

Cite this: *Chem. Commun.*, 2017, 53, 5231Received 27th February 2017,
Accepted 19th April 2017

DOI: 10.1039/c7cc01538k

rsc.li/chemcomm

In situ investigation of degradation at organometal halide perovskite surfaces by X-ray photoelectron spectroscopy at realistic water vapour pressure†

 Jack Chun-Ren Ke,^{ab} Alex S. Walton,^{bc} David J. Lewis,^d
 Aleksander Tedstone,^c Paul O'Brien,^c Andrew G. Thomas^{ab}* and
 Wendy R. Flavell^{ab}

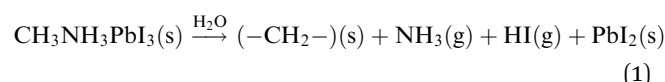
Near-ambient-pressure X-ray photoelectron spectroscopy enables the study of the reaction of *in situ*-prepared methylammonium lead iodide (MAPI) perovskite at realistic water vapour pressures for the first time. We show that MAPI decomposes directly to PbI₂, HI and NH₃ without formation of methylammonium iodide, allowing us to distinguish between alternative mechanisms for the atmospheric degradation reaction.

Organometal halide perovskite (OHP) absorbers have rapidly emerged as one of the most promising photovoltaic materials for use in solar cells due to their high visible to near-infrared light absorptivity,^{1–4} remarkably low exciton binding energy,⁵ and extraordinarily long diffusion length for electrons and holes.^{6,7} As a result, perovskite solar cells (PSCs) have achieved certified power conversion efficiencies of more than 20%,⁸ in just over 7 years from the initial discovery of the photovoltaic effect in TiO₂ coated with nanocrystalline perovskite.⁹ These efficiencies are now comparable to those of commercial polycrystalline silicon solar cells. However, OHPs show poor long-term stability, particularly when exposed to humid atmospheres which means that there are questions about the commercial exploitation of PSCs in solar panels.¹⁰ Improving the stability of OHP absorbers is of paramount importance, but only with insight into the degradation mechanisms at the atomic level can solutions thus be realised.

Numerous characterisation techniques have been applied to investigate the degradation mechanism, such as ultraviolet-visible (UV-Vis) spectroscopy,^{11–13} X-ray diffraction (XRD),^{11,13,14} Raman spectroscopy,¹⁵ and X-ray photoelectron spectroscopy

(XPS).^{16–18} This latter technique is useful since it allows *in situ* deposition of the films in an ultrahigh vacuum (UHV) environment,^{17,19–21} so that degradation of the films before measurement is ruled out. Several XPS studies have been carried out to determine the surface composition or film growth mechanism,^{19–27} and band alignment.^{28–30}

Despite these studies, the mechanism of degradation of pristine films at the atomic and nanometre scale remains imprecisely defined. Currently, two models of OHP degradation have been proposed. The first involves the water-catalysed loss of nitrogen, in the form of ammonia, and hydrogen iodide *via* the reaction (eqn (1)):¹⁷



Thermogravimetric analysis of bulk MAPI by Grätzel and co-workers suggests that the organic component of the perovskite material is lost *via* volatilisation of HI and CH₃NH₂ with the latter species lost at higher temperatures.³¹ The second proposal involves decomposition to methyl ammonium iodide (MAI) and lead iodide in the presence of liquid water (eqn (2)):¹⁶



Key difficulties in understanding the degradation mechanism of MAPI are the difficulty in determining the surface chemistry of the first few atomic layers of the pristine material without any exposure to ambient atmospheric conditions, and, conversely, the difficulty in measuring this same surface chemistry in real time under exposure to realistic environments. In this work we address both problems simultaneously by utilising Near-Ambient Pressure XPS (NAPXPS) to investigate the real-time degradation of an *in situ*-deposited MAPI film, exposed to water vapour pressures equivalent to atmospheric relative humidity (RH) of *ca.* 30% at a standard temperature of 25 °C, conditions which could be realistically be observed under ambient atmospheric environments. This yields an unparalleled mechanistic insight. The data unambiguously indicate a degradation pathway

