

The effect of degree of polymerization on intra- and interchain micellization of a tail-type cationic polysoap†

Paul A. FitzGerald, David M. McDonald and Gregory G. Warr*

Cite this: *Soft Matter*, 2013, 9, 2711

We have used Reversible Addition Fragmentation chain Transfer (RAFT) to polymerize the T-type surfactant monomer α,ω -methacryloylundecyltrimethylammonium bromide (MUTAB) to various degrees of polymerization, and thereby investigate how its self-assembly is affected. Small-angle neutron scattering (SANS) shows that the interchain aggregation into micelles with an approximately constant number of MUTAB monomer equivalents occurs at low degrees of polymerization, but that micelle elongation occurs when the degree of polymerization exceeds a critical value. In this regime interchain aggregation gives way to intrachain assembly into unimolecular or “unimer” micelles. As with conventional cationic surfactant solutions, addition of salicylate produces long, worm-like micelles containing many amphiphilic polymer chains at all degrees of polymerization. Oscillatory rheology reveals a transition from scission- to reptation-dominated relaxation as increasing polymer chain length also increases the distance between potential scission points. The measured relaxation times lie in the range of hundredths to a few seconds – thus demonstrating the rapidly equilibrating nature of these micellar systems even at the highest degrees of polymerization achieved.

Received 8th November 2012

Accepted 11th January 2013

DOI: 10.1039/c3sm27573f

www.rsc.org/softmatter

Introduction

Attempts to polymerize assemblies of surfactant monomers (surfmers) are found throughout the literature.^{1–14} Typically the aim is to kinetically trap some property of the original micelles, such as size¹ or morphology.^{15,16} This is rarely achieved, and most studies note changes due to the polymerization such as micelle growth.^{15,16} However, it is possible to retain some features of the original micelle such as the cross-sectional radius,^{12,13} and the polymerization is also often found to impart stability to the polymerized micelles against dilution and changes in temperature.^{15,16}

Apart from demonstrated stability to temperature and dilution, there are surprisingly few studies on the stability of polymerized micelles to other classic micelle morphology modifiers such as salt or additional surfactant. Most authors shy away from drawing firm conclusions about the thermodynamic nature of the polymerized structures, whether the final product is a polymer particle or an equilibrium polysoap micelle, or whether they contain one or many polysoap chains.^{15,16} Further, because traditional polymerization methods afford little in the way of controlling polymerization, it has previously not been possible to study the influence of molecular weight on the final structures formed. In this work we use Reversible Addition–

Fragmentation chain Transfer (RAFT) to synthesize a series of polysoaps with varying degree of polymerization, and elucidate the nature and relationship of both polymerized micelles and self-assembled polysoaps as a function of molecular weight.

Recently we have shown that micelles of the classic tail-polymerizable surfactant monomer MUTAB (T-type surfmer, Fig. 1) remain in thermodynamic equilibrium throughout the *in situ* polymerization process.^{17–19} It forms small spheroidal micelles before polymerization ($N_{\text{agg}} = 42$) and larger spheroids after polymerization with a four-fold increase in length, but only a minor increase in the cross-sectional radius. At intermediate polymer conversions, rod-like micelles form with lengths of more than 1500 Å, which collapse into polymerized spheroids as the monomer is consumed. These fully polymerized micelles are stable to both dilution and temperature changes, but remain in dynamic equilibrium in that they undergo facile and extensive structural rearrangement in response to the addition of, for example, salt or monomeric surfactant,^{14,19} or a change in temperature.¹⁴

In this work we compare the structure and properties of *in situ* polymerized aqueous MUTAB micelles with MUTAB polysoaps prepared *ex situ* using RAFT to control the degree of polymerization. We also compare the behaviour of these

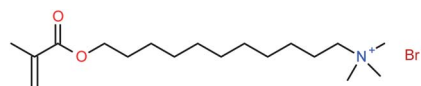


Fig. 1 Tail polymerizable surfactant monomer (surfmer) MUTAB.

School of Chemistry, F11, The University of Sydney, 2006, Australia. E-mail: gregory.warr@sydney.edu.au; Fax: +612 9351 3329; Tel: +612 9351 2106

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c3sm27573f



polymer chains in non-self-assembling and self-assembling solvents, drawing a distinction between inter- and intrachain association to form unimer micelles. Finally, we examine long, worm-like micelles made from polymerized MUTAB with added salicylate, and show how polymerization significantly affects solution rheology through micelle dynamics.

