quantification of the relative Ce^{3+} concentration is generally possible by simple linear combination of the spectra of reference materials³³ in which the Ce is either fully in the 3+ or 4+ oxidation state. The negligible role of substrate–film interactions on surface redox chemistry is established here by preparing and characterizing a (001)-oriented film of CZO28 on *r*-sapphire. As shown below, the surface characteristics of this sample were indistinguishable from those of CZO28 grown on YSZ. In addition to the ceria–zirconia films, a film of CeAlO_3 was measured to generate XANES Ce^{3+} reference spectra. This perovskite-structured oxide presents Ce^{3+} in a 12-fold oxygen coordinated environment, reasonably similar to that of the 8-fold coordination in the fluorite structure.

Film growth

Films, either with a uniform fixed composition ($x = 0, 0.15, 0.28$ and 0.48) or with a lateral gradient from $x = 0.0425$ to 0.48 , were grown by pulsed laser deposition (PLD). The composition space examined spans the solubility range of Zr into ceria (known to extend to about 50 mol% $\text{Zr}^{34,35}$). The film thickness was fixed at 220 nm such that growth strain at the film surface was fully relaxed,³⁶ yet self-absorption effects in the XANES measurements, which can become severe in thick samples, were negligible.³⁷ To minimize the number of targets required, in most cases films were grown from the end-members CeO_2 and CZO48, with intermediate compositions obtained using an alternating, monolayer-by-monolayer approach.³⁸ The CeO_2 target was synthesized from commercial powder (99.9% purity, Sigma Aldrich #211575). The powder was compacted under uniaxial pressure and sintered at 1500 °C for 10 h in still air.

CZO48 and CZO28 powders were prepared by a wet chemical route. Stoichiometric amounts of $\text{Ce}(\text{NO}_3)_3 \cdot 5.90\text{H}_2\text{O}$ (Alfa Aesar, 99.5% purity) and $\text{ZrO}(\text{NO}_3)_2 \cdot 5.98\text{H}_2\text{O}$ (Alfa Aesar, 99.9% purity) were dissolved in dilute HNO_3 and then added to distilled water. Ethylenediaminetetraacetic acid (EDTA) and citric acid were subsequently added to the solution as chelating agents. Ammonium hydroxide was then introduced dropwise to bring the pH of the mixture to ~ 10 . The solution was heated to 80 °C and held at this temperature under constant stirring until the liquid content was significantly reduced and a thick, reddish gel remained. The products were briefly heat treated at 350 °C in air, and then fully calcined at 700 °C for 8 h. Sintering was carried out by heat treatment of pressed compacts at 1500 °C for 8 h in still air. The intended compositions of the CZO materials were 25 and 45 mol% ZrO_2 ; the quoted compositions of 28 and 48% are based on post-synthesis electron dispersive X-ray spectroscopy (EDS) and X-ray fluorescence (XRF) chemical analyses.

Films of $\text{CeO}_{2-\delta}$ and CZO28 were grown using a PVD PLD/MBE 2300 instrument (KrF 248 nm excimer laser, 1.75 J cm^{-2} laser energy, 10 Hz repetition rate). The temperature of the substrate, single crystal $\text{Y}_{0.16}\text{Zr}_{0.84}\text{O}_{1.92}$ (YSZ) (001) (MTI Corp.) or *r*-cut Al_2O_3 (MTI Corp.), was fixed at 650 °C in the growth step, and the atmosphere was fixed at 30 mTorr of oxygen. At the conclusion of the deposition, the film was cooled at a rate of 10 °C min^{-1} . Compositionally graded films, as well as films of CZO15 and CZO48, were grown using a custom-built PLD/Laser-MBE System (Pascal Co., Ltd.) equipped with a KrF 248 nm excimer laser (Lambda COMPex Pro, 0.51 J cm^{-2} laser energy, 5 Hz repetition rate) and the capabilities for monolayer-by-monolayer compositional control.³⁸ During growth, the temperature of the $\text{Y}_{0.16}\text{Zr}_{0.84}\text{O}_{1.92}$ (YSZ) (001) single crystal substrates (MTI Corp.) was fixed at 640 °C, and the atmosphere was again fixed at 30 mTorr of oxygen. After deposition was complete, the films were cooled at a rate of 5 °C min^{-1} .

A 490 nm film of CeAlO_3 was grown using a target prepared from powders synthesized by solid state reaction as follows. Stoichiometric amounts of CeO_2 (99.9% purity, Sigma Aldrich) and Al_2O_3 (99.98% purity, Alfa Aesar CAS# 1334-28-1) were mixed by ball milling for 2 days. The resulting powder was heat-treated at 950 °C in still air for 12 h. Sintering was carried out at 1550 °C under flowing 3% H_2/Ar for 50 h. The film growth conditions were identical to those employed for CeO_2 and CZO28 with the exception of the use of vacuum as the atmosphere. After deposition and removal from the chamber, the CeAlO_3 film was annealed in flowing 3% H_2/Ar at 1600 °C for 12 h to ensure that the Ce was fully reduced, as subsequently verified by X-ray diffraction (no other phases were present) and X-ray photoelectron spectroscopy (XPS) about the Ce3d region.

The compositionally graded film was patterned into strips to facilitate the positioning of the X-ray beam for XANES measurements. The pattern of 19 strips (240 μm in width and separated by a spacing of 160 μm) was created using photolithography and ion milling. Specifically, the sample was coated with a photoresist (Shipley 1813) by spin coating (4000 rpm for 50 s), which was then baked at 100 °C for 2 min to drive off solvents and to solidify the film. Following exposure to UV

radiation for 12 s through a photomask, the photoresist was developed (Shipley 352) for 40 s. The sample then underwent ion milling for 40 min, removing 230 nm of oxide.