^a School of Physics and Astronomy, The University of Manchester, Oxford Road, Manchester, M13 9PL, UK. E-mail: wendy.flavell@manchester.ac.uk

^b Photon Science Institute, The University of Manchester, Oxford Road, Manchester, M13 9PL, UK. E-mail: andrew.g.thomas@manchester.ac.uk

^c School of Chemistry, The University of Manchester, Oxford Road, Manchester, M13 9PL, UK

^d School of Materials, The University of Manchester, Oxford Road, Manchester, M13 9PL, UK

† Electronic supplementary information (ESI) available: Experimental details and supporting results. See DOI: 10.1039/c7cc01538k



by reaction with water which involves complete loss of nitrogen from the film without the formation of MAI.

The MAPI film was prepared in UHV by vacuum deposition of PbCl_2 and MAI onto a $\text{SrTiO}_3(100)$ single crystal substrate, using the two-step route of Liu *et al.* (see ESI† for details of preparation and characterisation).³² Fig. 1a shows high-resolution NAPXPS spectra of the Sr 3d and Pb 4f core levels recorded from the MAPI film before water exposure, during water exposure at H_2O vapour pressures of 3 mbar (RH 10%) and 9 mbar (RH 30%), and after the water vapour has been removed and the system returned to UHV conditions. The intensity is normalised to the integrated area of the Sr $3d_{5/2}$ peak from the SrTiO_3 substrate, which allows comparison of the intensity of the Pb 4f peaks as a function of water exposure. All photoelectron binding energies are quoted to a precision of ± 0.05 eV. Prior to exposure to H_2O vapour and under UHV, the Pb $4f_{7/2}$ peak of the perovskite film is located at a binding energy (BE) of 138.60 eV, which is in excellent agreement with previous reports for the Pb in MAPI perovskite materials (denoted 'Pb-I' in Fig. 1).^{16,22,30,33} We note that there is a weaker Pb $4f_{7/2}$ peak at 136.80 eV, which is attributed to metallic lead.^{16,22,30,33} This has been observed for several perovskite materials, such as MAPI,^{16,22,27,33} MAPbI_2Cl ,³⁴ MAPbBr_3 .³³ Metallic lead is often observed in MAPI perovskite films, both for the standard solution-processed films^{35,36} and CVD-deposited films.²² In the study presented here the metallic Pb probably arises due to

Table 1 Stoichiometries of the film determined from UHV XPS at different stages, before and after water exposure. The metallic lead content is excluded from these quantifications. A full description of the quantification is given in the ESI

Status ID	I/Pb	N/Pb
UHV before water exposure	2.7 ± 0.1	0.9 ± 0.2
UHV after water exposure	2.4 ± 0.1	0.0 ± 0.2

decomposition of the PbCl_2 precursor vapour during the initial deposition step;²¹ some beam-induced decomposition cannot be ruled out, but is not the major source of elemental Pb (as we discuss further below). No Cl was observed in survey scans (see Fig. S1, ESI†). The Pb:I:N stoichiometry of the film determined from XPS, excluding the metallic Pb component, was found to be $1.0 \pm 0.1 : 2.7 \pm 0.1 : 0.9 \pm 0.2$ (Table 1), in good agreement with the nominal stoichiometry of MAPI. Quantification of the XPS signals is discussed further below and in the ESI† (Tables S1–S3).