Experimental section

MUTAB was synthesized as described previously.²⁰ The micellar polymerization was performed using UV polymerization at $2 \times$ CMC (~ 1.7 wt% in D_2O), under the same conditions as described previously. Conductivity measurements (see ESI, Fig. S1†) confirm that the polymerized MUTAB micelles do not disassemble on dilution to at least $1/20^{\text{th}}$ of the CMC of monomeric MUTAB, and indicates that micelle counterion binding is similar before and after polymerization.

A series of polymerized MUTAB samples were synthesised to varying degrees of polymerization, n , using the RAFT controlled polymerization technique.^{21,22} RAFT polymerizations of MUTAB were conducted on approximately 30 wt% MUTAB samples in methanol by thermal initiation at 70°C for 16 hours using azoisobutyronitrile (AIBN) initiator and cyanopropyl dithiobenzene (CPDB)^{21,22} RAFT control agent in the mole ratio $0.2 : 1 : n$ (AIBN : CPDB : MUTAB) where $n = 20, 40, 90, 200$ and 500 is the desired degree of polymerization. The polymer mixture was precipitated in ether to remove any residual monomer and dried under vacuum before being redissolved in either D_2O or deuterated methanol for SANS experiments. The RAFT polymers are capped by the small hydrophobic end groups of cyanopropyl and dithiobenzene, which add a negligible hydrophobic contribution to the polymeric backbone.

The actual degree of polymerization for the larger molecular weights was measured using SANS on 5 mg mL^{-1} samples in deuterated methanol brine (containing 0.1 M NaCl to screen any electrostatic interactions), and aqueous micellar solutions in D_2O under various solution conditions (see Results and discussion).

Small-angle neutron scattering (SANS) was performed on the 18 m SANS instrument at the HANARO reactor at the Korean Atomic Energy Research Institute (KAERO) and the NG3 SANS instrument at the NIST Center for Neutron Research (NCNR) at Gaithersburg MD,²³ using 2 mm path length cells with 5.14 \AA (HANARO) or 6.0 \AA (NCNR) neutrons collected on a 128×128 pixel, $640 \times 640\text{ mm}^2$ detector at 1.3 m and 9 m sample to detector distances, giving a combined q range of $0.008\text{--}0.3\text{ \AA}^{-1}$. Data reduction and analysis was performed in Igor Pro (WaveMetrics Inc., version 6.22A) using the NIST reduction and analysis macros (version 4.0).²⁴

SANS data for molecular weight determination was collected at an optimized concentration ($\sim 5\text{ mg mL}^{-1}$), which was large enough to give sufficient contrast yet low enough to exclude almost all polymer-polymer interactions. The optimized concentration was determined by collecting SANS data for the micelle polymerized sample at several concentrations between 1.0 and 20 mg mL^{-1} and then establishing the maximum concentration below which the normalized scattering (*i.e.* $I(q)/c$) overlay each other within experimental error.

Small-angle X-ray Scattering (SAXS) was performed on an Anton Paar SAXSess using 1 mm diameter quartz sample capillaries, and a 1 cm collimated line. Scattering was collected on image plates and desmeared using Otto Glatter's GIFT program to obtain the scattering intensity at $q = 0$ (*i.e.* $I(0)$).²⁵

Rheology data was collected at 25°C on an Anton Paar MCR 302 rheometer using a parallel plate geometry with a 50 mm diameter and 0.5 mm gap. Samples consisted of the polymerized MUTAB in D_2O with sodium salicylate ($0.82\text{ mol}\%$) to induce micelle elongation and network formation.