Structural and chemical characterization

X-ray diffraction (Rigaku ATXG, Cu K_{α}) was used to confirm phase formation of target materials (not shown) and to characterize the crystallographic features of the films. Specular θ - 2θ scans, azimuthal ϕ scans, and rocking curve measurements of the non-graded films were collected using this same instrument (Rigaku ATXG, Cu K_{α}). Micro-area X-ray diffraction patterns of the graded sample were measured using a Rigaku SmartLab diffractometer (Cu K_{α}) equipped with polycapillary focusing optics (CBO-f unit). The X-ray line source was focused to a spot 400 μm in diameter with an angular divergence of 0.4° . The composition gradient was aligned parallel to the goniometer axis to minimize the range of compositions detected. A Ni filter was placed in front of the detector to absorb Cu K_{β} radiation. Atomic force micrographs were recorded using a Bruker Dimension FastScan in tapping mode. The surface roughness of the observed features in the atomic force microscopy (AFM) images was extracted by post image analysis using Nanoscope Analysis software.³⁹

As is typical of ceria films on YSZ,^{40,41} all ceria-zirconia films prepared here were found to grow epitaxially. The specular θ - 2θ scans of the non-graded films, Fig. S1(a),[†] contained only the (00L) peaks permitted of a FCC lattice, whereas the ϕ azimuthal scans, collected for the non-graded films about the surface normal c axis with in-plane 2θ fixed at the (200) geometry, Fig. S1(b),[†] showed the expected four-fold symmetry. The lattice parameter mismatch between the films and the substrate results in measurable broadening in the rocking curves, Fig. S1(c and d),[†] particularly at low zirconium content. Nevertheless, the in-plane domain sizes are large, Table S1,[†] and the root-mean-square surface roughness, as measured by AFM, remained below 0.5 nm (Fig. S2[†]). In the compositionally graded film,

a shift in peak position to higher angle with increasing Zr content was observed, Fig. 1, reflecting the incorporation of the smaller zirconium ions into the CeO₂ structure. The positions of the (00L) peaks and the correlation between the lattice parameter and composition reported by Laguna *et al.*⁴² were used to infer the composition profile.

The chemical composition along the film thickness of the non-graded films was measured by angle-resolved XRF using synchrotron radiation at an excitation energy of 20480 eV, which is above the Zr K edge. Data were collected at the DuPont-Northwestern-Dow Collaborative Access Team (DND-CAT) 5BM-D station at the Advanced Photon Source (APS), Argonne National Laboratory, using an X-ray beam with a 15 μm (vertical) by 4 mm (horizontal) incident beam size. A Vortex-ME4 silicon drift detector was used to collect the fluorescence signal, where the critical angle was found to be 0.14° . The exit angle was fixed at $\approx 22^{\circ}$ and the incident angle varied between 0.11 and 0.175° , providing access to film depths in the range from ≈ 3 to ≈ 80 nm. The penetration depth changes somewhat with the dopant level for a given incident angle, but the measurement nevertheless directly provides depth-resolved composition information for each film. By limiting the maximum penetration depth to ≈ 85 nm, attenuation of the fluorescence signal as it travels out of the film due to self-absorption was avoided. In all cases, the film composition was uniform along the thickness direction, Fig. S3,[†] with average compositions of 14.5 ± 0.3 , 28.4 ± 1.1 and 47.5 ± 0.8 mol% ZrO₂, respectively, for films of target compositions 15, 28 and 48%.

Methods for oxidation state determination by XANES

Using these fully characterized films, we carried out XANES measurements at the Ce L₃ edge (2p-to-5d transition) in the fluorescence mode under five sets of conditions: in ambient air at room temperature, 800 $^{\circ}\text{C}$, and 1100 $^{\circ}\text{C}$; upon exposure to vacuum at 1100 $^{\circ}\text{C}$; and upon exposure to hydrogen at 800 $^{\circ}\text{C}$, as

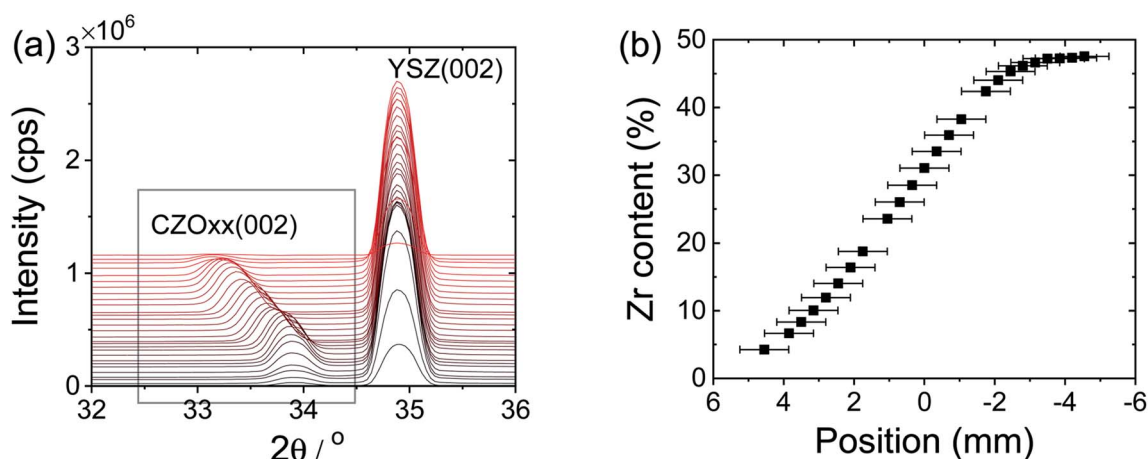


Fig. 1 X-ray diffraction study of graded ceria-zirconia film deposited on YSZ (001): (a) specular θ - 2θ scan, zoomed in to the region in the vicinity of the (002) peak; and (b) composition inferred from the refined lattice parameter.