Fig. 1(a) shows Pb 4f and Sr 3d spectra recorded during exposure to H_2O vapour. The Pb-I features decrease in intensity, whilst the metallic Pb features increase in intensity. The BE of the Pb $4f_{7/2}$ Pb-I feature shifts upwards slightly to 138.70 eV on exposure to 3 mbar H_2O and to 138.80 eV when the water exposure is increased to 9 mbar. This may be due to the formation of lead iodide (Pb $4f_{7/2}$ BE = 138.90 eV);¹⁶ however the peak cannot be resolved into individual components arising from MAPI and PbI_2 . The BE of the MAPI/ PbI_2 peak recorded after water exposure, when the water vapour is pumped away and the sample returned to ultra-high vacuum conditions (Fig. 1(a)-iv), is not shifted relative to the spectrum recorded at a water vapour pressure of 9 mbar (Fig. 1(a)-iii). This suggests that the presence of water vapour is not responsible for the observed binding energy shifts, *i.e.* the shifts are due to changes in the nature of the material. The data recorded following removal of the water vapour indicate that the changes induced by exposure to water vapour are not reversible. We observe an additional increase in the intensity of the metallic-lead-derived peak in the UHV spectrum taken following water exposure. In synchrotron studies this was attributed to beam damage,¹⁶ and it is well known that PbI_2 slowly decomposes to Pb and I_2 under illumination with visible light.³⁷ However, spectra taken at this stage from fresh points on the sample surface (see ESI,† Fig. S3) show a similar surface stoichiometry to that in Fig. 1(a)-iv, and in particular, similar amounts of metallic Pb. We conclude that the decomposition of PbI_2 into Pb and I_2 cannot be attributed solely to beam damage.

Fig. 1b shows I 3d core-level spectra, also normalised to the integrated area of the Sr $3d_{5/2}$ peak. The BE of the I $3d_{5/2}$ shifts towards higher BE, from 619.50 eV (UHV before) to 619.60 eV at 3 mbar water exposure and reaches 619.65 eV at 9 mbar water exposure. The change in BE matches the changes in the Pb 4f spectra. The I signal decreases significantly upon water exposure, suggesting it is lost from the surface, which is not consistent with a degradation mechanism resulting in the formation of MAI. We note that the intensity increases slightly when the water vapour is removed. This is a consequence of the lower kinetic energy of the I 3d electrons relative to Pb 4f. Lower kinetic energy electrons are

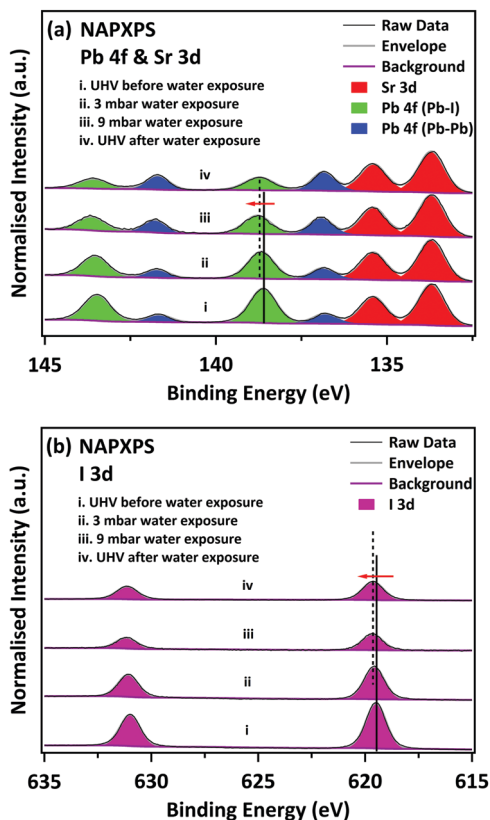


Fig. 1 NAPXPS spectra of (a) Pb 4f & Sr 3d, and (b) I 3d core levels under various conditions as noted (before/during/after water exposure). All the spectra are normalised to the corresponding integrated area of the Sr $3d_{5/2}$ peak.



less likely to be detected by the analyser than high energy ones as the gas pressure in the NAPXPS cell is increased.