Results and discussion

M_w and R_g of polyMUTAB in methanol

SANS patterns for the polymerized MUTAB samples in deuterated methanol – with 0.1 M NaCl added to screen any electrostatic interactions – are shown in Fig. 2. For ease of comparison, all intensities are normalized by concentration (w/v) and the backgrounds subtracted. The low q data was fit using the Zimm approximation,²⁶

$$\frac{1}{I(q)} = \frac{1}{I(0)} \left(1 + \frac{R_g^2}{3} q^2 \right) \quad (1)$$

where $I(q)$ is the scattering intensity, q is the scattering vector, R_g is the radius of gyration and $I(0)$ is the scattering at $q = 0$. Assuming negligible interactions between dissolved chains (see Experimental section), the (apparent) molecular weight (M_w) was calculated from $I(0)$ using²⁷

$$\frac{Kc}{I(0)} = \frac{1}{M_w} \quad (2)$$

where c is concentration, and the constant $K = 5.38 \times 10^{-3}\text{ cm}^2\text{ mol}^{-1}\text{ g}^{-2}$ was calculated from $K = v_p^2 \Delta\rho^2 / N_A$,²⁷ where v_p is the partial specific volume of the polymer ($\sim 1.0\text{ cm}^3\text{ g}^{-1}$), $\Delta\rho$ is the

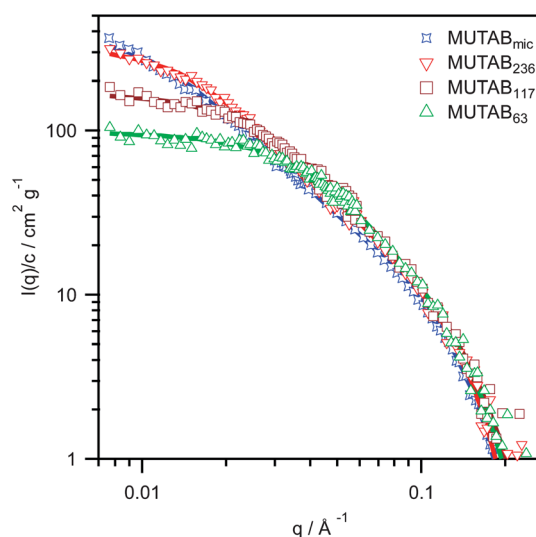


Fig. 2 SANS patterns for polymerized MUTAB samples in deuterated methanol + 0.1 M NaCl with degrees of polymerization of 63 (triangles), 117 (squares), 236 (upside down triangles) and the micelle polymerized sample ($n = 275$, stars). Fits are to the flexible cylinder model.^{30,31}



Table 1 Best-fit parameters for MUTAB_n solutions in d-methanol (Fig. 2) and D₂O (Fig. 4). N_{agg} is the number of polymerized MUTAB chains per micelle, and $N_{\text{agg,mon eq.}}$ is the number of MUTA⁺ monomer equivalents^a

Target n	MUTAB	20	40	90	200	500	Micellar
In d-methanol							
Zimm analysis							
Measured n	1	14	28	63	117	236	275
$R_g/\text{\AA}$	—	—	—	40	64	108	132
Flexible rod model							
Scale	—	—	—	0.0012	0.0012	0.0012	0.0009
Contour Length/ \AA	—	—	—	212	393	584	811
Kuhn Length/ \AA	—	—	—	37	39	82	98
Radius/ \AA	—	—	—	14	13	15	15
In D₂O							
Fits from prolate spheroid model							
Vol. frac.	0.011	0.012	0.017	0.019	0.02	0.019	0.017
$R_a/\text{\AA}$	26	26	27	36	54	79	101
$R_b/\text{\AA}$	14	17	18	19	18	19	19
$N_{\text{agg,mon eq.}}$	42	65	69	106	155	228	305
N_{agg}	42	4.7	2.4	1.7	1.3	1.0	1.1

^a Uncertainties are approximately $\pm 5\%$ for lengths (1–10 \AA) and $\pm 20\%$ for volume fractions.

scattering contrast ($5.69 \times 10^{-10} \text{ cm}^{-2}$),^{28,29} and N_A is the Avogadro constant.

The degree of polymerization, n , was calculated by dividing the measured M_w by the monomer molecular weight (MUTA⁺, 298 g mol⁻¹) and is given in Table 1 with corresponding R_g values from eqn (1) (full fit values given in Table S1, ESI[†]). The degree of polymerization could not be measured for the shorter polymers (target values of 20 and 40) using SANS and were instead measured with SAXS, using MUTAB₆₃ from SANS as a secondary standard.