Table 1 Experimental conditions and sample types for XANES measurements

Temp/atm (pO_2)	160 Torr (air)	~ 2.3 mTorr (vacuum)	7.4×10^{-14} mTorr ("hydrogen")
25 °C	Graded ^a , non-graded ^a		
800 °C	Graded ^a , CeO _{2-δ} non-graded ^a		Graded ^b , CeO _{2-δ} non-graded ^a
1100 °C	Non-graded ^a	Non-graded ^b	

^a Negligible Ce³⁺ in the bulk, at all compositions. ^b Increasing Ce³⁺ content in the bulk with increasing Zr content.¹⁰

summarized in Table 1. Data from the non-graded films were collected at the DND-CAT 5BM-D station at the APS using a Si (111) double crystal monochromator with an energy resolution of $\Delta E/E = 1.4 \times 10^{-4}$. The incident X-ray beam size was 0.05 (vertical) \times 5 (horizontal) mm². The graded film was measured at the undulator beamline 20-ID-C at the APS using linear polarized X-rays and a Si (111) monochromator with resolution $\Delta E/E = 1.3 \times 10^{-4}$. The full width at half maximum (FWHM) of the X-ray beam was ~ 6 μ m in both horizontal and vertical directions, focused using KB mirrors. Details of the beamline optics and instrumentation can be found elsewhere.⁴³ In both sets of experiments, the incident X-ray beam intensity was monitored using an ion chamber. For measurements at 5BM-D, the intensity was detuned to 60% of the maximum. The fluorescence signal was measured using Vortex-ME4 silicon drift detectors, set at 90° relative to the incident X-ray beam direction (2 detectors at 5BM-D, 1 detector at 20-ID-C). The X-ray incidence angle α was aligned to a precision better than 0.03° using a set of slits (0.05 \times 12 mm²) placed behind the sample. Positioning of the 19 CZO strips of the graded film was accomplished by setting the incident angle to $\sim 1^\circ$ and the incident beam energy to 6000 eV, and moving the stage along the gradient direction. XANES spectra, both surface and bulk, were then collected sequentially at each position of interest before moving to the next. The intensity of the Ce fluorescence signal was used to establish the beam position (Fig. S4†). Due to time constraints, only 10 of 19 strips could be studied.

To achieve surface sensitivity, XANES spectra were collected using an incidence angle that is smaller than the critical angle, θ_c , below which X-ray total reflection occurs. The measurement geometry was controlled by placement of the sample on a Huber goniometer mounted onto a 2-circle Huber stage (Fig. S5†). The excitation X-ray photon energy was set at 5733 eV for the non-graded films, and 6000 eV for the compositionally graded film, or 10 and 77 eV, respectively, above the Ce L₃ edge (5723 eV). At these excitation energies, the critical angle for the CZO compositions studied is $\approx 0.45^\circ$, as confirmed experimentally (Fig. S6†). At the incidence angle of $\alpha \approx 0.23^\circ$ used for surface-sensitive measurements, the penetration depths are estimated to be 2–3 nm, and the change in penetration depth on passing through the absorption edge is no more than about 1 nm.³⁷ Bulk sensitivity was achieved by setting $\alpha = 10^\circ$, at which the penetration depth exceeds the 220 nm thickness of the films. Hereafter, spectra at these two angles are respectively referred to as surface and bulk spectra, although it is to be noted that the 'bulk' measurement includes the surface region.

Environmental conditions were controlled using an Anton Paar (DHS1100) dome apparatus. Vacuum conditions were

achieved using a turbo pump (Pfeiffer D-35614 Asslar); the oxygen partial pressure is estimated to be 2.3 mTorr on the basis of a residual gas analysis measurement (RGA 200, Stanford Research Systems) of the exhaust stream from the dome apparatus. The measurement under hydrogen was performed using a mixture of 3.5% H₂ (balance He), bubbled through H₂O (held at 15 °C) so as to achieve an oxygen partial pressure of 7.4×10^{-14} mTorr, as implied by equilibrium between H₂O, H₂ and O₂. It is known that the bulk Ce³⁺ concentration in CeO₂ is negligible under all five of the conditions evaluated.⁴⁴ The bulk Ce³⁺ concentration is also negligible for *all compositions* for the measurements in *air*.¹² Under high temperature, and either vacuum or hydrogen, however, the Ce³⁺ concentration in the bulk increases with Zr content.¹² Thus, the five conditions provide access to a range of bulk oxidation states, the impacts of which on surface characteristics could then be explored.

The Ce³⁺ concentration was quantified by a linear combination routine using Ce³⁺ and Ce⁴⁺ reference spectra collected here as standards. All analyses were carried out using the Athena⁴⁵ software package. Each spectrum was normalized to unity by the edge height and the linear background removed. Example reference spectra, specifically from CeAlO₃ and CeO₂ films measured under ambient conditions and at $\alpha = 10^\circ$, are shown in Fig. 2. The intense peak at 5725 eV in the CeAlO₃ spectrum (ref. 46) is typical of materials bearing the Ce³⁺ ion including Ce(NO₃)₃ (ref. 47), Ce(OH)₃,⁴⁸ CePO₄,⁴⁹ Ce₂S₃,⁹ and CeF₃.⁵⁰ The XPS measurements (Fig. S7†) similarly revealed only the presence of Ce³⁺. The CeO₂ spectrum is composed of four