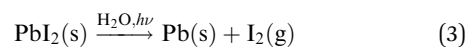
Fig. 2 shows the C 1s and Sr 3p spectra. The C 1s spectrum allows us to investigate the behaviour of the organic part of the perovskite under water vapour exposure. The Sr 3p_{3/2} peaks are located at 269.20 eV, consistent with Sr–O bonding.³⁸ The clean SrTiO₃ spectrum shows a very weak signal associated with a small amount of residual adventitious carbon. In the freshly-prepared MAPI film we observe two C species in the C 1s spectrum. The larger C 1s peak is located at a BE of 285.35 eV. This peak has been assigned both to CH₃I,²¹ and to C–C bonds (hydrocarbon).²¹ Ng *et al.* linked the formation of the C–C species to reaction with PbI₂,²² but it is perhaps more likely to occur due to decomposition of the precursors or simply due to outgassing of the sources during evaporation of the films. A shoulder is also observed at 286.85 eV, 1.5 eV higher than the main peak, corresponding to C–N bonding in the MAPI perovskite.^{16,17,21,34} However, as discussed below, we believe this feature also contains contributions from C=O/C–OH groups in the hydrocarbon contamination. We do not believe the MAI precursor contributes significantly to the higher binding energy feature here, since the core level shift in the C 1s feature between C–C/CH₃I and C–N in MAI is larger at 1.9 eV,²² than the 1.5 eV measured here. Upon exposure to water vapour, the relative intensities of the C 1s peaks change substantially. The higher energy peak, which contains contributions from C–N in the MAPI, is seen to decrease in intensity. During and after exposure to

9 mbar water vapour there is an increase in the relative intensity of the lower binding energy feature as C–N is converted to CH, consistent with the mechanism proposed in eqn (1).

Stronger evidence for the cleavage of the C–N bond is seen in the corresponding N 1s spectra in Fig. 2. Before exposure to water vapour, the N 1s spectrum of the pristine sample lies at a binding energy of 402.45 eV, in good agreement with the methylammonium group in MAPI.^{16,21} During and after exposure to water the N 1s peak is seen to dramatically decrease in intensity during water exposure, and after water exposure the N 1s signal cannot be reliably distinguished from the spectral noise. This confirms that at least some of the intensity in the C 1s peak at 286.75 eV arises from other carbon species (such as C=O/C–OH), since there is some remnant intensity, even after all of the N 1s signal is removed. During and after exposure to 9 mbar water vapour there is an increase in the relative intensity of the lower binding energy feature which may be due to conversion of C–N to C–H/C–C, consistent with the mechanism proposed in eqn (1).

The stoichiometry of the film before and after water exposure is calculated from the integrated peak areas, using CASAXPS sensitivity factors (Table 1).³⁹ These calculations are only performed for the sample under UHV conditions, since RSFs are not available for the analyser under near-ambient conditions. As noted above, the Pb:I:N: ratio of the as-deposited film is close to the expected stoichiometry. Following H₂O vapour exposure the Pb/I ratio decreases and N is completely lost from the surface. Although obtaining information on the stoichiometry of the C part of the material is difficult because of the hydrocarbon contamination, there are some observations which support the decomposition of MAPI to form NH₃. The decrease in intensity of the higher binding energy component of the C 1s spectrum is consistent with the overall loss of N (see ESI,† Table S2). In addition, we find that the total amount C at the surface remains constant, suggesting that C–N is converted to –CH₂– in accordance with eqn (1).

The evidence from NAPXPS strongly supports a decomposition mechanism similar to that suggested by Li *et al.* (eqn (1)),¹⁷ where nitrogen is lost from the perovskite *via* reaction with water to give NH₃(g). The results are consistent with conversion of lead in the MAPI film to PbI₂ (I/Pb is significantly reduced, Table 1), some of which is then further decomposed as shown in eqn (3).