The experimental degrees of polymerization for RAFT and the micelle polymerized MUTAB, and their corresponding radii of gyration in methanol brine are listed in Table 1. Degrees of polymerization are lower than the target values due to termination prior to complete conversion, with unreacted monomer removed during the precipitation step. Fig. 3 shows that the radii of gyration exhibit a power law dependence of $R_g \sim n^{0.81}$, which is

intermediate between the $n^{0.6}$ dependence expected for long excluded-volume polymers,³² and the $n^{1.0}$ expected for rigid rods.³²

The SANS patterns in methanol were fit to a flexible cylinder model^{30,31} as shown in Fig. 2, with the best-fit parameter values given in Table 1. The model provides good fits over the entire q range, and shows that there are only 5 to 10 Kuhn segments, each 40 to 100 \AA in length, per polymer chain. These are almost an order of magnitude larger than typical methacrylates, *e.g.* 7.2 \AA for poly(methyl methacrylate) in toluene,³³ probably due to electrostatic repulsion between monomers along the polymer chain. Indeed, the 6.2 \AA Debye length in a 0.1 M NaCl methanol solution at 25 °C is large enough to screen interactions between polymer chains 100–200 \AA apart, but not between adjacent charges along the chain, which are only 3–9 \AA apart. The large Kuhn lengths also suggest that the observed $R_g \sim n^{0.81}$ scaling behaviour is actually the chains behaving somewhere between rigid rods and an excluded volume polymer.

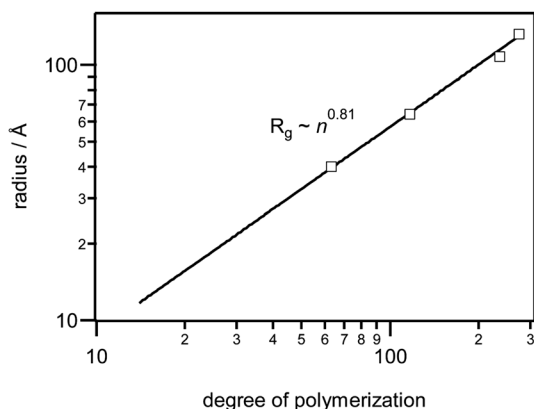


Fig. 3 Polymer chain radii of gyration in d-methanol brine.

polyMUTAB micelle morphologies in D₂O

SANS patterns for 2 wt% RAFT-polymerized MUTAB solutions in D₂O at various degrees of polymerization, n , are shown in Fig. 4, together with both micelle polymerized and monomeric MUTAB at 2xCMC of the monomer (1.7 wt%). The data were best fit as prolate spheroids³⁴ interacting through a screened Coulomb interaction,^{35,36} with key fit values given in Table 1 (full fit values in Table S2, ESI[†]).

The cross-sectional radius, $R_b = 18 \pm 1 \text{ \AA}$, is practically independent of n , (apart from monomeric MUTAB, which has been shown previously to have a smaller radius due to “hair-pinning” of the tails to allow the polar methacrylate group closer to the micelle surface).^{17,18} The polymerized radius is consistent with fully extended non-polar tails, which are



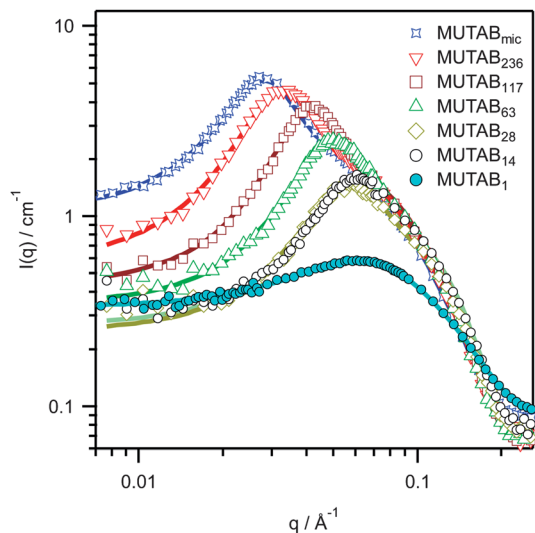


Fig. 4 SANS patterns of micellar solutions of 2 wt% MUTAB_n in D₂O. Solid lines show with fits to a prolate spheroid model with screened Coulomb interaction. Monomer, MUTAB₁, and micelle polymerized, MUTAB_{mic}, at $2 \times \text{CMC} = 1.72 \text{ wt}\%$ are also shown.¹⁸

indifferent to the polymethacrylate backbone and instead dictated by packing constraints. At low n the long axis, R_a , and hence the micelle size, is independent of the *degree of polymerization* *i.e.* $R_a = 26 \pm 1 \text{ \AA}$ at $n = 1, 14$ and 28 . However, for larger n the micelles elongate uniaxially, rapidly approaching the $R_a \sim n^{1.0}$ dependence expected for rigid rods.