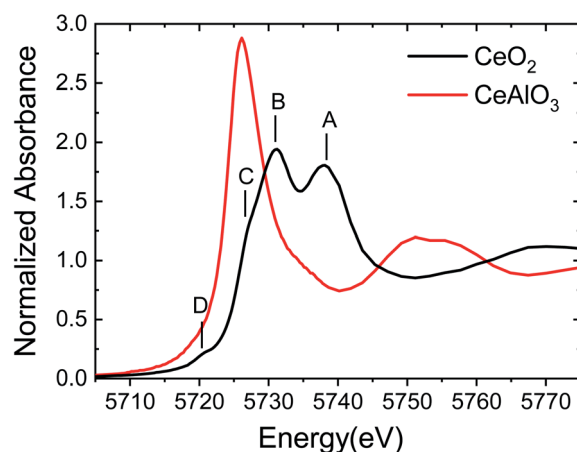


Fig. 2 Normalized Ce L₃ edge XANES spectra of the bulk regions of CeO₂ and CeAlO₃ films (as indicated), collected under ambient temperature and pressure.

elemental features: A, B, C, and D. All four features of this complex spectrum have been accounted for satisfactorily by Soldatov *et al.* as characteristics of fully oxidized CeO₂ using a full multiple-scattering approach.⁵¹ Even the controversial C feature,^{9,52,53} though frequently attributed to the presence of Ce³⁺, has been proven experimentally by high resolution X-ray emission spectroscopy to be a spectral characteristic associated with the Ce⁴⁺ cation.⁵⁴ Upon exposure to high temperature (800 and 1100 °C), the reference spectra are changed as a result of thermal broadening (Fig. S8†). In the case of ceria, the broadening has the apparent effect, at first glance, of increasing the intensity of the C feature, suggesting some thermal reduction and the generation of Ce³⁺ species. However, the integrated peak intensity associated with this feature (Table S2†) is statistically unchanged in response to the temperature changes, consistent with the thermogravimetric measurements that reveal that the Ce³⁺ concentration in bulk CeO₂ is negligible up to at least 1100 °C in air.⁴⁴

Meaningful implementation of linear combination fitting (LCF) methods for quantitative analysis required measurement of reference spectra for all compositions at the temperatures of

interest in order to account for thermal effects. The bulk XANES spectra of CZO films as measured in air at room temperature, at 800 °C, and at 1100 °C are shown in Fig. S9.† The general features of the Zr-bearing films are similar to those of CeO₂, with peaks A–D all visible. The heights of peaks A and B increase slightly with increasing Zr content. In addition, the high temperature exposure causes a dampening of features similar to that observed for CeO₂, with minimal changes to relative integrated peak intensities. Accordingly, we conclude that high-temperature induced spectral changes again result from thermal vibrations in the structures, without changes in oxidation state, in agreement with bulk thermogravimetric behavior.¹² Thus, these high temperature spectra serve as appropriate references for Ce in the 4+ oxidation state for the CZO composition of interest. With this understanding established, the spectra were represented according to

$$I(T, \text{CZO}_{xx}, \text{condition}) = f_1 \times I(T, \text{CZO}_{xx}, \text{air, bulk}) + f_2 \times I(T, \text{CeAlO}_3, \text{bulk}) \quad (1)$$

where I is the spectral intensity, T is the measurement temperature, and ‘condition’ specifies surface *vs.* bulk-

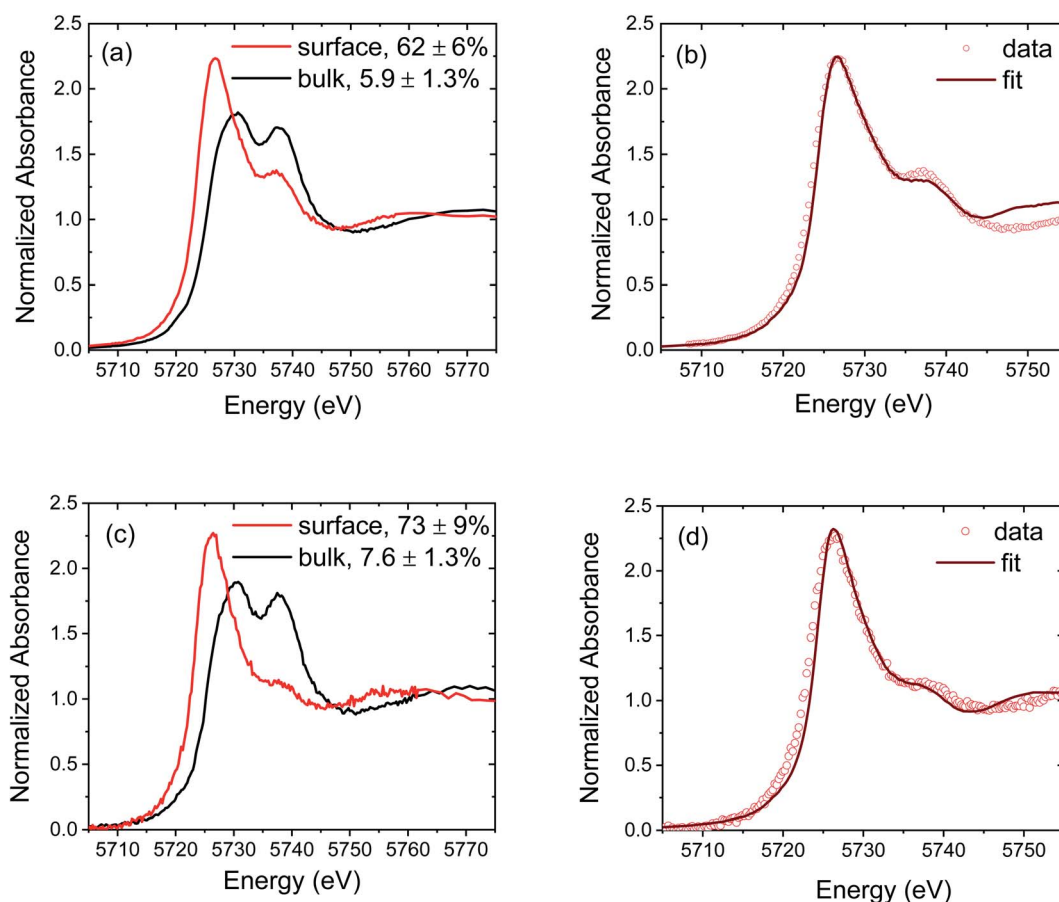


Fig. 3 Representative normalized Ce L₃ edge XANES spectra of ceria-based films: (a) comparison of the bulk and surface spectra of CeO_{2-δ} measured at 1100 °C under vacuum, (b) linear combination fitting result of the surface spectrum of (a) according to eqn (1), (c) comparison of the bulk and surface spectra of CZO_{4.25} measured at 800 °C under H₂/He, and (d) linear combination fitting result of the surface spectrum of (c) according to eqn (1). In (a) and (c), the concentration of Ce³⁺, relative to the total Ce content, obtained from the LCF analysis is given in the legend.