The observed changes in the C 1s spectra induced by water exposure are in agreement with the work of Li *et al.* suggesting that the formation of MAI (*via* eqn (2)) is not involved in the decomposition mechanism in the presence of water vapour.¹⁷ This is further confirmed by the complete loss of the N 1s feature, and a commensurate decrease in the intensity of the feature assigned to C–N in the C 1s spectrum. If MAI were formed we would expect simply to observe a shift to lower BE relative to the MAPI N 1s signal.²¹ The release of nitrogen must therefore occur directly from the MAPI film as stated in eqn (1), probably in the form of ammonia gas. We cannot rule out the formation of some CH₃PbI₃ by loss of ammonia during the

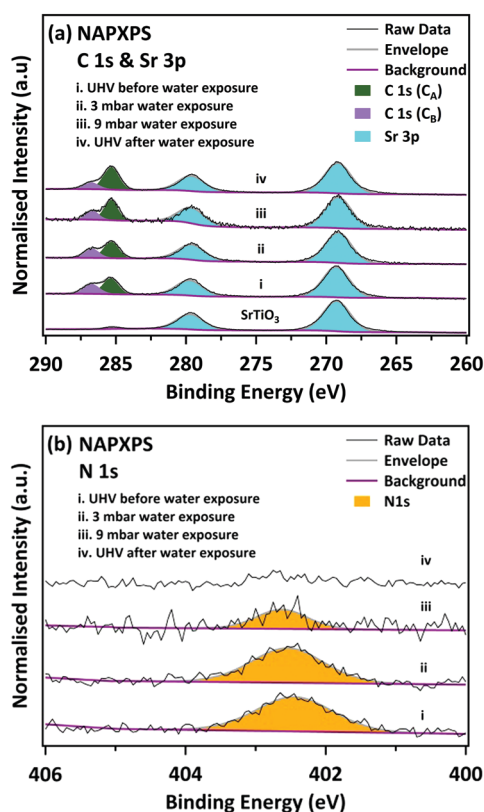


Fig. 2 NAPXPS spectra of (a) C 1s & Sr 3p, and (b) N 1s core levels before/during/after water exposure. All spectra are normalised to the corresponding integrated area of the Sr 3d_{5/2} peak.



degradation process, as its C 1s binding energy cannot be distinguished from that of MAPI.²¹ Indeed, this is suggested by the final Pb/I ratio of 2.4 ± 0.1 , higher than the 2.0 expected if only PbI₂ is produced.

In conclusion, XPS analysis suggests a stoichiometric MAPI thin film was successfully fabricated *via* vapour deposition, in good agreement with previous studies.^{40,41} The data clearly demonstrate the pathway of water-induced degradation of the MAPI film at relative humidity similar to those found in ambient air, using NAPXPS. It is shown that MAI is not formed during decomposition of MAPI during water exposure. The data suggest that at 9 mbar water vapour exposure, MAPI is almost completely transformed into a mixture of Pb, PbI₂ and hydrocarbon species, with a complete loss of nitrogen from the surface. The work further demonstrates the potential for NAPXPS in understanding the surface degradation of materials under atmospheric conditions. In this case it allows us to distinguish unambiguously between the proposed mechanistic pathways,^{16,17} potentially allowing methodologies for mitigating the degradation to be developed.

The authors thank the University of Manchester and EPSRC (UK) (grant number EP/K009710) for funding. C. R. Ke thanks the University of Manchester for the award of a President's Doctoral Scholarship.