The micelle aggregation numbers in monomer equivalents, $N_{\text{agg,mon eq.}}$, obtained by dividing the volume of the spheroidal micelles by the volume of a MUTA⁺ monomer, (496 \AA^3), is insensitive to n up to about 28-mers, but thereafter increases markedly. The reason for this is readily seen by considering the number of actual polyMUTAB units in each micelle, N_{agg} . Both values are listed in Table 1.

At low n , each micelle is formed by the association of several polymer chains (*i.e.* $N_{\text{agg}} > 1$), and the micelle aggregation number is determined by packing constraints. However, once the degree of polymerization of an individual polyMUTAB exceeds the equilibrium micelle aggregation number in monomer equivalents (*i.e.* when $n > N_{\text{agg,mon eq.}}$), the micelle size increases. Beyond this point intra-chain association occurs and the structure is best described as a unimer micelle, in which the average number of polymer chains per micelle is unity (see Table 1). The alkyl chain length still provides an effective packing constraint in the radial direction, so the micelles form increasingly elongated, rod-like structures. Thus it is packing constraints that cause the growth in the micelle rather than the effect of the polymer backbone conformation.

Viscoelastic behaviour with salicylate counterions

Addition of sodium salicylate (NaSal) to both monomeric and micelle-polymerized MUTAB samples produces long, worm-like micelles.¹⁹ This effect of certain aromatic anions is well known for conventional cationic surfactant systems such as

cetylpyridinium chloride, where strong binding of the salicylate produces rod like micelles and a viscoelastic solution.^{37,38} In light of the results above, we interpret this as a shift in the threshold for unimer micelle formation in the presence of salicylate to much higher n .

In polyMUTAB this transformation is accompanied by a dramatic increase in solution viscosity, although micelles of monomeric MUTAB are little affected. This can clearly be seen in Fig. 5, which shows the steady shear viscosity for 1.5 wt% solutions of MUTAB_n : NaSal (55 : 45) at various degrees of polymerization. All samples are shear thinning, with the limiting low-shear viscosity increasing with degree of polymerization. Relaxation times, taken as the inverse of the shear rate at which the viscosity falls to half its initial value, also increase markedly with n . These are listed in Table 2.

Storage and loss moduli obtained from oscillatory shear measurements are shown in Fig. 6, and are similar to those reported previously for entangled networks of worm-like micelles.^{37,39} [It was not possible to measure oscillatory rheology on low n samples, or the salicylate-free samples, because of their low viscosity.] Oscillatory data were fit using the Maxwell model⁴⁰ to extract the plateau modulus and terminal relaxation time, G_0 and τ , also listed in Table 2. Although the fits are not particularly good in Fig. 6c and d, the terminal relaxation time, τ , can still be unambiguously determined from the crossover frequency ($G' = G''$). As observed under steady shear, the

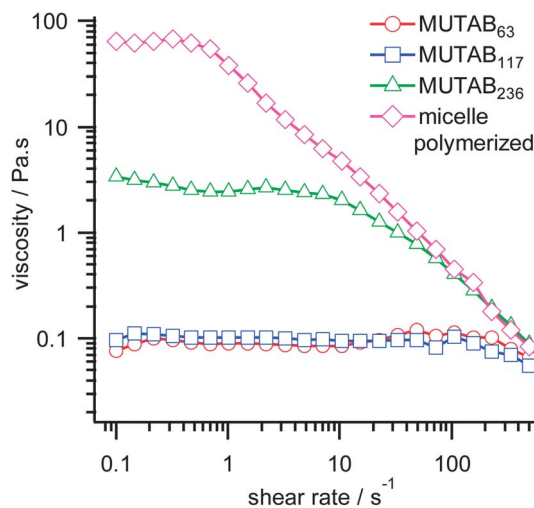


Fig. 5 Steady shear viscosity of 1.5 wt% MUTAB_n : NaSal (55 : 45 mole ratio) samples in D₂O.