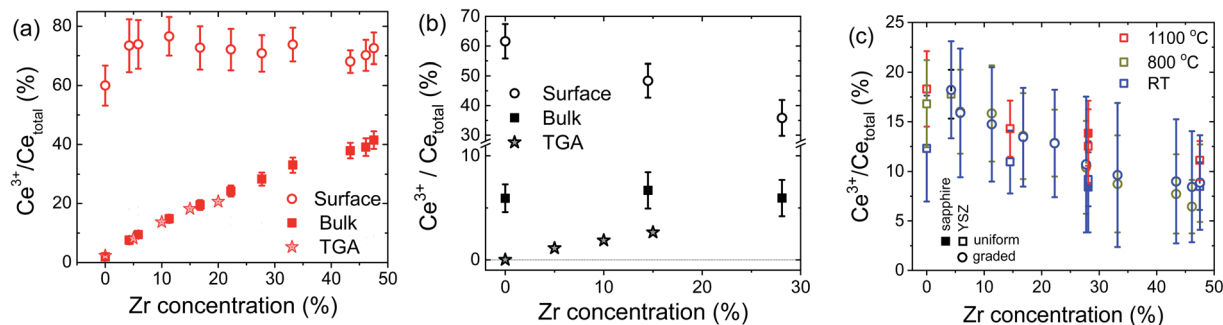


Fig. 4 Percentage of Ce in the 3+ oxidation state as a function of zirconium concentration: (a) at the surface and in the bulk as measured at 800 °C under H₂ and compared to the bulk TGA¹² prediction, (b) at the surface and in the bulk as measured at 1100 °C under vacuum ($p_{O_2} \approx 2.3$ mTorr), and compared to the bulk TGA¹² prediction; and (c) at the surface as measured in air at the temperatures indicated. The bulk Ce³⁺ concentrations in air, at all measurement temperatures, are taken to be zero (and provide the reference spectra for the LCF analysis).

dominated spectra as well as the gas atmosphere. The Ce³⁺ fraction is computed according to $f_2/(f_1 + f_2)$. The uncertainty in the Ce³⁺ concentration obtained from the LCF was estimated according to

$$\text{Uncertainty} = \frac{\sum |I_{\text{exp}} - I_{\text{fit}}|}{\sum |I_{\text{exp}}|} \quad (2)$$

where I_{exp} is the experimentally measured spectrum and I_{fit} is the LCF result according to eqn (1). The analysis is carried out over the energy range 20 eV below and 30 eV above the edge energy.

Results and discussion

Two pairs of representative spectra comparing bulk and surface-dominated regions of CeO_{2- δ} and CZO4.25 are presented in Fig. 3. The differences between bulk and surface regions are striking. The most intense feature in the surface spectra appears near 5725 eV, which coincides with the intense peak measured for CeAlO₃, while peaks A and B become markedly diminished relative to the analogous bulk spectra. The LCF analyses indicate a remarkable 10-fold enhancement of the Ce³⁺ concentration in the near-surface region relative to that in the bulk. This dramatic difference between surface and bulk reduction is in fact reflected under all of the measurement conditions and for all of the compositions, as summarized in Fig. 4. The complete set of spectra upon which this result is based are provided in Fig. S10 and S11† for the non-graded and graded films, respectively.

Several important trends are evident from Fig. 4. First, at 800 °C and under H₂, Fig. 4(a), the bulk Ce³⁺ concentration is in excellent quantitative agreement with that obtained from bulk thermogravimetric measurements of the oxygen non-stoichiometry.¹² The agreement not only gives confidence in the LCF methodology, but also implies that the surface-reduced region does not, within experimental uncertainty, extend beyond the 2–3 nm detected by XANES at $\alpha \approx 0.23^\circ$. Second, ~75% of the Ce in the near-surface region exists as Ce³⁺ under these conditions, irrespective of film composition. Accordingly, the almost 10-fold enhancement at low Zr content falls to less than a 2-fold enhancement at 48% Zr. In short, the enhanced reducibility of

the bulk with increasing Zr content does not translate into an enhanced reducibility at the surface.

Turning to the 1100 °C measurements under vacuum, Fig. 4(b), the XANES measurement indicates a composition independent Ce³⁺ bulk concentration of ~6%. Both the magnitude of the bulk Ce³⁺ concentration and its independence of Zr concentration are unexpected. Under these conditions, the bulk Ce³⁺ concentration is expected to be negligible for ceria and increase to only ~5% for CZO28 (from an extrapolation of the reported behavior of up to 20% Zr¹²). At the same time, the surface Ce³⁺ concentration determined by XANES falls steeply with Zr content, precisely the opposite of what one expects from the enhanced bulk reducibility. Again, bulk behavior cannot be assumed to describe surface trends.