Notes and references

- 1 F. Hao, C. C. Stoumpos, D. H. Cao, R. P. Chang and M. G. Kanatzidis, *Nat. Photonics*, 2014, **8**, 489–494.
- 2 N. J. Jeon, J. H. Noh, W. S. Yang, Y. C. Kim, S. Ryu, J. Seo and S. I. Seok, *Nature*, 2015, **517**, 476–480.
- 3 M. M. Lee, J. Teuscher, T. Miyasaka, T. N. Murakami and H. J. Snaith, *Science*, 2012, **338**, 643–647.
- 4 P. Qin, H. Kast, M. K. Nazeeruddin, S. M. Zakeeruddin, A. Mishra, P. Bäuerle and M. Grätzel, *Energy Environ. Sci.*, 2014, **7**, 2981–2985.
- 5 A. Miyata, A. Mitioglu, P. Plochocka, O. Portugall, J. T.-W. Wang, S. D. Stranks, H. J. Snaith and R. J. Nicholas, *Nat. Phys.*, 2015, **11**, 582–587.
- 6 S. D. Stranks, G. E. Eperon, G. Grancini, C. Menelaou, M. J. Alcocer, T. Leijtens, L. M. Herz, A. Petrozza and H. J. Snaith, *Science*, 2013, **342**, 341–344.
- 7 A. A. Zhumekenov, M. I. Saidaminov, M. A. Haque, E. Alarousu, S. P. Sarmah, B. Murali, I. Dursun, X.-H. Miao, A. L. Abdelhady and T. Wu, *ACS Energy Lett.*, 2016, **1**, 32–37.
- 8 M. A. Green, K. Emery, Y. Hishikawa, W. Warta and E. D. Dunlop, *Prog. Photovoltaics: Res. Appl.*, 2016, **24**, 905–913.
- 9 A. Kojima, K. Teshima, Y. Shirai and T. Miyasaka, *J. Am. Chem. Soc.*, 2009, **131**, 6050–6051.
- 10 T. A. Berhe, W.-N. Su, C.-H. Chen, C.-J. Pan, J.-H. Cheng, H.-M. Chen, M.-C. Tsai, L.-Y. Chen, A. A. Dubale and B.-J. Hwang, *Energy Environ. Sci.*, 2016, **9**, 323–356.
- 11 G. Niu, W. Li, F. Meng, L. Wang, H. Dong and Y. Qiu, *J. Mater. Chem. A*, 2014, **2**, 705–710.
- 12 F. Matsumoto, S. M. Vorpahl, J. Q. Banks, E. Sengupta and D. S. Ginger, *J. Phys. Chem. C*, 2015, **119**, 20810–20816.
- 13 J. Yang, B. D. Siempelkamp, D. Liu and T. L. Kelly, *ACS Nano*, 2015, **9**, 1955–1963.
- 14 Y. Han, S. Meyer, Y. Dkhissi, K. Weber, J. M. Pringle, U. Bach, L. Spiccia and Y.-B. Cheng, *J. Mater. Chem. A*, 2015, **3**, 8139–8147.
- 15 M. Ledinsky, P. Löper, B. Niesen, J. Holovsky, S.-J. Moon, J.-H. Yum, S. De Wolf, A. Fejfar and C. Ballif, *J. Phys. Chem. Lett.*, 2015, **6**, 401–406.
- 16 B. Philippe, B.-W. Park, R. Lindblad, J. Oscarsson, S. Ahmadi, E. M. Johansson and H. K. Rensmo, *Chem. Mater.*, 2015, **27**, 1720–1731.
- 17 Y. Li, X. Xu, C. Wang, C. Wang, F. Xie, J. Yang and Y. Gao, *J. Phys. Chem. C*, 2015, **119**, 23996–24002.
- 18 W. Huang, J. S. Manser, P. V. Kamat and S. Ptasincka, *Chem. Mater.*, 2015, **28**, 303–311.
- 19 Y. Li, X. Xu, C. Wang, C. Wang, F. Xie, J. Yang and Y. Gao, *AIP Adv.*, 2015, **5**, 097111.
- 20 X. Zhou, X. Li, Y. Liu, F. Huang and D. Zhong, *Appl. Phys. Lett.*, 2016, **108**, 121601.
- 21 L. Liu, J. A. McLeod, R. Wang, P. Shen and S. Duhm, *Appl. Phys. Lett.*, 2015, **107**, 061904.