Table 2 Relaxation times estimated from steady shear data, τ_s , and fit values from Maxwell model, τ and G_0

Degree of polymerization, n	τ_s/s (steady shear)	τ/s (crossover frequency)	G_0/Pa	$d/\text{\AA}$
63	—	0.016	2.6	1170
117	0.0019	0.019	2.9	1120
236	0.04	0.27	2.8	1140
275	0.8	2.3	6.2	872



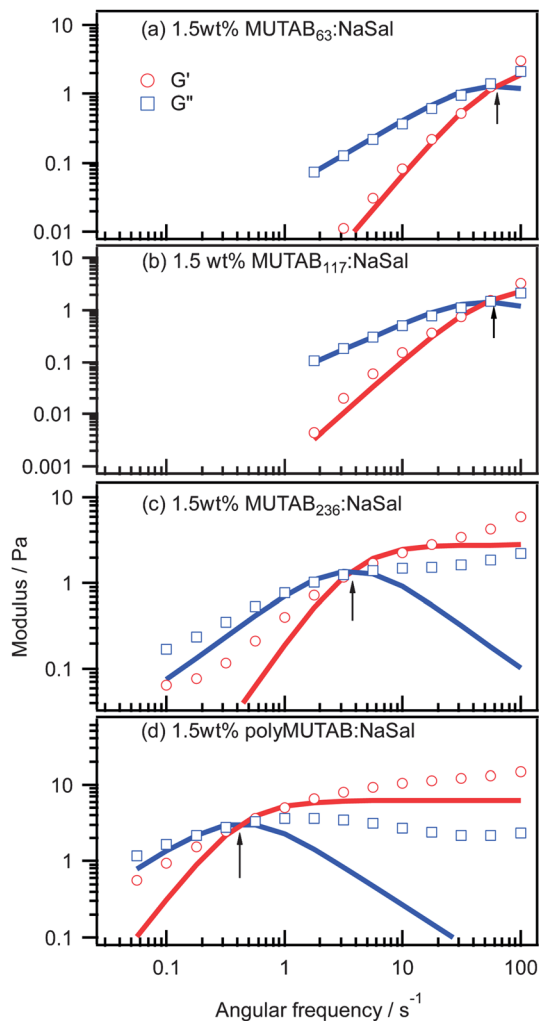


Fig. 6 Oscillatory rheology on 1.5 wt% MUTAB_n : NaSal (55 : 45 mole ratio) samples in D_2O . Arrows indicate crossover frequency.

relaxation time increases with degree of polymerization up to and including the micelle polymerized system ($n = 275$). Relaxation times in these systems are not more than a few seconds even for the largest degree of polymerization.

Relaxation in such living polymer systems occurs by a combination of reptation and chain scission.⁴¹ If a system is reptation dominated (*i.e.* $\tau_{\text{rep}} \ll \tau_{\text{break}}$) then it can be treated as a classical polymer which, for worm-like micelles yields an exponential relaxation spectrum.⁴¹ If chain scission is fast (*i.e.* $\tau_{\text{rep}} > \tau_{\text{break}}$) then two regions are obtained for the oscillatory rheology data: $t \geq \tau_{\text{break}}$ (low frequency) yields Maxwellian behaviour with a terminal relaxation time $\tau \sim (\tau_{\text{rep}}\tau_{\text{break}})^{1/2}$, and $t < \tau_{\text{break}}$ (high frequency) which is dominated by Rouse motion and produces a minimum in G'' .

At low degrees of polymerization, *i.e.* MUTAB_{63} : NaSal and MUTAB_{117} : NaSal, both G' and G'' are well-described by the Maxwell model in the accessible frequency range. This indicates single relaxation times and stress relaxation dominated by micelle breaking. For longer n , *i.e.* MUTAB_{236} : NaSal and the $\text{MUTAB}_{\text{mic}}$: NaSal systems, the terminal relaxation times increase by two orders of magnitude, and the Maxwell model

fails to describe G' and G'' even at low frequencies. This indicates a transition to reptation dominated stress relaxation and the emergence of a relaxation time spectrum.

The plateau modulus, G_0 , changes little with degree of polymerization (Table 2). This is not surprising for an entangled network of long, worm-like micelles, for which $G_0 = kT/d^3$, where d is the distance between entanglement points.^{42,43} Here d is approximately 1100 Å for all measured RAFT- MUTAB_n : NaSal polymer systems, and 870 Å for the micelle polymerized $\text{MUTAB}_{\text{mic}}$: NaSal system. This also suggests that the observed change in low-shear viscosity is a direct consequence of increasing relaxation time.