The disagreement between bulk TGA and bulk XANES results at 1100 °C under vacuum conditions warrants some discussion. The possibility that the discrepancy arises from errors in thermogravimetry is considered small. Taking specifically the ceria film, the XANES result of 5.9% Ce³⁺ and an assumption that these species are fully charge-compensated by oxygen vacancies imply a CeO_{1.97} stoichiometry. A change in oxygen content from CeO₂ to CeO_{1.97} would be accompanied by a change in mass of 0.28%, a value routinely detected by gravimetry. Similarly, because of the agreement between the LCF analysis and the TGA results under H₂ at 800 °C, even for compositions with low Ce³⁺ concentrations, we consider the likelihood of the disagreement arising from artifacts of the XANES analysis methodology to be small. An alternative possibility that is also immediately ruled out is a strain-induced discrepancy. Because the lattice constants of the films are larger than those of YSZ, particularly at low Zr content, one might expect the films to be under compressive strain in the near vicinity of the substrate. Reduction, however, induces chemical expansion,⁵⁵ and therefore any compression would be expected to decrease the extent of film reduction, not increase it beyond the expected thermodynamic value, as is observed. With these possibilities ruled out, we suggest that at 1100 °C, in contrast to the behavior at 800 °C, the surface-reduced region may extend somewhat into the bulk, contributing to the Ce³⁺ signal detected at an incident angle of 10°. Extension of the reduced region to a depth of just 15–20 nm would yield film-averaged Ce³⁺ concentrations that are consistent with the bulk

TGA results. Convolution of bulk and extended surface contributions, with opposite dependences on Zr content, could thus account for the relative insensitivity of the Ce^{3+} concentration to composition in the bulk-dominated results.

Turning to the surface reduction in air, Fig. 4(c), the trends here are as surprising as the results under more strongly reducing conditions. As already noted, in air the bulk has negligible Ce^{3+} content, irrespective of Zr concentration and temperature,¹² and the bulk spectra were used to establish the Ce^{4+} reference spectra at the three measurement temperatures. The surface Ce^{3+} concentration in air is found to be significant at all temperatures, reaching almost 20% for ceria. Furthermore, the surface reduction decreases with increasing Zr content, similar to the behavior under vacuum at 1100 °C, demonstrating yet again that the surface trends are precisely the opposite of what might be anticipated based on bulk reducibility. Another highly unexpected result is the rather slight sensitivity to temperature. There is no statistical difference between measurements at ambient temperature and 800 °C (graded film), while there is a slight increase in surface reduction on increasing the temperature from 25 to 1100 °C (uniform composition films). Shown in Fig. 4(c) are also results from selected measurements performed using CZO28 on *r*-sapphire, from which it is apparent that the nature of the substrate has a negligible impact on the measured surface characteristics.

The changes in oxidation state in the films due to the changes in environmental conditions were found to be reversible, as evidenced for example, from the behavior of the uniform-composition CZO48 film (Fig. S12†). The spectra measured at the initiation of the experiment under ambient conditions (25 °C, air) and again after exposure to high temperature, vacuum conditions were identical to one another, for both the surface and bulk regions of the film. The absence of beam-induced film damage, already implied by the reversibility, was directly ascertained by comparing the sequence of spectra measured to generate the integrated spectra used in the analysis. The results, for example from $\text{CeO}_{2-\delta}$ measured at 1100 °C under vacuum (Fig. S13†), reveal the absence of any temporal evolution.

The enhanced reduction of bulk CZO over ceria has been attributed to the smaller coordination number of Zr relative to that of Ce, with 7-fold coordination of the Zr^{4+} cation being globally accommodated by a decrease in the Ce oxidation state from 4+ to 3+.¹¹ In the surface region, it is plausible that the disruption to the crystalline structure inherently accommodates 7-fold coordination of Zr^{4+} without requiring as great a loss of oxygen and thus without requiring as great a change in the Ce oxidation state. If catalytic activity is correlated with the concentration of surface oxygen vacancies, as many have proposed,^{4,29,30} and Ce^{3+} surface concentration is indeed proportional to the vacancy concentration, these results suggest that activity *decreases* with increasing Zr content. To date, there has been no systematic report describing the surface catalytic activity of $\text{Ce}_{1-x}\text{Zr}_x\text{O}_2$ materials as a function of composition in the absence of convoluting microstructural effects (see, for example, the studies of Aneggi *et al.*,⁵⁶ Bulfin *et al.*,⁵⁷ and Piumetti *et al.*⁵⁸). Hence, it is unknown whether introducing Zr,

which is desirable for increasing the thermal stability and bulk oxygen storage capacity, has a detrimental impact on area-specific catalytic reaction rates.

Summary and conclusions

Using XANES measurements in combination with a thin film geometry, the surface and bulk Ce^{3+} concentrations in a series of $\text{Ce}_{1-x}\text{Zr}_x\text{O}_2$ materials have been quantified under a range of operationally relevant conditions. The findings are summarized as follows:

- Under highly reducing conditions (800 °C, 3.5% H_2), the bulk Ce^{3+} concentration is in quantitative agreement with prior literature values determined from thermogravimetric analysis. Specifically, the Ce^{3+} concentration (relative to the amount of Ce) increases from 0 to ~41% with an increase in Zr content from 0 to 48 mol%.

- Under these same conditions, the surface is highly reduced. The Ce^{3+} concentration is ~75% and is approximately independent of Zr concentration. Thus, the enhanced bulk reducibility imparted to ceria by Zr is not reflected in the surface properties.

- Under less extreme atmospheric conditions (1100 °C in a vacuum; 1100 °C, 800 °C, and RT in air), the surface is again substantially more reduced than the bulk. Remarkably, the surface Ce^{3+} concentration *decreases* with increasing Zr content, in opposition to the trend in the bulk.

- Because all XANES results reported here are those obtained after equilibrium has been achieved, both the surface and bulk concentrations are thermodynamic values, free of any kinetic effects.

- Possible artifacts such as beam damage, or a through-plane gradient in Zr content, or residual stress in the films have all been experimentally ruled out as possible explanations for these results.