- 22 T.-W. Ng, C.-Y. Chan, M.-F. Lo, Z. Q. Guan and C.-S. Lee, *J. Mater. Chem. A*, 2015, **3**, 9081–9085.
- 23 G. R. Kumar, A. D. Savariraj, S. Karthick, S. Selvam, B. Balamuralitharan, H.-J. Kim, K. Viswanathan, M. Vijaykumar and K. Prabakar, *Phys. Chem. Chem. Phys.*, 2016, **18**, 7284–7292.
- 24 J. Emarat, T. Schnier, N. Pourdavoud, T. Riedl, K. Meerholz and S. Olthof, *Adv. Mater.*, 2016, **28**, 553–559.
- 25 M.-C. Jung, Y. M. Lee, H.-K. Lee, J. Park, S. R. Raga, L. K. Ono, S. Wang, M. R. Leyden, B. D. Yu and S. Hong, *Appl. Phys. Lett.*, 2016, **108**, 073901.
- 26 H. Yu, F. Wang, F. Xie, W. Li, J. Chen and N. Zhao, *Adv. Funct. Mater.*, 2014, **24**, 7102–7108.
- 27 H. Xie, X. Liu, L. Lyu, D. Niu, Q. Wang, J. Huang and Y. Gao, *J. Phys. Chem. C*, 2015, **120**, 215–220.
- 28 T. Ding, R. Li, W. Kong, B. Zhang and H. Wu, *Appl. Surf. Sci.*, 2015, **357**, 1743–1746.
- 29 E. M. Miller, Y. Zhao, C. C. Mercado, S. K. Saha, J. M. Luther, K. Zhu, V. Stevanović, C. L. Perkins and J. van de Lagemaat, *Phys. Chem. Chem. Phys.*, 2014, **16**, 22122–22130.
- 30 R. Lindblad, D. Bi, B.-W. Park, J. Oscarsson, M. Gorgoi, H. Siegbahn, M. Odelius, E. M. Johansson and H. K. Rensmo, *J. Phys. Chem. Lett.*, 2014, **5**, 648–653.
- 31 A. Dualeh, P. Gao, S. I. Seok, M. K. Nazeeruddin and M. Grätzel, *Chem. Mater.*, 2014, **26**, 6160–6164.
- 32 M. Liu, M. B. Johnston and H. J. Snaith, *Nature*, 2013, **501**, 395–398.
- 33 R. Lindblad, N. K. Jena, B. Philippe, J. Oscarsson, D. Bi, A. Lindblad, S. Mandal, B. Pal, D. D. Sarma and O. Karis, *J. Phys. Chem. C*, 2015, **119**, 1818–1825.
- 34 B. Conings, L. Baeten, C. De Dobbelaere, J. D'Haen, J. Manca and H. G. Boyen, *Adv. Mater.*, 2014, **26**, 2041–2046.
- 35 W. Zhang, S. Pathak, N. Sakai, T. Stergiopoulos, P. K. Nayak, N. K. Noel, A. A. Haghighirad, V. M. Burlakov, A. Sadhanala and W. Li, *Nat. Commun.*, 2015, **6**, 10030.
- 36 D. Bi, C. Yi, J. Luo, J.-D. Décoppet, F. Zhang, S. M. Zakeeruddin, X. Li, A. Hagfeldt and M. Grätzel, *Nat. Energy*, 2016, **1**, 16142.
- 37 R. Dawood, A. Forty and M. Tubbs, *Proc. R. Soc. A*, 1965, **284**, 272–288.
- 38 H. Seyama and M. Soma, *J. Chem. Soc., Faraday Trans. 1*, 1984, **80**, 237–248.
- 39 N. Fairley, *CasaXPS manual 2.3.15*, Acolyte Science, 2009.
- 40 O. Malinkiewicz, A. Yella, Y. H. Lee, G. M. Espallargas, M. Graetzel, M. K. Nazeeruddin and H. J. Bolink, *Nat. Photonics*, 2014, **8**, 128–132.
- 41 O. Malinkiewicz, C. Roldán-Carmona, A. Soriano, E. Bandiello, L. Camacho, M. K. Nazeeruddin and H. J. Bolink, *Adv. Energy Mater.*, 2014, **4**, 1400345.