As seen from the structures of MUTAB micelles, above a degree of polymerization around 30–60, individual polyMUTA⁺ chains undergo intrachain association into short rods or prolate spheroidal unimer micelles. Previous SANS studies, together with the observation of entanglement distances, d , much larger than the dimension of any polyMUTAB chains or unimer micelles confirms that addition of salicylate leads once again to multiple chains associating into long, worm-like micelles. By increasing the degree of polymerization of polyMUTA⁺ chains, the number of potential scission points within each micelle is decreased. This increases τ_{break} , and gives rise to a transition from scission- to reptation-dominated stress relaxation, beyond which the micelles behave like conventional polymers.

Conclusions

We have synthesised a series of polysoaps to different degrees of polymerization using RAFT controlled polymerization of the T-type cationic MUTAB surfmer and compared their self-assembly structure and dynamics with an *in situ* micelle polymerized sample.

The assembly of these polysoaps remains subject to traditional surfactant packing constraints. At low degrees of polymerization, interchain association of several MUTAB_n polysoap molecules in aqueous solution form micelles with an approximately constant number of MUTAB monomer equivalents. However, when the degree of polymerization reaches and ultimately exceeds that aggregation number, the assembly transitions into larger unimer micelles. These contain a single polyMUTAB chain formed by intrachain association of its alkyl tails packed into an elongated spheroid or short, rod-like structure with a cross-sectional radius determined by the MUTA⁺ alkyl chain length.

In the presence of salicylate, these chains assemble into long, worm-like micelles whose dynamics depend strongly on the degree of polymerization of the polysoap. This leads to a transition from scission- to reptation-dominated stress relaxation with increasing degree of polymerization, allowing their terminal relaxation time and zero-shear viscosity to be varied by over two orders of magnitude.

Acknowledgements

We gratefully acknowledge the Australian Institute of Nuclear Science and Engineering (AINSE) for travel funding, Dr Suraj



Sharma and Dr Shao Cong Dai for assistance with the rheology measurements, and Professor Sebastien Perrier and Dr Khwanrat Chatjaroenporn for assistance with the RAFT polymerizations. We acknowledge the support of the National Institute of Standards and Technology, U.S. Department of Commerce, and the Korea Atomic Energy Research Institute (KAERI) in providing the neutron research facilities used in this work. This work utilized facilities supported in part by the National Science Foundation under Agreement no. DMR-0944772. This work was supported by Discovery Grants from the Australian Research Council.