Significantly, these findings stand in direct contrast to computational predictions of enhanced surface reducibility in $\text{Ce}_{1-x}\text{Zr}_x\text{O}_2$ with increasing Zr content. Uncovering the connection, if any, between the surface redox state and catalytic activity in this material class will be the subject of future studies. Beyond the ceria–zirconia system, this work cautions against the use of bulk descriptors, in particular, bulk defect concentrations, as predictors of catalytic activity. Furthermore, the methodology developed here opens up new possibilities for directly measuring and quantifying surface characteristics of redox active oxides under technologically relevant conditions in conjunction with a high-throughput configuration.

Conflicts of interest

There are no conflicts to declare.

Acknowledgements

This work was supported primarily by NSF DMR-1505103, with additional partial support provided by the Institute for Catalysis in Energy Processes at Northwestern University under DOE

Grant DE-FG02-03ER15457 and MRSEC under NSF Grant DMR-1720139. The XANES and XRF experiments of uniform films were performed at the DuPont-Northwestern-Dow Collaborative Access Team (DND-CAT) located at Sector 5 of the Advanced Photon Source (APS). DND-CAT is supported by Northwestern University, E.I. DuPont de Nemours & Co., and The Dow Chemical Company. This research used resources of the Advanced Photon Source, a U.S. Department of Energy (DOE) Office of Science User Facility operated for the DOE Office of Science by Argonne National Laboratory under Contract No. DE-AC02-06CH11357. Sector 20 operations are supported by the US Department of Energy and the Canadian Light Source. Part of this work made use of the Pulsed Laser Deposition Shared Facility and the X-ray Diffraction Facility at the Materials Research Center at Northwestern University, supported by the National Science Foundation MRSEC program (DMR-1720139) and the Soft and Hybrid Nanotechnology Experimental (SHyNE) Resource (NSF ECCS-1542205).

References

- 1 A. Trovarelli, *Catal. Rev.*, 1996, **38**, 439–520.
- 2 R. Di Monte and J. Kaspar, *Catal. Today*, 2005, **100**, 27–35.
- 3 S. Liu, X. Wu, D. Weng and R. Ran, *J. Rare Earths*, 2015, **33**, 567–590.
- 4 J. Paier, C. Penschke and J. Sauer, *Chem. Rev.*, 2013, **113**, 3949–3985.
- 5 T. X. T. Sayle, S. C. Parker and C. R. A. Catlow, *Surf. Sci.*, 1994, **316**, 329–336.
- 6 E. Mamontov, T. Egami, R. Brezny, M. Koranne and S. Tyagi, *J. Phys. Chem. B*, 2000, **104**, 11110–11116.
- 7 P. Fornasiero, G. Balducci, R. DiMonte, J. Kaspar, V. Sergo, G. Gubitosa, A. Ferrero and M. Graziani, *J. Catal.*, 1996, **164**, 173–183.
- 8 C. E. Hori, H. Permana, K. Y. S. Ng, A. Brenner, K. More, K. M. Rahmoeller and D. Belton, *Appl. Catal., B*, 1998, **16**, 105–117.
- 9 F. Zhang, C.-H. Chen, J. M. Raitano, J. C. Hanson, W. A. Caliebe, S. Khalid and S.-W. Chan, *J. Appl. Phys.*, 2006, **99**, 084313.
- 10 T. Kim, J. M. Vohs and R. J. Gorte, *Ind. Eng. Chem. Res.*, 2006, **45**, 5561–5565.
- 11 M. Kuhn, S. R. Bishop, J. L. M. Rupp and H. L. Tuller, *Acta Mater.*, 2013, **61**, 4277–4288.
- 12 Y. Hao, C.-K. Yang and S. M. Haile, *Chem. Mater.*, 2014, **26**, 6073–6082.
- 13 J. R. Scheffe and A. Steinfeld, *Energy Fuels*, 2012, **26**, 1928–1936.
- 14 J. C. Conesa, *Surf. Sci.*, 1995, **339**, 337–352.
- 15 G. Balducci, J. Kašpar, P. Fornasiero, M. Graziani and M. S. Islam, *J. Phys. Chem. B*, 1998, **102**, 557–561.
- 16 Z. Yang, T. K. Woo, M. Baudin and K. Hermansson, *J. Chem. Phys.*, 2004, **120**, 7741–7749.
- 17 Z. X. Yang, Y. W. Wei, Z. M. Fu, Z. S. Lu and K. Hermansson, *Surf. Sci.*, 2008, **602**, 1199–1206.
- 18 Z. X. Yang, Z. M. Fu, Y. W. Wei and K. Hermansson, *Chem. Phys. Lett.*, 2008, **450**, 286–291.
- 19 Z. Yang, T. K. Woo and K. Hermansson, *J. Chem. Phys.*, 2006, **124**, 7.
- 20 M. Nolan, S. C. Parker and G. W. Watson, *Surf. Sci.*, 2005, **595**, 223–232.
- 21 Z. A. Feng, F. El Gabaly, X. Ye, Z.-X. Shen and W. C. Chueh, *Nat. Commun.*, 2014, **5**, 4374.
- 22 F. Zhang, P. Wang, J. Koberstein, S. Khalid and S. W. Chan, *Surf. Sci.*, 2004, **563**, 74–82.
- 23 S. Deshpande, S. Patil, S. Kuchibhatla and S. Seal, *Appl. Phys. Lett.*, 2005, **87**, 3.
- 24 P. Dutta, S. Pal, M. S. Seehra, Y. Shi, E. M. Eyring and R. D. Ernst, *Chem. Mater.*, 2006, **18**, 5144–5146.
- 25 S. Kim, R. Merkle and J. Maier, *Surf. Sci.*, 2004, **549**, 196–202.
- 26 S. C. DeCaluwe, M. E. Grass, C. Zhang, F. E. Gabaly, H. Bluhm, Z. Liu, G. S. Jackson, A. H. McDaniel, K. F. McCarty, R. L. Farrow, M. A. Linne, Z. Hussain and B. W. Eichhorn, *J. Phys. Chem. C*, 2010, **114**, 19853–19861.
- 27 W. C. Chueh, A. H. McDaniel, M. E. Grass, Y. Hao, N. Jabeen, Z. Liu, S. M. Haile, K. F. McCarty, H. Bluhm and F. El Gabaly, *Chem. Mater.*, 2012, **24**, 1876–1882.
- 28 Z. L. Zhao, M. Uddi, N. Tsvetkov, B. Yildiz and A. F. Ghoniem, *Phys. Chem. Chem. Phys.*, 2017, **19**, 25774–25785.
- 29 Z. X. Yang, T. K. Woo, M. Baudin and K. Hermansson, *J. Chem. Phys.*, 2004, **120**, 7741–7749.
- 30 E. Aneggi, J. Llorca, M. Boaro and A. Trovarelli, *J. Catal.*, 2005, **234**, 88–95.
- 31 A. Trovarelli and J. Llorca, *ACS Catal.*, 2017, **7**, 4716–4735.
- 32 V. K. Paidi, L. Savereide, D. J. Childers, J. M. Notestein, C. A. Roberts and J. van Lierop, *ACS Appl. Mater. Interfaces*, 2017, **9**, 30670–30678.
- 33 G. Bunker, *Introduction to XAFS. A Practical Guide to X-ray Absorption Fine Structure Spectroscopy*, Cambridge University Press, 2010.
- 34 A. I. Kozlov, D. H. Kim, A. Yezerets, P. Andersen, H. H. Kung and M. C. Kung, *J. Catal.*, 2002, **209**, 417–426.
- 35 P. Duran, M. Gonzalez, C. Moure, J. R. Jurado and C. Pascual, *J. Mater. Sci.*, 1990, **25**, 5001–5006.
- 36 A. Fluri, D. Pergolesi, A. Wokaun and T. Lippert, *Phys. Rev. B*, 2018, **97**, 125412.
- 37 M. Birkholz, *Thin Film Analysis by X-ray Scattering*, Wiley-VCH, 2006.
- 38 H. Koinuma and I. Takeuchi, *Nat. Mater.*, 2004, **3**, 429.
- 39 M. Charnley, J. Gilbert, O. G. Jones and N. P. Reynolds, *Bio-Protoc.*, 2018, **8**, 2732.
- 40 W. C. Chueh, Y. Hao, W. Jung and S. M. Haile, *Nat. Mater.*, 2012, **11**, 155–161.
- 41 W. C. Chueh and S. M. Haile, *Phys. Chem. Chem. Phys.*, 2009, **11**, 8144–8148.
- 42 O. H. Laguna, A. Perez, M. A. Centeno and J. A. Odriozola, *Appl. Catal., B*, 2015, **176**, 385–395.
- 43 S. M. Heald, J. O. Cross, D. L. Brewster and R. A. Gordon, *Nucl. Instrum. Methods Phys. Res., Sect. A*, 2007, **582**, 215–217.
- 44 R. J. Panlener, R. N. Blumenthal and J. E. Garnier, *J. Phys. Chem. Solids*, 1975, **36**, 1213–1222.
- 45 B. Ravel and M. Newville, *J. Synchrotron Radiat.*, 2005, **12**, 537–541.

- 46 A. Martínez-Arias, M. Fernández-García, L. N. Salamanca, R. X. Valenzuela, J. C. Conesa and J. Soria, *J. Phys. Chem. B*, 2000, **104**, 4038–4046.
- 47 J. Zhang, Z. Y. Wu, T. Liu, T. D. Hu, Z. H. Wu and X. Ju, *J. Synchrotron Radiat.*, 2001, **8**, 531–532.
- 48 J. El Fallah, S. Boujana, H. Dexpert, A. Kiennemann, J. Majerus, O. Touret, F. Villain and F. Lenormand, *J. Phys. Chem.*, 1994, **98**, 5522–5533.
- 49 A. Bianconi, A. Marcelli, M. Tomellini and I. Davoli, *J. Magn. Mater.*, 1985, **47–48**, 209–211.
- 50 G. Kaindl, G. K. Wertheim, G. Schmiester and E. V. Sampathkumaran, *Phys. Rev. Lett.*, 1987, **58**, 606–609.
- 51 A. V. Soldatov, T. S. Ivanchenko, S. Della Longa, A. Kotani, Y. Iwamoto and A. Bianconi, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1994, **50**, 5074–5080.
- 52 G. Kaindl, G. Schmiester, E. V. Sampathkumaran and P. Wachter, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 1988, **38**, 10174–10177.
- 53 P. Nachimuthu, W. C. Shih, R. S. Liu, L. Y. Jang and J. M. Chen, *J. Solid State Chem.*, 2000, **149**, 408–413.
- 54 G. Gasperi, L. Amidani, F. Benedetti, F. Boscherini, P. Glatzel, S. Valeri and P. Luches, *Phys. Chem. Chem. Phys.*, 2016, **18**, 20511–20517.
- 55 D. Marrocchelli, S. R. Bishop, H. L. Tuller and B. Yildiz, *Adv. Funct. Mater.*, 2012, **22**, 1958–1965.
- 56 E. Aneghi, C. de Leitenburg, J. Llorca and A. Trovarelli, *Catal. Today*, 2012, **197**, 119–126.
- 57 B. Bulfin, F. Call, J. Vieten, M. Roeb, C. Sattler and I. V. Shvets, *J. Phys. Chem. C*, 2016, **120**, 2027–2035.
- 58 M. Piumetti, S. Bensaid, D. Fino and N. Russo, *Appl. Catal., B*, 2016, **197**, 35–46.