References

- 1 A. J. Hyde and D. J. M. Robb, *J. Phys. Chem.*, 1963, **67**, 2089–2092.
- 2 C. E. Larrabee and E. D. Sprague, *J. Polym. Sci., Polym. Lett. Ed.*, 1979, **17**, 749–751.
- 3 S. M. Hamid and D. C. Sherrington, *Br. Polym. J.*, 1984, **16**, 39–45.
- 4 C. M. Paleos, *Chem. Soc. Rev.*, 1985, **14**, 45–67.
- 5 S. Hamid and D. Sherrington, *J. Chem. Soc., Chem. Commun.*, 1986, 936–938.
- 6 P. Dais, C. M. Paleos, G. Nika and A. Malliaris, *Makromol. Chem.*, 1993, **194**, 445–450.
- 7 D. Cochon, F. Candau and R. Zana, *Macromolecules*, 1993, **26**, 5755–5764.
- 8 D. Cochon, R. Zana and F. Candau, *Macromolecules*, 1993, **26**, 5765–5771.
- 9 S. R. Kline, *Langmuir*, 1999, **15**, 2726–2732.
- 10 S. Kline, *J. Appl. Crystallogr.*, 2000, **33**, 618–622.
- 11 M. Abe, K. Tsubone, T. Koike, K. Tsuchiya, T. Ohkubo and H. Sakai, *Langmuir*, 2006, **22**, 8293–8297.
- 12 S. Liu, Y. I. Gonzalez, D. Danino and E. W. Kaler, *Macromolecules*, 2005, **38**, 2482–2491.
- 13 Z. Y. Zhu, Y. I. Gonzalez, H. X. Xu, E. W. Kaler and S. Y. Liu, *Langmuir*, 2006, **22**, 949–955.
- 14 P. A. FitzGerald and G. G. Warr, *Adv. Colloid Interface Sci.*, 2012, **179**, 14–21.
- 15 M. Summers and J. Eastoe, *Adv. Colloid Interface Sci.*, 2003, **100–102**, 137–152.
- 16 H.-P. Hentze and E. W. Kaler, *Curr. Opin. Colloid Interface Sci.*, 2003, **8**, 164–178.
- 17 K. Chatjaroenporn, R. W. Baker, P. A. FitzGerald and G. G. Warr, *J. Colloid Interface Sci.*, 2009, **336**, 449–454.
- 18 K. Chatjaroenporn, R. W. Baker, P. A. FitzGerald and G. G. Warr, *Langmuir*, 2010, **26**, 11715–11719.
- 19 P. A. FitzGerald, K. Chatjaroenporn, X. L. Zhang and G. G. Warr, *Langmuir*, 2011, **27**, 11852–11859.
- 20 J. Michas, C. M. Paleos and P. Dais, *Liq. Cryst.*, 1989, **5**, 1737–1745.
- 21 T. P. Le, G. Moad, E. Rizzardo and S. H. Thang, WO9801478A1, 1998.
- 22 J. Chiefari, Y. K. Chong, F. Ercole, J. Krstina, J. Jeffery, T. P. T. Le, R. T. A. Mayadunne, G. F. Meijs, C. L. Moad, G. Moad, E. Rizzardo and S. H. Thang, *Macromolecules*, 1998, **31**, 5559–5562.
- 23 C. J. Glinka, J. G. Barker, B. Hammouda, S. Krueger, J. J. Moyer and W. J. Orts, *J. Appl. Crystallogr.*, 1998, **31**, 430–445.
- 24 S. R. Kline, *J. Appl. Crystallogr.*, 2006, **39**, 895–900.
- 25 O. Glatter, *J. Appl. Crystallogr.*, 1977, **10**, 415–421.
- 26 J. P. Cotton, in *Neutron, X-ray and Light Scattering*, ed. P. Lindner and T. Zemb, Elsevier Science Publishers, Amsterdam, 1991, p. 11.
- 27 P. Lindner and T. Zemb, *Neutrons, X-rays and light: scattering methods applied to soft condensed matter*, Elsevier Science, Amsterdam, 2002.
- 28 A. Munter, <http://www.ncnr.nist.gov/resources/sldcalc.html>, accessed 30 October 2012.
- 29 V. F. Sears, *Neutron News*, 1992, **3**, 26–37.
- 30 J. S. Pedersen and P. Schurtenberger, *Macromolecules*, 1996, **29**, 7602–7612.
- 31 W. R. Chen, P. D. Butler and L. J. Magid, *Langmuir*, 2006, **22**, 6539–6548.
- 32 P.-G. de Gennes, *Scaling Concepts in Polymer Physics*, Cornell University Press, London, 1979.
- 33 R. Kirste and O. Kratky, *Z. Phys. Chem.*, 1962, **31**, 363–374.
- 34 *Structure Analysis by Small-angle X-ray and Neutron Scattering*, ed. L. A. Feigin and D. I. Svergun, Plenum Press, New York, 1987.
- 35 J.-P. Hansen and J. B. Hayter, *Mol. Phys.*, 1982, **46**, 651–656.
- 36 J. B. Hayter and J. Penfold, *Mol. Phys.*, 1981, **42**, 109–118.
- 37 H. Rehage and H. Hoffmann, *J. Phys. Chem.*, 1988, **92**, 4712–4719.
- 38 M. E. Cates and S. J. Candau, *J. Phys.: Condens. Matter*, 1990, **2**, 6869–6892.
- 39 A. Khatory, F. Lequeux, F. Kern and S. J. Candau, *Langmuir*, 1993, **9**, 1456–1464.
- 40 J. D. Ferry, *Viscoelastic Properties of Polymers*, Wiley, New York, 1961.
- 41 M. E. Cates, *Macromolecules*, 1987, **20**, 2289–2296.
- 42 M. Doi and S. F. Edwards, *The Theory of Polymer Dynamics*, Clarendon Press, Oxford, 1986.
- 43 F. Kern, F. Lequeux, R. Zana and S. J. Candau, *Langmuir*, 1994, **10**, 1714–1723.

